

Renormalisation scheme dependence of the trilinear Higgs coupling in extended scalar sectors

Herbstschule of High-Energy Physics, Bad Honnef 2025

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Motivation: Trilinear Coupling

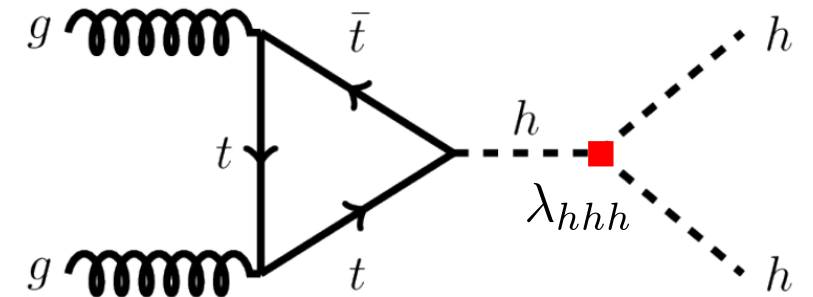
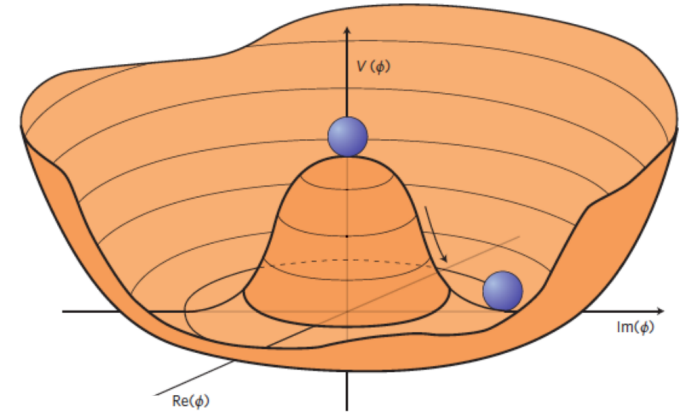
[Bahl, Braathen, Gabelmann, Weiglein 2305.03015]

- Trilinear Higgs coupling determines the shape of the Higgs potential:

- crucial for understanding dynamics of electroweak phase transitions

$$V_{\text{SM}} \supset \frac{m_h^2}{2} h^2 + \underbrace{\frac{3m_h^2}{v}}_{\lambda_{hhh}^{\text{SM},0}/6} h^3 + \frac{3m_h^2}{v^2} h^4$$

- sensitive to BSM physics
 - exp. limit: $-1.2 < \frac{\lambda_{hhh}}{\lambda_{hhh}^{\text{SM},0}} < 7.2$ at 95% C.L.
[ATLAS, 2406.09971]
[CMS, 2407.13554]



Singlet Extension of the Standard Model (SSM)

[Bahl, Braathen, Gabelmann, Weiglein 2305.03015]

- we consider the real singlet extension of the SM

$$V(\Phi, S) = \mu^2 |\Phi| + \frac{\lambda}{2} |\Phi|^2 + \kappa_{SH} (\Phi^\dagger \Phi) S + \frac{\lambda_{SH}}{2} (\Phi^\dagger \Phi) S^2 + \frac{m_S^2}{2} S^2 + \frac{\kappa_S}{3} S^3 + \frac{\lambda_S}{2} S^4$$

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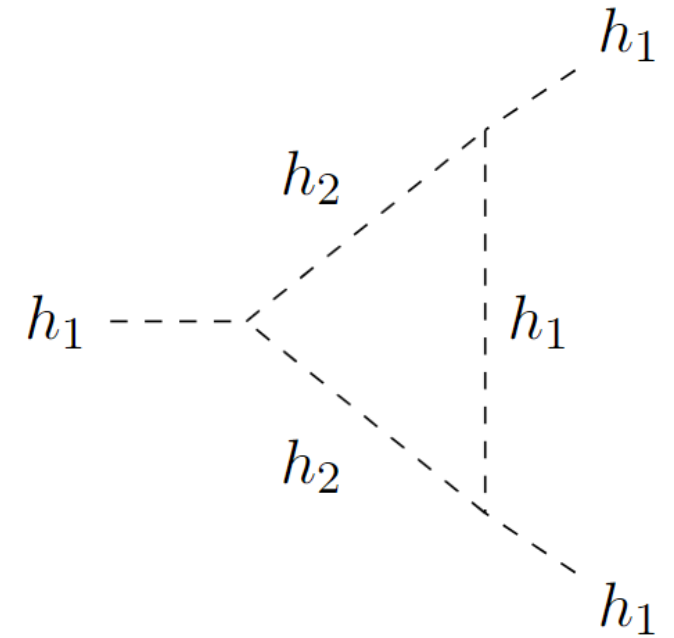
$$V(\Phi, S) = \mu^2 |\Phi|^2 + \frac{\lambda}{2} |\Phi|^4 + \kappa_{SH} (\Phi^\dagger \Phi) S + \frac{\lambda_{SH}}{2} (\Phi^\dagger \Phi) S^2 + \frac{m_S^2}{2} S^2 + \frac{\kappa_S}{3} S^3 + \frac{\lambda_S}{2} S^4$$

- after symmetry breaking: $\Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{2} G^+ \\ v+h+iG \end{pmatrix}$ $S = s + v_S$
- the fields mix to 2 CP-even eigenstates $h_{1,2}$
- the Lagrangian parameters are traded for $m_{h_1}, m_{h_2}, \alpha, v, v_S, \kappa_S, \kappa_{SH}$
- trilinear coupling at tree level: $\lambda_{h_1 h_1 h_1} = 3m_{h_1}^2 \left(\frac{\cos^3 \alpha}{v} + \frac{\sin^3 \alpha}{v_S} \right)$

5 new independent parameters

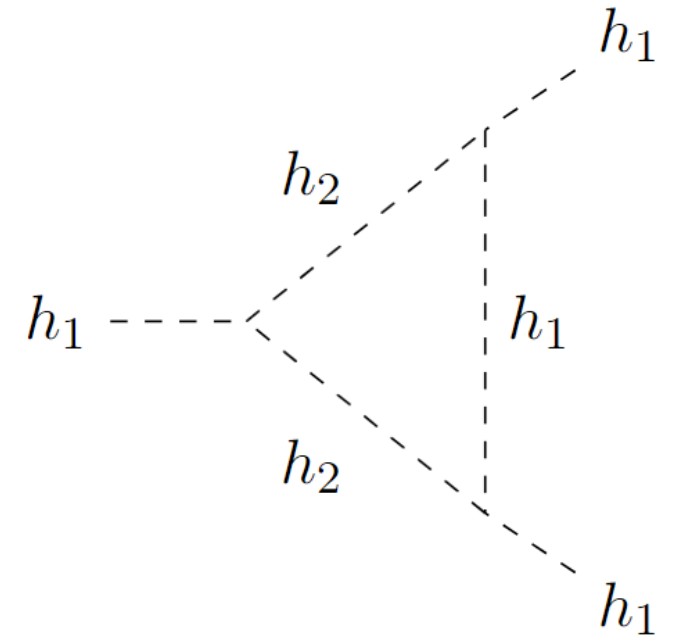
Motivation: Precision Calculations

- reliable and numerically stable theoretical predictions of physical observables are imperative
 - important for comparison with experimental data
 - needed for consistent constraints on parameter space of BSM models



Problem: Precision Calculations

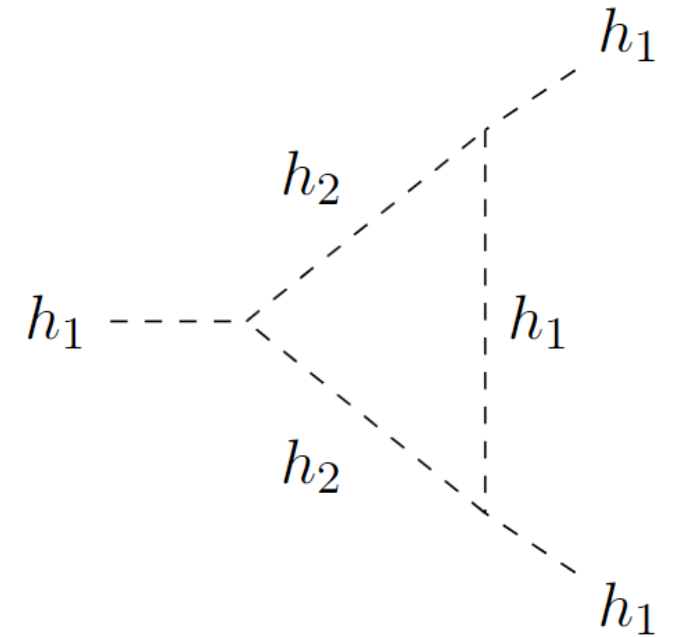
- integrals over loop momenta diverge
extracted by dimensional regularisation 
 $D = 4 - 2\epsilon$



Problem: Precision Calculations

- integrals over loop momenta diverge
extracted by dimensional regularisation ✓
 $D = 4 - 2\epsilon$
- final expression contains divergence for $D \rightarrow 4$
absorbed in renormalised parameters ✓

understanding of renormalisation scheme
dependence needed



Choice of Renormalization scheme

On shell renormalisation

- renormalised parameters correspond to measurable quantities
- impose a renormalisation condition $\rightarrow$ this subtracts the divergence

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Minimal subtraction renormalisation

- no conditions $\rightarrow$ only divergent part is subtracted

$$\delta m^2 = a \frac{1}{\epsilon} + \cancel{b\epsilon^0} + c\epsilon$$

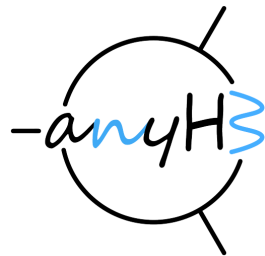
- μ dependence $\rightarrow$ running of parameters

$$m_B = m^{\overline{\text{MS}}}(\mu) + \delta m^{\overline{\text{MS}}}(\mu) = m^{\text{OS}} + \delta m^{\text{OS}}$$

Predicting the Trilinear Coupling: anyH3

[Bahl, Braathen, Gabelmann, Weiglein; 2305.03015]

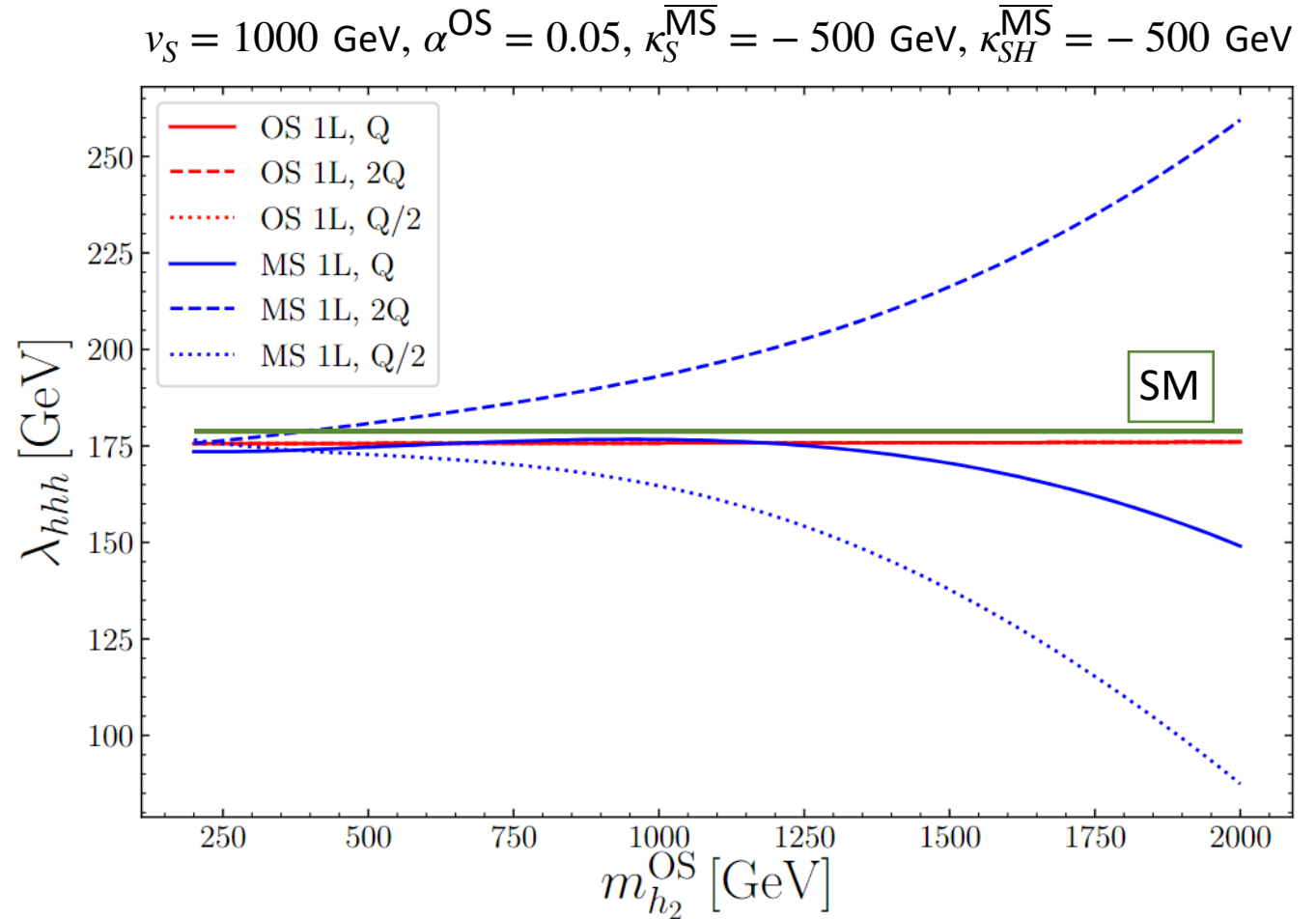
- Python code anyH3 allows fast computation of λ_{hhh} at full one-loop level
- customize:
 - parameters/couplings
 - renormalisation
 - tadpole treatment



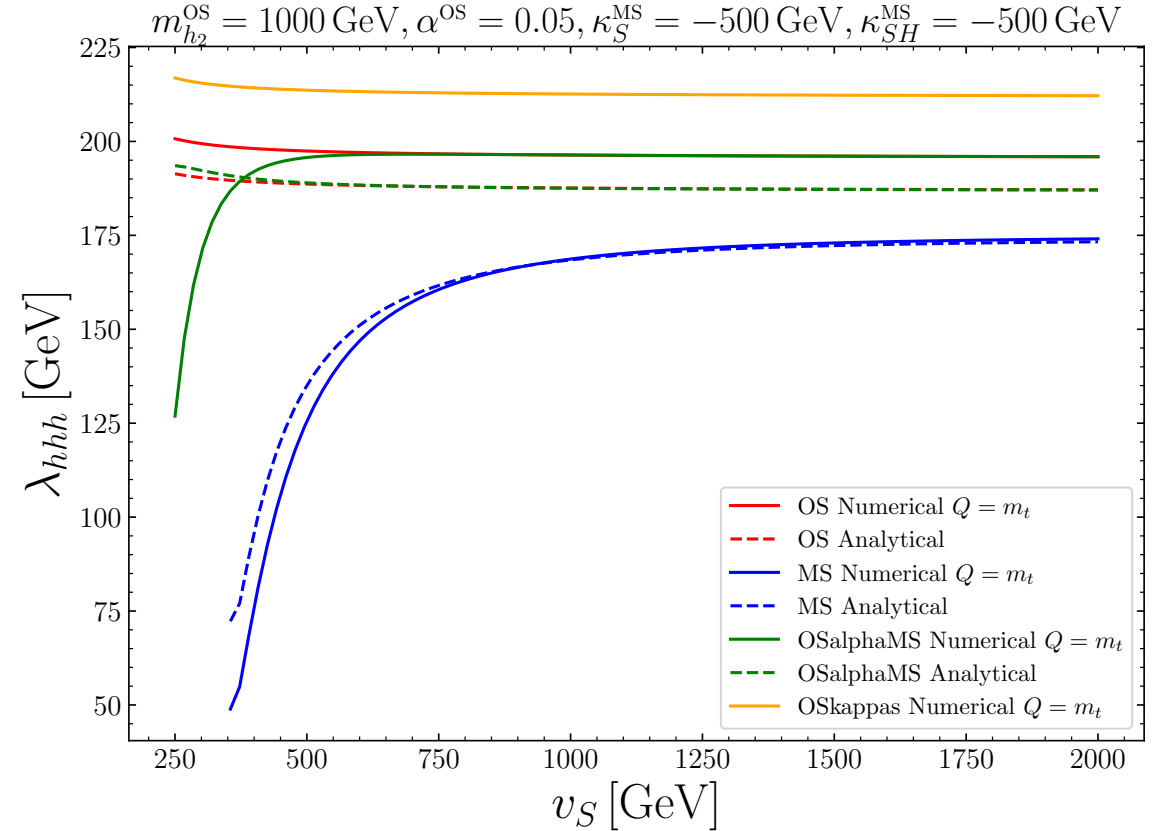
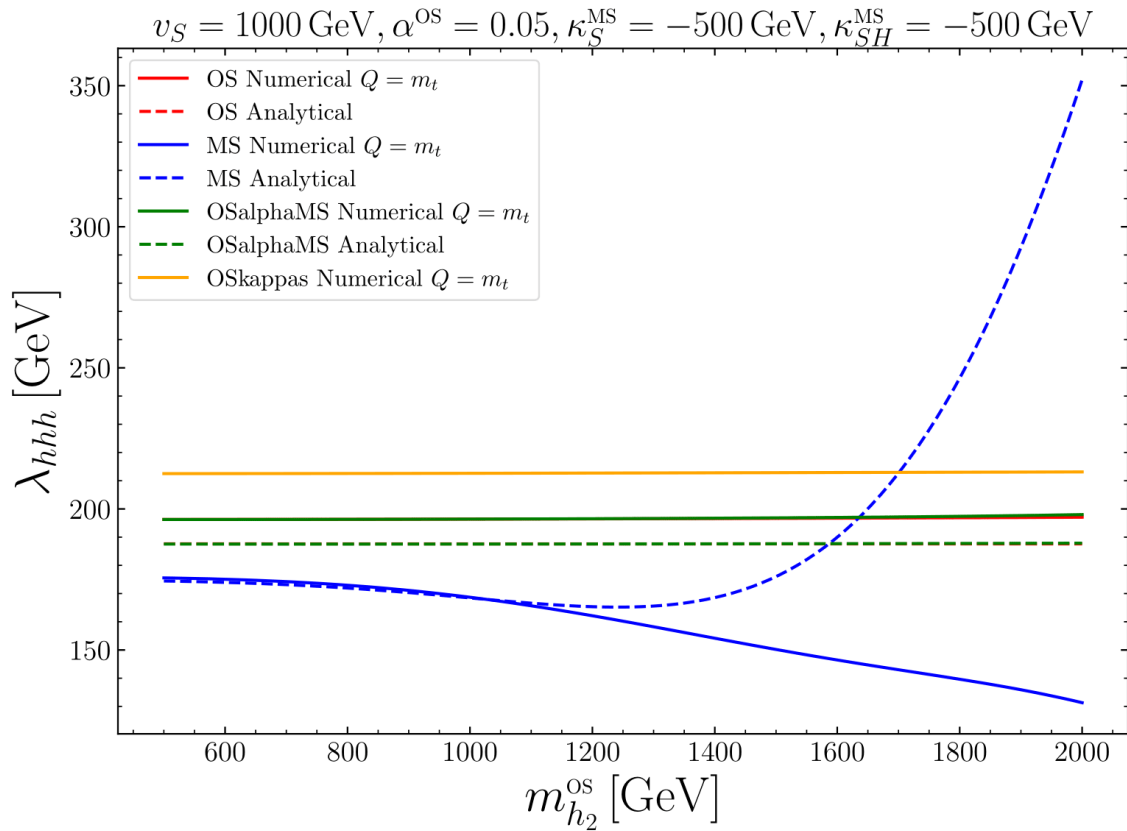
$$\begin{aligned}
 \lambda_{hhh} = & \underbrace{\text{tree-level: } \lambda_{hhh}^{(0)}}_{\text{tree-level: } \lambda_{hhh}^{(0)}} + \underbrace{\text{one-particle irreducible: } \delta_{\text{genuine}}^{(1)} \lambda_{hhh}}_{\text{one-particle irreducible: } \delta_{\text{genuine}}^{(1)} \lambda_{hhh}} \\
 & + \underbrace{\text{external leg corrections: } \delta_{\text{WFR}}^{(1)} \lambda_{hhh}}_{\text{external leg corrections: } \delta_{\text{WFR}}^{(1)} \lambda_{hhh}} + \underbrace{\text{tadpoles: } \delta_{\text{tad. WFR}}^{(1)} \lambda_{hhh} + \delta_{\text{tadpoles}}^{(1)} \lambda_{hhh}}_{\text{tadpoles: } \delta_{\text{tad. WFR}}^{(1)} \lambda_{hhh} + \delta_{\text{tadpoles}}^{(1)} \lambda_{hhh}} \\
 & + \underbrace{\text{renormalisation: } \delta_{\text{CT}}^{(1)} \lambda_{hhh}}_{\text{renormalisation: } \delta_{\text{CT}}^{(1)} \lambda_{hhh}}
 \end{aligned}$$

Scheme Dependence of the Trilinear Coupling

- implemented conversion
- varying Q provides an estimate of size of higher order effects
- MS band contains OS values



Comparison of different schemes



Summary

- trilinear Higgs coupling as an important observable
- first tests of renormalisation scheme dependence in SSM
- further steps: investigate full parameter space, criteria definition for the schemes, extension to other BSM models

Goal: find algorithmic criteria for which scheme to use depending on region in parameter space $\longrightarrow$ deliver a robust, consistent prediction for λ_{hhh}

Backup Slides

Tadpole Treatment

- Tadpole-free $\overline{\text{MS}}$ scheme:

$$T_h = t_h + \delta^{(1)} t_h = 0 \Rightarrow t_h = -\delta^{(1)} t_h = - \left. \frac{\partial V^{(1)}}{\partial h} \right|_{\text{min}}$$

- OS tadpole renormalisation:

$$t_h = 0 \quad \delta_{\text{CT}}^{(1)} t_h = -\delta^{(1)} t_h$$

- $\overline{\text{MS}}$ tadpole renormalisation at the tree-level minimum (FJ):

$$t_h = 0 \text{ and } \overline{\text{MS}} \text{ renormalised tadpoles}$$

Renormalisation Schemes for SSM

- full OS scheme:

$$\delta\alpha = \frac{1}{2} \frac{\Sigma_{h_1 h_2}(p^2=m_{h_1}^2) + \Sigma_{h_1 h_2}(p^2=m_{h_2}^2)}{m_{h_2}^2 - m_{h_1}^2} \quad \delta v_S = \cos\alpha \frac{\delta^{(1)} t_{h_2}}{m_{h_2}^2} + \sin\alpha \frac{\delta^{(1)} t_{h_1}}{m_{h_1}^2}$$

- MS scheme: all parameters renormalised $\overline{\text{MS}}$
- hybrid scheme: OS masses and VEV but $\overline{\text{MS}}$ mixing angle

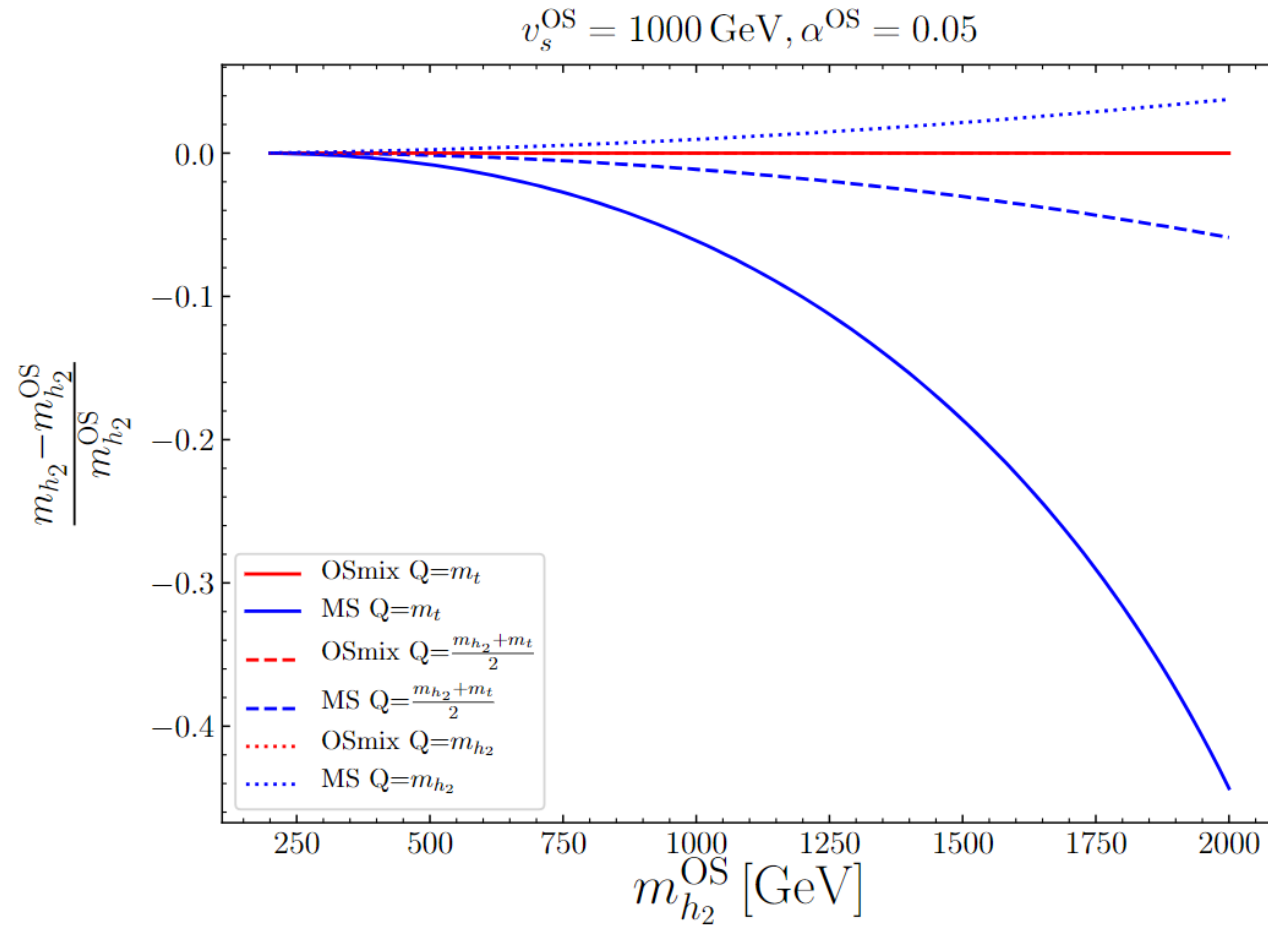
$$\delta v_S = \cos\alpha \frac{\delta^{(1)} t_{h_2}}{m_{h_2}^2} + \sin\alpha \frac{\delta^{(1)} t_{h_1}}{m_{h_1}^2}$$

Analytic Formulas

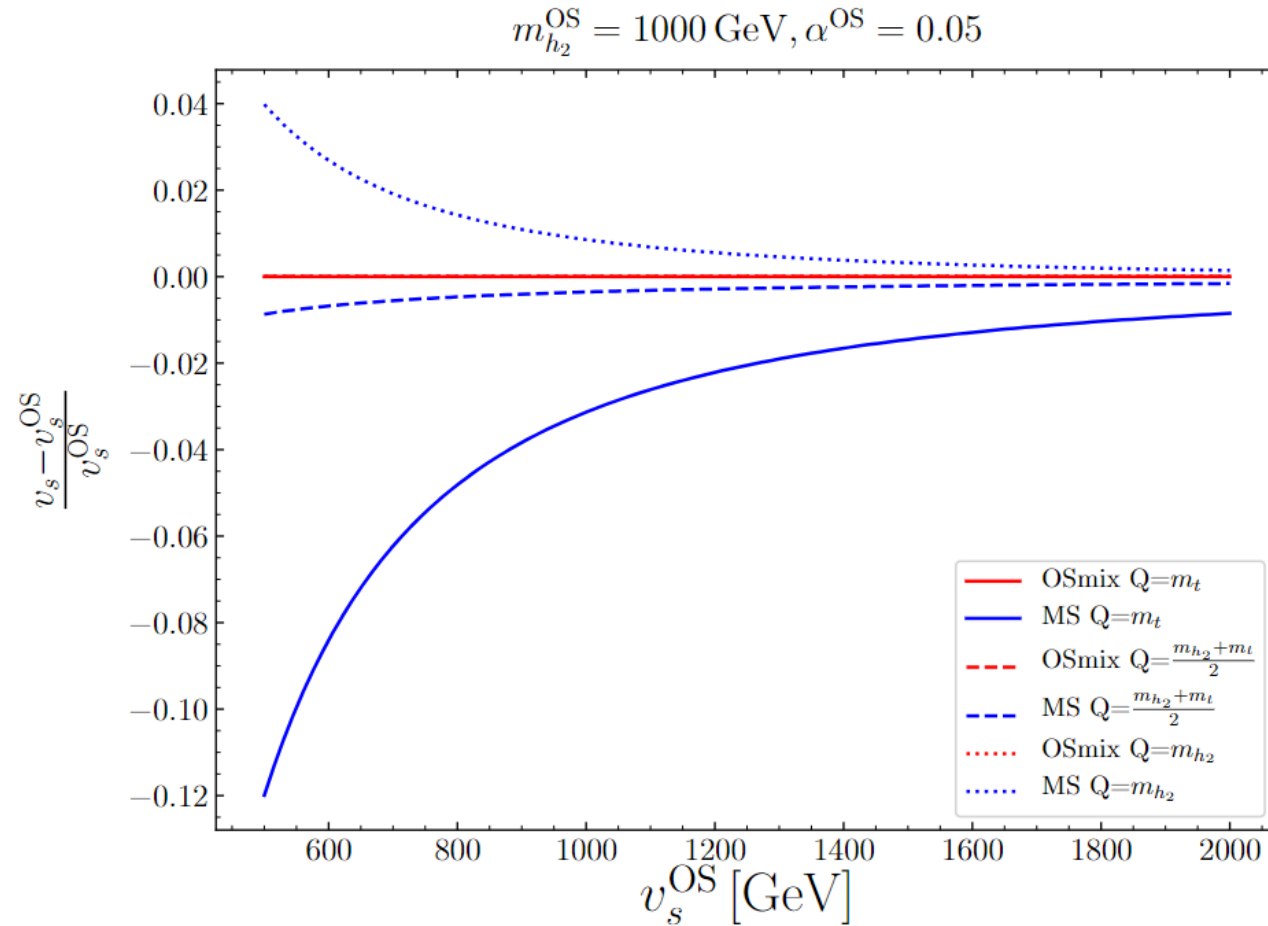
$$\lambda_{hhh}^{OS,(1)} = \frac{\kappa_{SH}^2 v^2}{64 m_{h_2}^2 \pi^2 v_S^4} \left[6\alpha m_{h_2}^2 (-4 + \sqrt{3}\pi) v_S - 2\kappa_{SH} v_S v \right. \\ \left. - \alpha (-3 + \sqrt{3}\pi) (2\kappa_S v_S^2 - 3\kappa_{SH} v^2) \right]$$

$$\lambda_{hhh}^{MS,(1)} = \frac{1}{64 m_{h_2}^2 \pi^2 v_S^4 v^2} \left\{ 6\alpha m_{h_2}^6 v_S^3 + 6\kappa_{SH} m_{h_2}^4 v_S^3 v \right. \\ - 6\alpha \kappa_{SH}^2 v_S (-2m_{h_2}^2 + \kappa_S v_S) v^4 - 2\kappa_{SH}^3 v_S v^5 + 9\alpha \kappa_{SH}^3 v^6 \\ - 6 \left[\kappa_{SH} m_{h_2}^2 v_S^2 v (m_{h_2}^2 v_S - \kappa_{SH} v^2) \right. \\ + \alpha (m_{h_2}^6 v_S^3 + 2\kappa_{SH} m_{h_2}^2 v_S^2 (2m_{h_2}^2 - \kappa_S v_S) v^2 \\ \left. + \kappa_{SH}^2 v_S (9m_{h_2}^2 - 2\kappa_S v_S) v^4 + 3\kappa_{SH}^3 v^6) \right] \cdot \ln \left(\frac{m_{h_2}^2}{\mu^2} \right) \left. \right\}$$

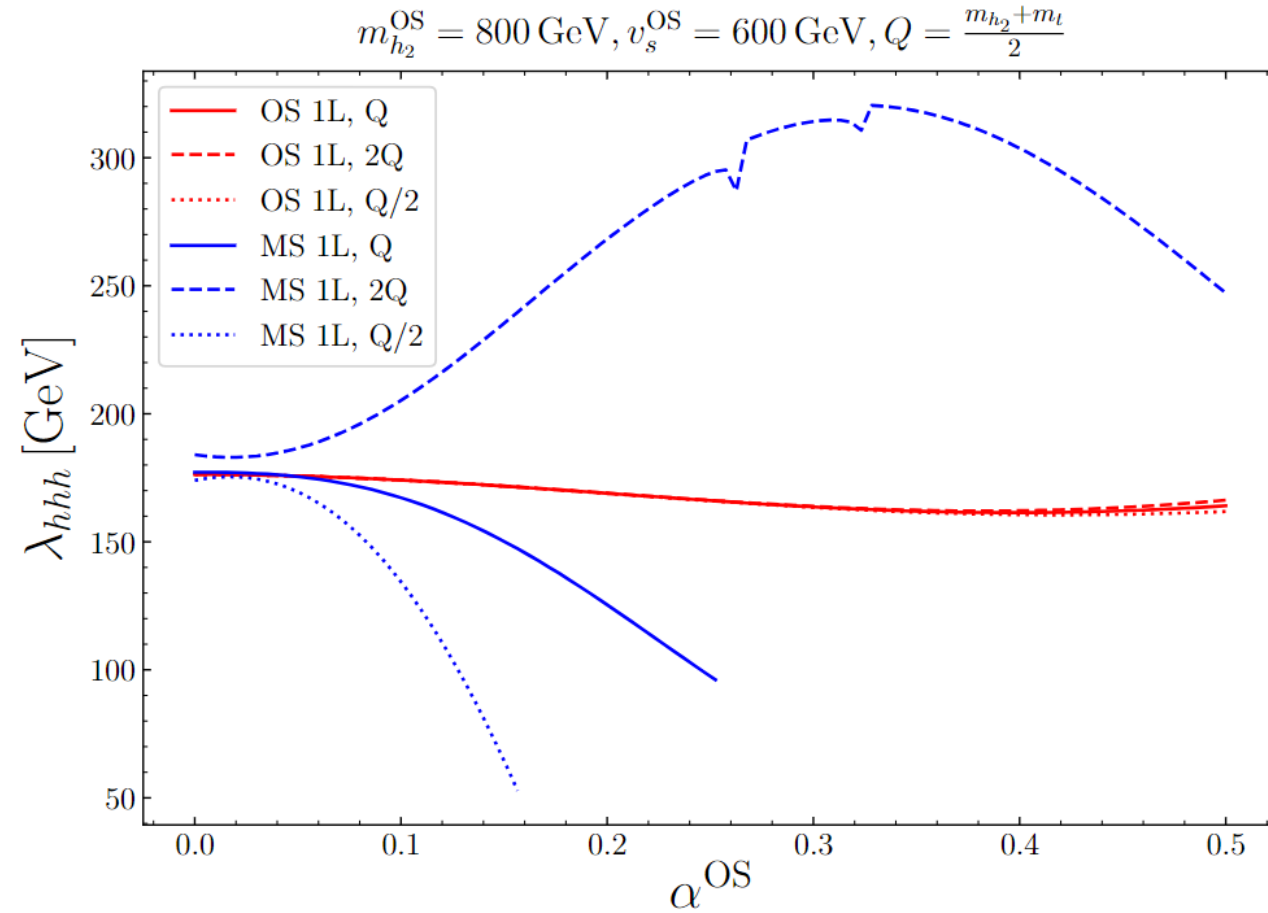
Change of BSM mass



Change of singlet VEV



Dependence on the mixing angle



Reparametrisation of Lagrangian

$$\lambda = \frac{c_\alpha^2 m_{h_1}^2 + s_\alpha^2 m_{h_2}^2}{v^2} - \frac{t_\phi}{v^3}$$

$$\lambda_{SH} = \frac{(m_{h_1}^2 - m_{h_2}^2) c_\alpha s_\alpha}{vv_S} - \frac{\kappa_{SH}}{v_S}$$

$$\lambda_S = \frac{(m_{h_1}^2 + m_{h_2}^2) v_S + (m_{h_2}^2 - m_{h_1}^2) v_S c_{2\alpha} - 2t_S - 2\kappa_S v_S^2 + \kappa_{SH} v^2}{8v_S^3}$$

$$M_S^2 = \frac{6t_S - \kappa_{SH} v^2 - 2\kappa_S v_S^2 - (m_{h_1}^2 + m_{h_2}^2) v_S + (m_{h_2}^2 - m_{h_1}^2) (v s_{2\alpha} - v_S c_{2\alpha})}{4v_S}$$

$$\mu^2 = \frac{3}{2} \frac{t_\phi}{v} - \frac{1}{2} \kappa_{SH} v_S - \frac{1}{2} (m_{h_1}^2 c_\alpha^2 + m_{h_2}^2 s_\alpha^2) - \frac{v_S}{2v} (m_{h_1}^2 - m_{h_2}^2) c_\alpha s_\alpha.$$