

# Unitarity Triangle Angles Explained: a Predictive New Quark Mass Matrix Texture

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### Outline of Talk



- Mysteries of the quark mass and mixing spectra
- Weak Interaction flavour structure
- SM origins of masses and mixings
- Historical efforts to explain

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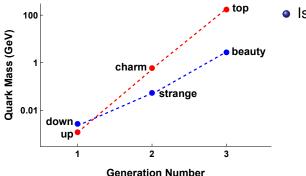


- Mysteries of the quark mass and mixing spectra
- Weak Interaction flavour structure
- SM origins of masses and mixings
- Historical efforts to explain
- Mysteries of the Unitarity Triangle
- The new mass matrix texture
- Confronting the data
- Symmetries of the texture
- Discussion and conclusions

## Mystery of Quark Mass Spectra



• Quark masses show marked hierarchical structure:



Is quasi-"geometric":

$$\frac{m_c}{m_t} \simeq 0.0035$$

$$\frac{m_u}{m_t} \sim 0.0020$$

$$\frac{m_u}{m_c} \simeq 0.0020$$

$$\frac{m_s}{m_b} \simeq 0.020$$

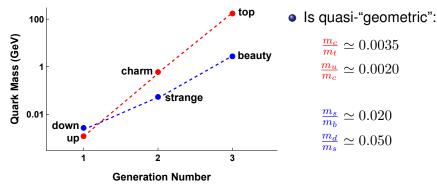
$$\frac{m_d}{m_s} \simeq 0.050$$

Noted by very many authors

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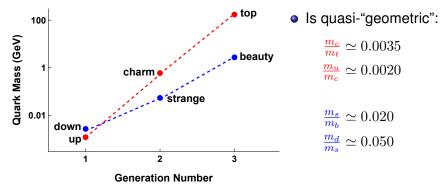


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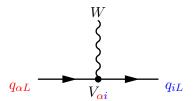


- Noted by very many authors
- Masses not predicted in the SM
- Hierarchy certainly not explained within SM
- BSM, Froggatt-Neilsen mechanism has had some success



• CKM quark mixing matrix:

$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$





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$$\sim \begin{pmatrix} 1 & \lambda & A\lambda^3(\overline{\rho} - i\overline{\eta}) \\ -\lambda & 1 & A\lambda^2 \\ A\lambda^3(1 - \overline{\rho} - i\overline{\eta}) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4),$$

where 
$$\lambda \equiv |V_{us}| \simeq 0.22$$
  
 $A, \overline{\rho} \text{ and } \overline{\eta} \lesssim \mathcal{O}(1)$ 



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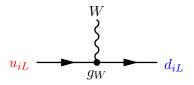
- Elements not predicted by the SM
- Strong hierarchy certainly not explained within SM
- But masses and mixings both arise in the Yukawa/Mass matrices

### The Weak Interaction



- In the gauge theory
- 3 generations of quarks:

$$\begin{pmatrix} u_1 \\ d_1 \end{pmatrix}_L \quad \begin{pmatrix} u_2 \\ d_2 \end{pmatrix}_L \quad \begin{pmatrix} u_3 \\ d_3 \end{pmatrix}_L$$

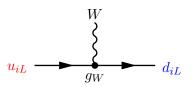


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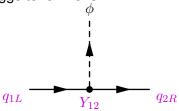
- ullet Write  $\underline{u}_w = (u_1, u_2, u_3)^T$  and  $\underline{d}_w = (d_1, d_2, d_3)^T$
- $W^{\pm}$  couplings initially flavour-diagonal:

$$\mathcal{L}_W \sim g_W \, \underline{\bar{u}_{wL}} \cdot \underline{d}_{wL} \, W^+ + H.C.$$

# Masses and Mixings from SM



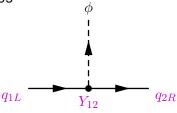
 Fermion Masses and Mixings have common origin in (Yukawa) couplings of the Higgs to fermions



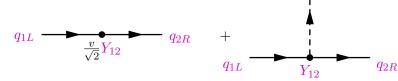
# Masses and Mixings from SM



 Fermion Masses and Mixings have common origin in (Yukawa) couplings of the Higgs to fermions



- After SSB,  $\phi \to \frac{v}{\sqrt{2}} + H$
- Diagram splits to give:



Н

## Taking all Permutations



Recall

$$\underline{u}_w = (u_1, u_2, u_3)^T$$
 and  $\underline{d}_w = (d_1, d_2, d_3)^T$ 

After SSB, Lagrangian for the quark masses is (dropping L/R labels):

$$\mathcal{L}_{Mass} \sim \frac{v}{\sqrt{2}} \, \underline{\bar{u}}_w \cdot Y_u \cdot \underline{u}_w + \frac{v}{\sqrt{2}} \, \underline{\bar{d}}_w \cdot Y_d \cdot \underline{d}_w$$

Identify

$$\frac{v}{\sqrt{2}}Y_u \equiv M_u$$
 and  $\frac{v}{\sqrt{2}}Y_d \equiv M_d$ 

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 and  $\frac{v}{\sqrt{2}}Y_d \equiv M_d$ 

- ullet  $M_u$  and  $M_d$  are clearly not diagonal
- Can choose basis where they are Hermitian without observable consequences.

## **Physical Particles?**



- Identified as eigenstates of  $M_u$  and  $M_d$
- So, diagonalise to find them:

$$(u,c,t)^T \equiv \underline{u} = U_u \cdot \underline{u}_w$$
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Then (chiral labels dropped):

$$\mathcal{L}_{M+W} = \underline{\bar{u}} \cdot D_{u} \cdot \underline{u} + \underline{\bar{d}} \cdot D_{d} \cdot \underline{d} + g_{W} \, \underline{\bar{u}} \cdot U_{u} \cdot U_{d}^{\dagger} \cdot \underline{d} \, W^{+} + \dots$$

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where

$$egin{aligned} D_u &\equiv U_u \cdot M_u \cdot U_u^{\dagger} = \mathrm{diag}(m_u, m_c, m_t) \ D_d &\equiv U_d \cdot M_d \cdot U_d^{\dagger} = \mathrm{diag}(m_d, m_s, m_b) \ &= V_{CKM} \equiv U_u \cdot U_d^{\dagger} \ &= \mathrm{unitary}. \end{aligned}$$

Thus masses and mixings both originate in the MMs

### Historical Context



- So, can mass ratios and mixings be related?
- One way is with "texture zeroes" pioneering idea by Harald Fritzsch (1976-78)
- E.g.  $M_d \equiv M_F(m_b, a_d, b_d)$

$$= m_b \begin{pmatrix} 0 & a_d & 0 \\ a_d^* & 0 & b_d \\ 0 & b_d^* & 1 \end{pmatrix} : \text{diagonalise} \Rightarrow \begin{vmatrix} |a_d| = \sqrt{\frac{m_d m_s}{m_b m_b}} \\ |b_d| = \sqrt{\frac{m_s}{m_b}} \sim 0.1044$$

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Diagonalised by:

$$U_d^{\dagger} \sim \begin{pmatrix} 1 & s_1^d & s_1^d s_2^d \\ -s_1^d & 1 & s_2^d \\ 0 & -s_2^d & 1 \end{pmatrix} \text{ with } \begin{aligned} s_1^d &\equiv \sin \theta_{12}^d \simeq \sqrt{\frac{m_d}{m_s}} \sim 0.224 \\ s_2^d &\equiv \sin \theta_{23}^d \simeq \sqrt{\frac{m_s}{m_b}} \sim 0.14 \end{aligned}$$

Already somewhat encouraging.

### Fritzsch Texture Prediction for $V_{us}$



- ullet BUT  $s_2^d$ ,  $s_3^d$  too big, AND should treat  $M_u$  and  $M_d$  alike
- Do by writing  $M_u = M_F(m_t, a_u, b_u)$  [NB. 8 params for 10 obs  $\checkmark$  ]
- Since

$$V_{CKM} = U_u U_d^{\dagger}(\dagger \Rightarrow \text{ inverse for unitary matrix}),$$

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In complex case, phase enters (gives CP-violation):

$$\tilde{\delta} \equiv arg(\mathbf{a_u}) - arg(\mathbf{a_d})$$

Together give:

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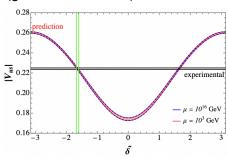
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Good fit ✓





• Here, another phase enters:

$$\bar{\beta} \equiv arg(\frac{\mathbf{b_u}}{\mathbf{a}}) - arg(\frac{\mathbf{b_d}}{\mathbf{a}})$$

One finds:

$$\begin{split} V_{cb} \sim A\lambda^2 &= |s_2^d - s_2^u e^{i\bar{\beta}}| \\ &= |\sqrt{\frac{m_s}{m_b}} - \sqrt{\frac{m_c}{m_t}} e^{i\bar{\beta}}| \end{split}$$

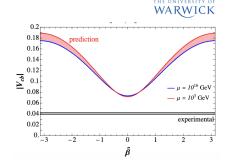
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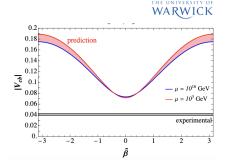
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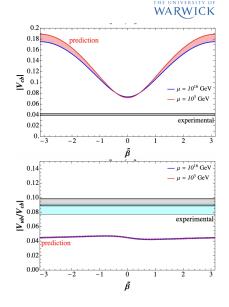
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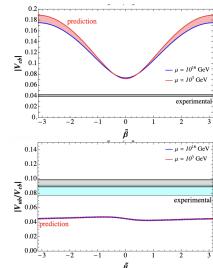
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 Figs from B. Belfatto and
 Z. Berezhiani, arXiv: 2305.00069. Recent approach to revive Fritzsch using non-Hermitian MMs (but 10 pars )



## The Unitarity Triangle



$$V_{CKM} \equiv \begin{matrix} V_u U_d^{\dagger} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \text{ is unitary}.$$

ie. complex dot-product of every pair of columns (or rows) is zero.
 E.g.

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$

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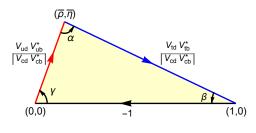


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 E.g.

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•  $\Rightarrow$  triangle in complex plane (normalise by  $1/|V_{cd}V_{cb}^*|$ ):



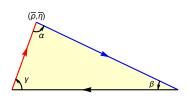
- Base length unity
- 2 parameters, choose:
   2 angles or
   top vertex = ρ + iη
- Area =  $\frac{1}{2}\overline{\eta}$
- All CP-violating observables  $\propto$  Area

# Mysteries of the Unitarity Triangle



- Sides/Angles of UT are arbitrary in SM
- But measured angles:

$$\alpha = (91.6 \pm 1.4)^{\circ}$$
  
 $\beta = (22.6 \pm 0.4)^{\circ}$   
 $\gamma = (65.7 \pm 1.3)^{\circ}$ 



### consistent with "special" values:

$$(\alpha, \beta, \gamma) \simeq (\frac{\pi}{2}, \frac{\pi}{8}, \frac{3\pi}{8}) \equiv (\alpha_0, \beta_0, \gamma_0).$$

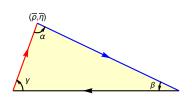
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- Seems striking!
- Coincidence or smoking gun?
- → Test as clue to what lies behind.



### **Build Special Angles into a Texture**



$$M_q^{HS} \equiv n_q \begin{pmatrix} c_q \lambda_q^4 & b \lambda_q^3 & 0 \\ b \lambda_q^{*3} & b \lambda_q^2 & A_0 \lambda_q^2 \\ 0 & A_0 \lambda_q^{*2} & 1 \end{pmatrix}, \quad \begin{matrix} q = \mathbf{u}, \mathbf{d}, \\ \lambda_q \text{ complex} \\ \arg(\lambda_q) \text{ unobservable} \\ (A_0, b, c_u, c_d) \leq \mathcal{O}(1) \end{matrix}$$

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Complex ratio is fixed constant:

$$\frac{\lambda_u}{\lambda_d} \equiv -i \tan \frac{\pi}{8}$$

Controls angles of the UT (see later)

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- $ho |\lambda_u/\lambda_d| \simeq 0.41$  controls relative strength of "u" and "d" mass hierarchies
- Complex sum is fitted parameter close to  $\lambda$ :

$$|\boldsymbol{\lambda_d} + \boldsymbol{\lambda_u}| \equiv \lambda_0 = \boldsymbol{\lambda} + \mathcal{O}(\boldsymbol{\lambda}^3).$$



## Build Special Angles into a Texture



$$M_q^{HS} \equiv n_q \begin{pmatrix} c_q \lambda_q^4 & b \, \boldsymbol{\lambda_q^3} & 0 \\ b \, \boldsymbol{\lambda_q^{*3}} & b \lambda_q^2 & A_0 \boldsymbol{\lambda_q^2} \\ 0 & A_0 \boldsymbol{\lambda_q^{*2}} & 1 \end{pmatrix}, \quad \boldsymbol{\lambda_q} \text{ complex}$$

$$(A_0, b, c_u, c_d) \lesssim \mathcal{O}(1)$$

Complex ratio is fixed constant:

$$\frac{\frac{\lambda_u}{\lambda_d} \equiv -i \tan \frac{\pi}{8}}{}$$

- Controls angles of the UT (see later)
- ightharpoonup arg  $\lambda_u/\lambda_d=-i$ , is sole source of CP violation
- $|\lambda_u/\lambda_d| \simeq 0.41$  controls relative strength of "u" and "d" mass hierarchies
- Complex sum is fitted parameter close to  $\lambda$ :

$$|\lambda_d + \lambda_u| \equiv \lambda_0 = \lambda + \mathcal{O}(\lambda^3).$$

Describes 10 observables with 7 real parameters



# Leading-order Solution (Quark Masses)



■ Diagonalise → masses:

$$D_q = U_q M_q^{HS} U_q^{\dagger} = m_3^q \begin{pmatrix} (c_q - b)\lambda_q^4 & 0 & 0 \\ 0 & b\lambda_q^2 & 0 \\ 0 & 0 & 1 \end{pmatrix}, \quad q = \mathbf{u}, \mathbf{d},$$

• Good for mass hierarchy  $(\frac{\lambda_u}{\lambda_u}, \frac{\lambda_d}{\lambda_d} << 1)$   $\checkmark$ 

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- 3 free parameters (at LO):  $b, c_u, c_d$  (to fit 4 mass ratios)
- ⇒ one constraint/prediction (LO):

$$\frac{m_c}{m_t} \frac{m_b}{m_s} = \left| \frac{\lambda_u}{\lambda_d} \right|^2 = \tan^2 \frac{\pi}{8} = \begin{cases} 0.172 \ (LO) \\ 0.176 \ (NLO) \end{cases} \text{ c.f. } 0.177 \pm 0.002 \ (exp) \checkmark$$

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• Fits any  $m_u$ ,  $m_d \checkmark$  (no prediction here).



## Leading-order Solution (Quark Mixing)



- ullet Diagonalised by  $2 \times 2$  (complex) rotations in 23 and 12 spaces.
- Small entries induced in the 13 elements of  $U_q$ :

$$U_q \simeq \begin{pmatrix} 1 & \pm \lambda_q & A_0 \lambda_q^3 \\ \mp \lambda_q^* & 1 & -A_0 \lambda_q^2 \\ 0 & A_0 \lambda_q^{*2} & 1 \end{pmatrix}, \quad q = \mathbf{u}, \mathbf{d}.$$

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• Combine  $U_u$  and  $U_d$ :

$$\Rightarrow V_{CKM} = \mathbf{U_u} \mathbf{U_d}^{\dagger} \simeq \begin{pmatrix} 1 & \lambda_0 & A_0 \lambda_0^2 \mathbf{\lambda_u} \\ -\lambda_0 & 1 & A_0 \lambda_0^2 \\ A_0 \lambda_0^2 \mathbf{\lambda_d^*} & -A_0 \lambda_0^2 & 1 \end{pmatrix}$$

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C.f. Wolfenstein form:

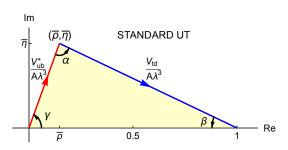
$$V_{CKM} = \begin{pmatrix} 1 & \lambda & A\lambda^3(\overline{\rho} - i\overline{\eta}) \\ -\lambda & 1 & A\lambda^2 \\ A\lambda^3(1 - \overline{\rho} - i\overline{\eta}) & -A\lambda^2 & 1 \end{pmatrix} \Rightarrow \begin{cases} \lambda \simeq \lambda_0 \checkmark \\ A \simeq A_0 \checkmark \\ (\overline{\rho} + i\overline{\eta}) \simeq \frac{\lambda_u^*}{\lambda_0} \end{cases}$$

## The UT Angles



#### • Have deduced that:

$$egin{aligned} oldsymbol{\lambda_u^*} &\simeq \lambda(\overline{
ho} + i\overline{\eta}) = rac{V_{ub}^*}{A\lambda^2} \ oldsymbol{\lambda_d^*} &\simeq \lambda(1 - \overline{
ho} - i\overline{\eta}) = rac{V_{td}}{A\lambda^2} \ &\Rightarrow \gamma \simeq \arg oldsymbol{\lambda_u^*} \ eta \simeq \arg oldsymbol{\lambda_d} \ & ext{and} \ lpha \simeq \arg (-rac{oldsymbol{\lambda_u}}{oldsymbol{\lambda_d}}) \end{aligned}$$



## The UT Angles

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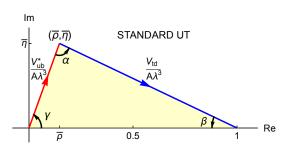
• Have deduced that:

$$\lambda_{u}^{*} \simeq \lambda(\overline{\rho} + i\overline{\eta}) = \frac{V_{ub}^{*}}{A\lambda^{2}}$$

$$\lambda_{d}^{*} \simeq \lambda(1 - \overline{\rho} - i\overline{\eta}) = \frac{V_{td}}{A\lambda^{2}}$$

$$\Rightarrow \gamma \simeq \arg \lambda_{u}^{*}$$

$$\beta \simeq \arg \lambda_{d}$$



- Recall, HS texture asserts  $\frac{\lambda_u}{\lambda_d} = -i \tan \frac{\pi}{8}$ 
  - ightharpoonup  $\Rightarrow$   $\alpha \simeq \frac{\pi}{2}$   $\checkmark$

and  $\alpha \simeq \arg(-\frac{\lambda_u}{\lambda_u})$ 

- $ightharpoonup \Rightarrow \tan \beta = \left| \frac{\lambda_u}{\lambda_d} \right|$  (see Figure).
- $ightharpoonup 
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- Data from PDG
- Renormalise to common scale  $(\mu = m_t)$
- Fit using full numerical diagonalisation



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- $\rightarrow$  poor fit:  $\chi^2/dof \simeq 100/3!$
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- Disaster?
- Not necessarily!
- Because these quantities "run" with renormalisation scale
- ~ 13% from weak to GUT scales:  $A(\uparrow), m_c/m_t(\uparrow)$  and  $m_s/m_b(\downarrow)$ .
- $\lambda$ ,  $\alpha$ ,  $\beta$ ,  $m_u/m_c$  and  $m_d/m_s$  are  $\sim$  invariant.
- $\bullet \Rightarrow \text{vary } \mu$

## Some Details of the Fit



- Fit  $\chi^2/\text{d.o.f} \simeq 1.01/2$
- Best fit renormalisation scale:

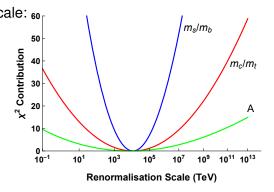
$$\mu \sim (0.3 \rightarrow 3) \times 10^4 \text{ TeV}$$

- Fitted values of the free parameters:
  - $\lambda_0 = 0.22646$
  - $A_0 = 0.854$
  - b = 0.462
  - $c_u = 0.344$
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- Three curves minimise at common scale ~ 10<sup>4</sup> TeV



## **Fit Predictions**



Observable	Input Renormali-	Fitted Value
	sed to $\mu = 10^4$ TeV	at $\mu = 10^4$ TeV
$ m_u/m_c  (\times 10^3)$	$2.00 \pm 0.05$	2.00
$\left  m_d / m_s \right  (\times 10^2)$	$4.97 \pm 0.06$	4.97
$\frac{m_c/m_t}{(\times 10^3)}$	$3.46 \pm 0.03$	3.46
$m_s/m_b  (\times 10^2)$	$1.968 \pm 0.008$	1.968
$\lambda$	$0.2250 \pm 0.0007$	0.2250
A	$0.88 \pm 0.02$	0.88
$\overline{ ho}$	$0.159 \pm 0.009$	0.152
$\overline{\eta}$	$0.352 \pm 0.007$	0.348
UT Angles		Prediction from Fit
$\frac{\alpha}{\alpha}$ (°)	$91.6 \pm 1.4$	$91.30 \pm 0.02$
eta (°)	$22.6 \pm 0.4$	$22.3 \pm 0.1$
$\gamma$ (°)	$65.7 \pm 1.3$	$66.4 \pm 0.1$

Fitted values in table are predictions



## The Leading Order UT (LO-UT)



Define useful complex constants:

$$\mathbf{z_0} \equiv \mathbf{\lambda_u^*}/\lambda_0 = is_0 e^{-i\beta_0} = \rho_0 + i\eta_0,$$

$$\mathbf{\overline{z_0}} \equiv \mathbf{\lambda_d^*}/\lambda_0 = c_0 e^{-i\beta_0} = 1 - \mathbf{z_0},$$

#### where

$$s_0 \equiv \sin \beta_0; \quad c_0 \equiv \cos \beta_0; \quad \eta_0 = s_0 c_0 = \frac{1}{2\sqrt{2}} \quad \text{and} \quad \rho_0 = s_0^2.$$

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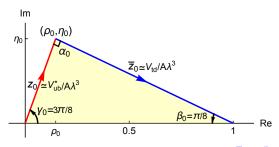
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Use to construct LO-UT



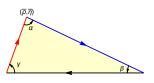
# Symmetries of the $M_q^{HS}$



- Properties of the *paired system*  $(M_u, M_d)$ , rather than of either in isolation
- Could be viewed as consequence of forms, or, preferably, as ab initio symmetries which constrain  $(M_u, M_d)$  forms
- Outlined below

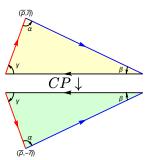


- *CP*:
- Under CP, all complex numbers in the MMs are complex-conjugated



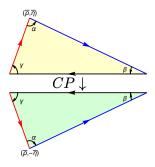


- *CP*:
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- Observable effect is to flip orientation of UT in complex plane  $(\overline{\eta} \to -\overline{\eta})$
- Unless  $\overline{\eta} = 0$  (CP is conserved)

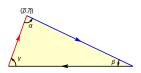




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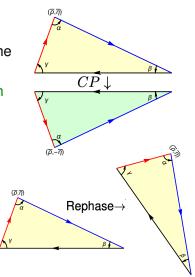


- Rephasing:
- Simultaneous *phase changes* of  $M_u$  and  $M_d$  unobservable



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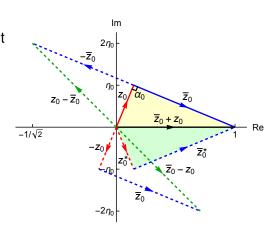
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- Rephasing:
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- UT simply rotates in complex plane
- (Physical) shape and size invariant



# Symmetry for $lpha_0=rac{\pi}{2}$



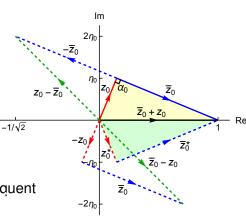
- In HS texture, simple sign change of z<sub>0</sub> (or of z̄<sub>0</sub>, but not both), flips orientation of the UT (see fig →)
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- But iff  $\alpha = \pm \frac{\pi}{2}$



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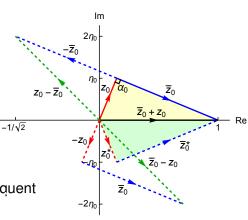
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- Equivalent to CP transformation
- Can be reversed by a subsequent actual CP transformation



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- But iff  $\alpha = \pm \frac{\pi}{2}$
- Equivalent to CP transformation
- Can be reversed by a subsequent actual CP transformation
- Symmetry is good to all orders



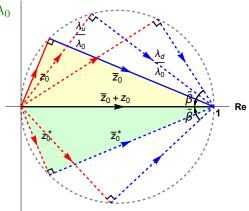
# Symmetry for $\beta_0 = \frac{\pi}{8}$



• First consider  $\beta_0 = \tilde{\beta} \neq \frac{\pi}{8}$  (fig $\rightarrow$ ) Im keeping  $\alpha = \frac{\pi}{2}$  and  $\lambda_u + \lambda_d = \lambda_0$ 

Clearly now

$$\left|\frac{\lambda_u}{\lambda_d}\right|=\tan\tilde{\beta},$$
 and  $-\frac{\pi}{2}<\tilde{\beta}<\frac{\pi}{2}$ 



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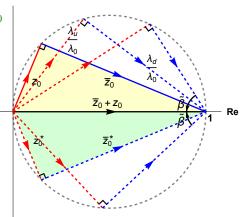
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and 
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 (\*)

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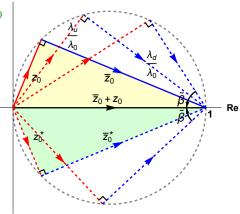
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 (\*)

- Iff  $\tilde{\beta} = \frac{\pi}{8}$ , the result is just a CP transformation
- $\Rightarrow$  to fix  $\beta_0 = \frac{\pi}{8}$  require symmetry under transformation (\*) followed by CP flip





- Proposed geometric-hierarchical MM texture
- Mass hierarchy "slopes" are related to UT sides
- ullet Symmetries constrain forms  $ightarrow lpha \simeq rac{\pi}{2}$  and  $eta \simeq rac{\pi}{8}$



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c.f.  $0.177 \pm 0.002$  (exp)

Precise predictions of UT angles:

$$\alpha - \frac{\pi}{2} = (1.30 \pm 0.02)^{\circ} \text{ c.f. } (1.6 \pm 1.4)^{\circ} \text{ (exp)}$$

$$\beta - \frac{\pi}{8} = (-0.2 \pm 0.1)^{\circ} \text{ c.f. } (0.1 \pm 0.4)^{\circ} \text{ (exp)}$$

$$\gamma - \frac{3\pi}{8} = (-1.1 \pm 0.1)^{\circ}$$
 c.f.  $(-1.8 \pm 1.3)^{\circ}$  (exp)



# **Backup Slides**



## Isospin Reflection Symmetry?



Can re-write texture:

$$M_q^{HS} \equiv n_q \begin{pmatrix} c' \lambda_q^4 & b' \lambda_q^3 & 0 \\ b' \lambda_q^{*3} & b' \lambda_q^2 & A_0 \lambda_q^2 \\ 0 & A_0 \lambda_q^{*2} & 1 \end{pmatrix} \pm d \lambda_q^4 I$$

 $(b' \simeq b)$ . Still get good fit to data.

- First (leading) matrix solely responsible for quark mass differences and mixing parameters.
- Second (small) matrix is  $I_z$ -dependent "pedestal" on quark masses. Symmetric under a generation-SU(3) symmetry.
- All coefficients ( $\lambda_0$ ,  $A_0$ , b', c', d) symmetric under isospin reflection operator  $u \leftrightarrow d$ .
- Symmetry broken (only) by  $\lambda_q$ ,  $n_q$  and the sign of d.



# Analytic NLO Solutions: 1) Mixing Parameters VARWICK

• We give here the algebraic NLO solutions of the texture:

$$\begin{split} \lambda &= \lambda_0 \left( 1 + f_{\lambda} \lambda_0^2 \right) + \mathcal{O}(\lambda_0^5) \\ A &= A_0 \left\{ 1 + \left[ \frac{1}{4} (3b - 2\rho_0) - 2f_{\lambda} \right] \lambda_0^2 \right\} + \mathcal{O}(\lambda_0^4) \\ \overline{\rho} &= \rho_0 \left( 1 + c_0 f_{\rho} \lambda_0^2 \right) + \mathcal{O}(\lambda_0^4) \\ \overline{\eta} &= \eta_0 \left\{ 1 + \left[ s_0 f_{\rho} + \frac{1}{2} (1 - 5b) \right] \lambda_0^2 \right\} + \mathcal{O}(\lambda_0^4), \end{split}$$
 where  $f_{\lambda} = \frac{3}{4} f_A - \frac{5}{4} + \eta_0 \delta_c,$ 

$$f_{A} = \frac{1}{2} \left[ A_0^2 + \frac{1}{4} (c_A + c_A) \right] \qquad \delta_A = \frac{1}{2} (c_A - c_A)$$

where 
$$f_{\lambda} = \frac{1}{4}f_A - \frac{1}{4} + \eta_0 \sigma_c$$
,  $f_A = \frac{1}{b} \left[ A_0^2 + \frac{1}{2} (c_d + c_u) \right], \quad \delta_c = \frac{1}{b} (c_d - c_u)$  and  $f_{\rho} = \frac{1}{s_0} \left[ -\frac{1}{2} f_A + \frac{7}{4} b - \frac{1}{2} \delta_c \right] + s_0 (1 + \delta_c)$ .

• NLO corrections above, as fractions of LO terms are respectively:  $-5.8 \times 10^{-3}, +2.6\%, +3.6\%$  and -1.8% (using fitted param values from table).

# Analytic NLO Solutions: 2) Mass Ratios



For the quark mass ratios, we find:

$$\begin{split} \frac{m_1^q}{m_2^q} &= -\lambda_q^2 (1 - r_q) \left\{ 1 + \left[ r_A \frac{(2 - r_q)}{(1 - r_q)} - 2 \right] \lambda_q^2 \right\} + \mathcal{O}(\lambda_q^6) \\ \frac{m_2^q}{m_3^q} &= b \lambda_q^2 \left[ 1 + (1 - r_A) \lambda_q^2 \right] + \mathcal{O}(\lambda_q^6), \end{split}$$

where  $r_q = \frac{c_q}{b}$  and  $r_A = \frac{A_0^2}{b}$ .

- NLO corrections to mass ratios  $m_c/m_t$ ,  $m_s/m_b$ ,  $m_u/m_c$ ,  $m_d/m_s$  as fractions of LO terms are (resp.)  $-4.3 \times 10^{-3}$ , -2.5%, +4.3%, and +4.5% (using fitted param values from table).
- All results compatible with full numerical results reported in table.