# High-precision measurement of the $B_d^0$ lifetime and other lifetime measurements in ATLAS

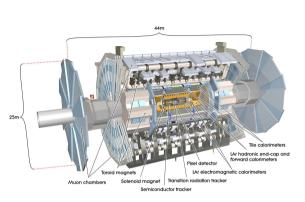
Maria Smizanska, Lancaster University On behalf of the ATLAS collaboration

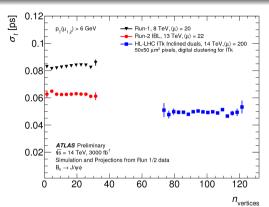
More than Lifetime, Siegen, 22-26 Sept 2025

#### Outline

- ATLAS detector and B-physics data
- Precision measurement of the  $B_d^0$  meson lifetime performed at the ATLAS using Run2 data Eur. Phys. J. C (2025) 85:736.  $\square$ .
- Some other lifetime measurements in ATLAS
  - Measurement of the relative width difference  $\Delta \Gamma_d / \Gamma_d$  of the  $B_d^0 \overline{B_d^0}$  system with the ATLAS detector JHEP06(2016)081
  - Measurement of effective lifetime in  $B^0_{(s)} \to \mu^+ \mu^-$ JHEP 09 (2023) 199  $\square$
  - Measurement of the CP-violating phase, the width difference  $\Delta\Gamma_s$  between the meson mass eigenstates and the average decay width  $\Gamma_s$  in the  $B_s \to J/\psi \phi$  decay. Eur. Phys. J. C 81 (2021) 342  $\Box$
- HL-LHC and B-physics performance in ATLAS

#### ATLAS detector and features important for high precision B-physics measurements



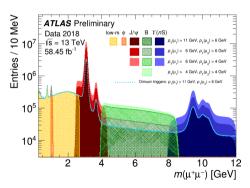


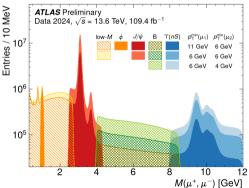
Time resolution of  $B_S^0 \to J/\psi \phi$  for different numbers of reconstructed PV in the same bunch crossing ATL-PHYS-PUB-2013-010  $\overline{C}^0$ 

- ullet Inner Detector: PIX, SCT and TRT,  ${m 
  ho_{
  m T}} > 0.5\,{
  m GeV},\,|\eta| < 2.5$ 
  - Run2: new IBL 25% improvement of time resolution with respect to Run1.
  - Time, mass resolutions remain stable within increasing pileup in Run 2
- Muon Spectrometer: triggering ( $|\eta|$ < 2.4), precision tracking ( $|\eta|$ < 2.7)

#### Data-taking conditions for di-muon final states

- ATLAS has collected 139 fb<sup>-1</sup> of data in Run 2, and 25 fb<sup>-1</sup> in Run 1
- Typical triggers di-muons with p<sub>T</sub> thresholds of either 4 GeV or 6 GeV or 11 GeV (vary over run periods)
  - Low  $p_T$  thresholds were activated in the end of fills when the instantaneous luminosity decreases.
  - Additional trigger selections are applied, e.g. on di-muon masses, targeting different analyses, as shown in Fig.
- For Run2 events accepted for analysis, an average number of pp interactions per bunch crossing (pile-up) was 31.
- Trigger tracking: transverse impact parameter d0 < 10mm on all tracks, no trigger cuts from lower d0 side.





#### Precision measurement of the $B_d^0$ meson lifetime at the ATLAS

ATLAS

#### A new record for precision on B-meson lifetimes

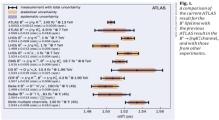
As direct searches for physics beyond the Standard Model continue to push frontiers at the LHC, the b-hadron physics sector remains a crucial source of insight for testing established theoretical models

The ATLAS collaboration recently published a new measurement of the Bo lifetime using B°→1/wK° decays from the entire Run-2 dataset it has recorded at 13 TeV. The result improves the precision of previous world-leading measurements by the CMS and LHCb collaborations by a factor of two. Studies of b-hadron lifetimes probe

our understanding of the weak interaction. The lifetimes of b-hadrons can be systematically computed within the heavy-quark expansion (HOE) framework, where b-hadron observables are expressed as a perturbative expansion in inverse powers of the b-quark mass.

ATLAS measures the "effective" Belifetime, which represents the average decay time incorporating effects from mixing and CP contributions, as r(Bo) = 1.5053 ± 0.0012 (stat.) ± 0.0035 (syst.) ps. The result is consistent with previous measurements published by ATLAS and other experiments, as summarised in figure 1. It also aligns with theoretical predictions from HOE and lattice OCD, as well as with the experimental world average

The analysis benefitted from the large Run-2 dataset and a refined trigger selection, enabling the collection of an extensive sample of 2.5 million B° → IAvK\* decays. Events with a I/w meson decaying into two muons with sufficient transverse momentum are cleanly identified in the ATLAS Muon Spectrometer by the first-level hardware trigger. In the nextlevel software trigger, exploiting the full detector information, these muons are



then combined with two tracks meas- tematic uncertainties, with the largest ured by the Inner Detector, ensuring they contributions arising from the correoriginate from the same vertex

through a two-dimensional unbinned distribution. maximum-likelihood fit, utilising the measured B°-candidate mass and decay with the average decay width (Γ.) of time, and accounting for both signal and the light and heavy B .- meson mass background components. The limited eigenstates, also measured by ATLAS. hadronic particle-identification capa- to determine the ratio of decay widths bility of ATLAS requires careful model- as E./E. = 0.9905 + 0.0022 (stat.) + 0.0036 ling of the significant backgrounds from (syst.) + 0.0057 (ext.). The result is conother processes that produce I/w mesons. sistent with unity and provides a strin-The sensitivity of the fit is increased by gent test of OCD predictions, which also estimating the uncertainty of the decay-support a value near unity. time measurement provided by the ATLAS tracking and vertexing algorithms on a Further reading per-candidate basis. The resulting life- ATLAS Collab. 2024 arXiv:2411.09962.

lation between Bo mass and lifetime, The B\*-meson lifetime is determined and ambiguities in modelling the mass

ATLAS combined its measurement

time measurement is limited by sys- ATLAS Collab. 2021 Eur. Phys. J. C 81 342.

 "Precision measurement of the B<sub>d</sub><sup>0</sup> meson lifetime performed at the ATLAS using Run2 data " Eur. Phys. J. C (2025) 85:736.

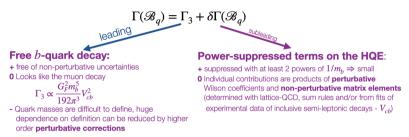
- ATLAS B<sup>0</sup> lifetime is the most precise measurement to date
- Also measured the decay rate  $\Gamma_d$  of  $B^0$
- $\Gamma_d/\Gamma_s$  determined using ATLAS  $\Gamma_s$ measured in  $B_s^0 \to J/\psi \phi$  Eur. Phys. J. C 81 (2021) 342 🕜
- Recognised by ATLAS ATLAS Briefing 25 November 2024 7
- Recognised by CERN LHC CERN Seminar and Article in CERN Courier p. 15). 🖸

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#### Physics motivation

• In the Heavy Quark Expansion (HQE) theory the decay rate  $\Gamma$  = 1/  $\tau$  is calculated:



- $\Gamma$  suffers from huge uncertainties of the leading term  $\Gamma_3$ ;  $\Gamma_d = 0.63^{+0.11}_{-0.07}~\mathrm{ps}^{-1}$  Lenz, et al. 2023  $\Box$
- In rates ratios the leading term  $\Gamma_3$  exactly cancels leaving just small uncertainties from sub-leading terms;  $\Gamma_d/\Gamma_s = 1.003 \pm 0.006$  Lenz, Piscopo, Rusov 2023
- Lifetimes measurements also test models of New Physics (BSM) contributing to both leading and sub-leading terms. Excluding BSM effects larger than 5%, will already constrain many BSM scenarios Lenz2021
- Since the theory predictions are more precise for ratios we also extracted  $\Gamma_d/\Gamma_s$ .

Derivation of decay rate  $\Gamma_d$  and ratio  $\Gamma_d/\Gamma_s$  from measured  $B_d^0$  lifetime.

#### $\Gamma_d$

The lifetime we measured in  $B^0 \to J/\psi K^{*0}$  is the effective lifetime  $\tau_{B_d^0}$  related to decay widths:  $\Gamma_L$ ,  $\Gamma_H$  of the light and heavy mass eigenstates of  $B_d^0 - \overline{B_d^0}$  system via: Dunietz, Fleischer, Nierste, 2000  $\square$ 

$$\tau_{B^0} = \frac{1}{\Gamma_d} \frac{1 + y^2}{1 - y^2},\tag{1}$$

where  $\Gamma_d = (\Gamma_L + \Gamma_H)/2$ ;  $y = \Delta \Gamma_d/(2\Gamma_d) = (\Gamma_L - \Gamma_H)/(2\Gamma_d)$ .

Using the value of y from Heavy Flavour Averaging group (HFLAV) Ref.HFLAV:2023  $\mathcal{C}$ , our  $\tau_{\mathcal{B}_d^0}$  value and Eq(1) allows  $\Gamma_d$  to be extracted.

#### $\Gamma_d/\Gamma_s$

ATLAS measured  $\Gamma_s = 0.6703 \pm 0.0014 (\text{stat.}) \pm 0.0018 (\text{syst.}) \text{ ps}^{-1} \text{ from } B_s^0 \to J/\psi \phi \text{ Eur. Phys. J. C 81}$  (2021) 342  $\square$ . This result combined with  $B_d^0 \to J/\psi K^*$  allowed us to determine the ratio  $\Gamma_d/\Gamma_s$ .

#### Extracting $B_d^0$ lifetime: Maximum-likelihood fit

2D unbinned maximum-likelihood fit applied simultaneously to mass and proper decay time of  $B_d^0$  candidates. Likelihood Function formula:

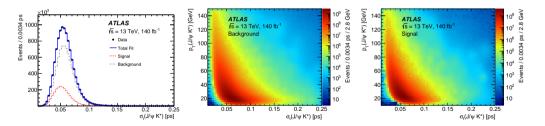
$$\ln L = \sum_{i=1}^{N} w(t_i) \ln[f_{\text{sig}} \mathcal{M}_{\text{sig}}(m_i) \mathcal{T}_{\text{sig}}(t_i, \sigma_{t_i}, p_{\text{T}_i}) + (1 - f_{\text{sig}}) \mathcal{M}_{\text{bkg}}(m_i) \mathcal{T}_{\text{bkg}}(t_i, \sigma_{t_i}, p_{\text{T}_i})]$$

- $f_{\rm sig}$  fraction of signal events in the total number of events, N.
- ullet  $\mathcal{M}_{sig}$ ,  $\mathcal{T}_{sig}$  mass and time signal PDFs
- ullet  $\mathcal{M}_{bkg}$ ,  $\mathcal{T}_{bkg}$  mass and time Backround PDFs
- The mass  $m_i$ , the proper decay time  $t_i$ , its uncertainty  $\sigma_i$  and the  $B_d^0$  candidate transverse momentum  $p_{T_i}$  are the values measured from the data for each event i.
- Weight w<sub>i</sub> accounts for event selection efficiency, based on MC.

### Conditional probability of the time uncertainty $\sigma_{t_i}$ and $\rho_{T_i}$ of $B_d^0$

- Per-candidate time errors  $\sigma_{t_i}$ , extracted from data in the vertex fit of each  $B_d^0$  are used for deconvolution of proper decay times  $t_i$ .
- $\sigma_{t_i}$  are different for signal and background, see fig Left. Same it true for  $\rho_T$  of  $B_d^0$ .
- To account for these differences, the 2D probability terms  $C_j(\sigma_{t_i}, p_{\Gamma_i})$  are introduced into Signal and Bg Time PDF terms  $\mathcal{T}$ . Method first used in G.Punzi

$$\mathcal{T}_j(t_i,\sigma_{t_i},p_{\Gamma_i}) = P_j(t_i,\sigma_{t_i},p_{\Gamma_i}) \cdot C_j(\sigma_{t_i},p_{\Gamma_i}), \tag{2}$$



Left:  $\sigma_{t_i}$  for Signal and Bg. Middle and Right: 2D probability terms  $C_j(\sigma_{t_i}, p_{\Gamma_i})$  for Bg. and Signal. The sPlot technique Pivk:2004ty  $rac{1}{2}$  is used to separate Signal and Bg.

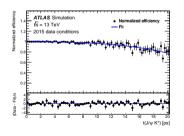
#### Efficiencies and corrections

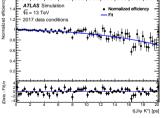
- Trigger, offline reconstruction, event selections bias the reconstructed proper-decay time distribution.
- Trigger and offline tracking impose  $|d_0| < 10$  mm, for all four final-state tracks of  $B^0 \to J/\psi K^{*0}$ , resulting in inefficiency at large times.
- Inefficiencies determined by signal MC, passed through simulation of detector response & triggers, offline tracking and event selections - as data.

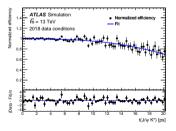
Ratio of proper decay time distributions before and after the whole chain, were fitted by:

$$1/w(t_i) = p_0 \cdot [1 - p_1 \cdot (\operatorname{Erf}((t_i - p_3)/p_2) + 1)]$$
 (3)

Erf is error function,  $p_0$ ,  $p_1$ ,  $p_2$ ,  $p_3$  determined in the fit. Weights  $w(t_i)$  are used to re-weight each event in Likelihood fit.







#### Systematic uncertainties

Source of uncertainty	Systematic uncertainty [ps]
ID alignment	0.00108
Choice of mass window	0.00104
Time efficiency	0.00130
Best-candidate selection	0.00041
Mass fit model	0.00152
Mass-time correlation	0.00229
Proper decay time fit model	0.00010
Conditional probability model	0.00070
Fit model test with pseudo-experiments	0.00002
Total	0.0035

ullet The statistical error of this  $B_d^0$  lifetime measurement is 0.0012 (stat.) ps.

#### Systematic uncertainties: Inner detector alignment

- ID misalignment effects: dominated by global length scale biases originating from ID geometry radial and longitudinal distortions along track trajectory.
- These manifest themselves as a shift in the reconstructed masses of known resonances, e.g.  $J/\psi \to \mu^+\mu^-$  Eur. Phys. J. C 80 (2020) 1194 7.

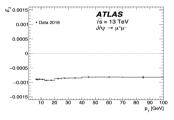


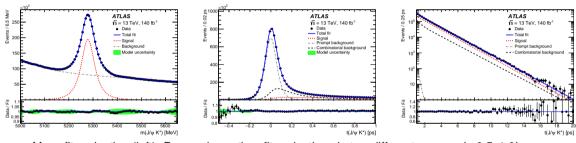
Fig showing a bias of the  $J/\psi \to \mu^+\mu^-$  mass, measured as a relative difference from PDG value, as a function of a muon track  $p_{\Gamma}$ .

- In our analysis  $B_d^0$  tracks are re-fitted with the  $J/\psi$  mass constrained to PDG, effectively removing the misalignment effect. The impact of misalignment estimated by alternative fit:  $B_d^0$  vertex re-fitted without PDG constrain.
- Additionally, to account for the momentum bias affecting hadrons from  $K^{*0}$  their  $p_T$  are altered by -0.085% Eur. Phys. J. C 80 (2020) 1194  $\checkmark$ .
- The two effects summed in quadrature, give the systematics: 0.9  $\sigma_{stat}$

### Results: The $B^0$ effective lifetime and the mass and time projections of the likelihood fit

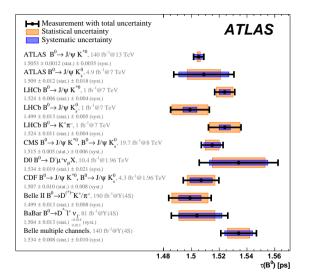
The  $B^0$  effective lifetime value measured with a total of 2450 500  $\pm$  2400  $B^0 \to J/\psi K^{*0}$  signal events The measured effective lifetime is

$$\tau = 1.5053 \pm 0.0012$$
 (stat.)  $\pm 0.0035$  (syst.) ps.



- Mass fit projection (left). Proper decay time fit projections in two different ranges: (-0.5; 1.0) ps (Middle) and (1; 20) ps (Right).
- Solid blue line total fit, dashed red line signal.
- The lower panels: ratio of data point to the fit value. The green band the envelope of model variations included in the systematic uncertainty, the bars on the data points indicate statistical uncertainties. Plot -right the model variation band too small to be visible.

#### A comparison ATLAS $B^0$ lifetime with the latest results of other experiments.



- The current ATLAS result in  $B^0 \to J/\psi K^{*0}$  channel.
- The previous ATLAS result Phys. Rev. D 87 (2013) 032002  $\overline{C}$  in the  $B^0 \to J/\psi K_0^0$  channel.
- Latest LHCb results JHEP 04 (2014) 114  $\mathbb{C}^*$  in  $B^0 \to J/\psi K^{*0}$  and  $B^0 \to J/\psi K^0_S$  decays and Phys. Lett. B 736 (2014) 446  $\mathbb{C}^*$  in the  $B^0 \to K^+\pi^-$ .

- $e^+e^-$  colliders: Belle II Phys. Rev. D 107 (2023) L091102  $\mathbb C^*$  in the  $B^0 \to D^{(*)-}K^+/\pi^+$  channel and the last result from BaBar Phys. Rev. D 73 (2006) 012004  $\mathbb C^*$  in the  $B^0 \to D^{*+}\ell^-\bar{\nu}_\ell$ . Belle PhysRevD.71.07990, 2005  $\mathbb C^*$  this combination includes  $B^0_d$  decays to  $D^{*-}\ell^+\nu$ ,  $D^*-\pi^+$ ,  $D^-\pi^+$ ,  $D^*-\rho^+$ ,  $J/\psi K^{*0}$ ,  $J/\psi K^0_o$ .

### Results: Determination of the $B_d^0$ average decay width $\Gamma_d$ and the ratio $\Gamma_d/\Gamma_s$

#### $\Gamma_d$

• We determine  $\Gamma_d$  from our measured effective lifetime  $\tau_{B^0}$ , using Eq (1) and input values  $2y = \Delta \Gamma_d / \Gamma_d = 0.001 \pm 0.010$  from HFLAV:2023  $\Box$ :

$$\Gamma_d = 0.6639 \pm 0.0005 \text{ (stat.)} \pm 0.0016 \text{ (syst.)} \pm 0.0038 \text{ (ext.)} \text{ ps}^{-1}$$

The uncertainty denoted 'ext.' originates from the HFLAV.

• The value  $\Gamma_d$  is in agreement with HQE theory of  $0.63^{+0.11}_{-0.07}~\mathrm{ps}^{-1}$  Lenz et al. 2023  $\checkmark$ 

#### $\Gamma_d/\Gamma_s$

• Using  $\Gamma_s = 0.6703 \pm 0.0014$  (stat.)  $\pm 0.0018$  (syst.) ps<sup>-1</sup> measured by the ATLAS Eur. Phys. J. C 81 (2021) 342  $\checkmark$  the resulting  $\Gamma_d/\Gamma_s$  ratio is

$$\Gamma_d/\Gamma_s = 0.9905 \pm 0.0022 \text{ (stat.)} \pm 0.0036 \text{ (syst.)} \pm 0.0057 \text{ (ext.)}$$

- The statistical, systematic and external uncertainties are propagated from the quantities above. In  $\Gamma_d/\Gamma_s$  systematic uncertainties of the ATLAS measurements of  $\tau_{B^0}$  and  $\Gamma_s$  primarily come from different sources. They are therefore treated as uncorrelated.
- $\Gamma_d/\Gamma_s$  agrees with theory HQE and lattice QCD models prediction.

# Measurement of the relative width difference $\Delta\Gamma_d/\Gamma_d$ of the $B_d^0$ - $B_d^0$ system with ATLAS JHEP06(2016)081 $\Box$

- The value of  $\Delta\Gamma_d/\Gamma_d$  is obtained by comparing the decay-time distributions of  $B^0 \to J/\psi K^{*0}$  and  $B^0 \to J/\psi K^0_S$  decays.
- Ratio of measured Proper-decay lengths distributions of the two processes is fitted using corresponding decay rate functions, with  $\Delta \Gamma_d/\Gamma_d$  as a free parameter of the fit and using inputs values from  $1/\Gamma_d$ ,  $\Delta m_s$  and  $\sin 2\beta$  available at that time 2014.
- Using 25.1  $fb^{-1}$  of 2011 and 2012 data the result is  $\Delta\Gamma_d/\Gamma_d = (-0.1 \pm 1.1 \text{ (stat.)} \pm 0.9 \text{ (syst.)})x$  10  $^{-2}$ .
- This is in agreement with theory SM value: (0.42  $\pm$  0.08) x 10<sup>-2</sup> JHEP06(2016)081  $\red C$ Lenz, Nierste.
- The word average measurement is currently (2025)  $\Delta\Gamma_d/\Gamma_d = (-0.1 \pm 1.0) \times 10^{-2}$ . The list of individual measurements with ATLAS result still leading in precision.

$3.4 \pm 2.3 \pm 2.4$	<sup>1</sup> SIRUNYAN	2018BY	CMS	pp at 8 TeV
$-0.1 \pm 1.1 \pm 0.9$	<sup>2</sup> AABOUD	2016G	ATLS	$\it pp$ at 7, 8 TeV
$-4.4 \pm 2.5 \pm 1.1$	<sup>3</sup> AAIJ	2014E	LHCB	$\it pp$ at 7 TeV
$1.7 \pm 1.8 \pm 1.1$	<sup>4</sup> HIGUCHI	2012	BELL	$e^+ \; e^-  ightarrow \varUpsilon(4S)$
$0.8 \pm 3.7 \pm 1.8$	<sup>5</sup> AUBERT,B	2004C	BABR	$e^+  e^-  o \varUpsilon(4S)$

#### Measurement of effective lifetime in $B_s^0 \to \mu^+\mu^-$ JHEP 09 (2023) 199

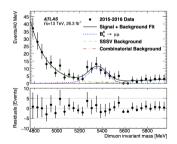
- In Standard Model (SM) only the *CP*-odd heavy-mass eigenstate in the  $B_s^0 \bar{B}_s^0$  pair decays into  $\mu^+\mu^-$  M. Beneke JHEP 10 (2019) 232 and errata JHEP 11 (2022) 099.
- The effective  $B^0_{(s)} o \mu^+ \mu^-$  lifetime is defined as

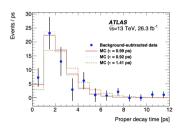
$$\tau_{\mu\mu} = \frac{\int_{0}^{\infty} t \Gamma(B_{s}^{0}(t) \to \mu\mu) \, dt}{\int_{0}^{\infty} \Gamma(B_{s}^{0}(t) \to \mu\mu) \, dt}, \text{ where: } \Gamma\left(B_{s}\left(t\right) \to \mu\mu\right) = \Gamma\left(B_{s}^{0}\left(t\right) \to \mu\mu\right) + \Gamma\left(\bar{B}_{s}^{0}\left(t\right) \to \mu\mu\right)$$

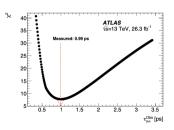
- and t is the proper decay time of the  $B_s^0$  and  $\bar{B}_s^0$  mesons.
- In the SM hypothesis  $au_{\mu\mu}$  coincides with the lifetime of the heavy  $B_s^0$  eigenstate  $au_{B_s^H}$ .
- The experimental average of the  $B^0_s$   $\bar{B}^0_s$  lifetimes and their difference Phys. Rev. D 107 (2023) 052008  $\ref{2008}$  yields the prediction  $au^{SM}_{\mu\mu}=(1.624\pm0.009)\,$  ps, with new physics effects perturbing it at most by the difference between the heavy and light eigenstate lifetimes (0.193 ps ).

## Measurement of effective lifetime in $B^0_{(s)} \to \mu^+ \mu^-$ , cont 1

- Data from 2015-2016 are used in this measurement.
- Un-binned maximum likelihood fit to candidates in the [4766 5966] MeV mass region (Left Fig) , yielding 58  $\pm$  13 (stat. only)  $\mathcal{B}^0_{(s)} \to \mu^+\mu^-$  signal events.
- Signal and backgrounds weights calculated from the result of the mass fit are used to construct the proper decay time data histogram background-subtracted employing per-event weights calculated according to the *sPlot* technique (Middle Fig).
- The lifetime measurement is obtained by minimising the binned  $\chi^2$  between the data histogram and lifetime-dependent pure signal MC templates extracted from MC simulated samples, as illustrated in the Middle and Right Fig.

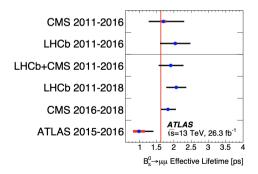






# Measurement of effective lifetime in $B^0_{(s)} o \mu^+\mu^-$ , cont 2

- The result using 2015-2016 data is  $\tau_{\mu\mu}=0.99^{+0.42}_{-0.07}$  (stat.)  $\pm$  0.17 (syst.) ps. It is consistent with the SM prediction  $\tau^{\rm SM}_{\mu\mu}=$  (1.624  $\pm$  0.009) ps Phys. Rev. D 107 (2023) 052008  $\ref{2008}$  as well as with the other available experimental results.
- The systematic errors are dominated by data-MC discrepancies, followed by uncertainties in backgrounds lifetime modelling.



Measurement of the CPV phase, width difference  $\Delta\Gamma_s$  and the average decay width  $\Gamma_s$  in  $B_s \to J/\psi \phi$ . Eur. Phys. J. C 81 (2021) 342 🗸

#### Physics motivation

- $B_s^0 \to J/\psi \phi$  is used to measure CP-violation phase  $\phi_s$  potentially sensitive to New Physics (NP)
  - In SM  $\phi_s$  is related to the CKM elements  $\phi_s \simeq 2 \arg[-(V_{ts}V_{tb}^*)/(V_{cs}V_{cb}^*)]$  and predicted with high precision PhysRevD.91.073007
- The same analysis also measure  $\Gamma_s^L$  and  $\Gamma_s^H$  the decay widths of  $B_s^0$  mass eigenstates and  $\Delta\Gamma_s = \Gamma_s^L \Gamma_s^H$ .  $\Delta\Gamma_s$  was calculated in SM Lenz-et-al 2020  $\Gamma$  and new experimental results are needed to eventually get sensitivity to NP models that predict smaller  $\Delta\Gamma_s$  than SM LenzOct2021  $\Gamma$

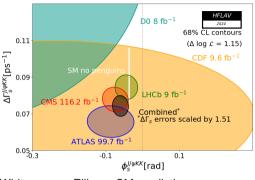
#### More analysis in preparation

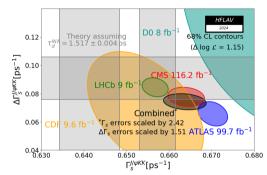
- Total Run2 data including remaning 58.5 fb<sup>-1</sup>
- Followed by Run3 analysis and Run2+Run3 comb.

# ATLAS result with partial Run2 data 80.5 fb<sup>-1</sup> combined with 19.2 fb<sup>-1</sup> of Run1 data

	Solution (a)					
Parameter	Value	Statistical	Systematic			
		uncertainty	uncertainty			
$\phi_s$ [rad]	-0.087	0.036	0.021			
$\Delta\Gamma_s$ [ps <sup>-1</sup> ]	0.0657	0.0043	0.0037			
$\Gamma_s$ [ps <sup>-1</sup> ]	0.6703	0.0014	0.0018			
$ A_{\parallel}(0) ^2$	0.2220	0.0017	0.0021			
$ A_0(0) ^2$	0.5152	0.0012	0.0034			
$ A_S ^2$	0.0343	0.0031	0.0045			
$\delta_{\perp}$ [rad]	3.22	0.10	0.05			
δ <sub>∥</sub> [rad]	3.36	0.05	0.09			
$\delta_{\perp} - \delta_{S}$ [rad]	-0.24	0.05	0.04			

#### $B_s^0 \to J/\psi \phi$ latest world combination 2024





White narrow Ellipse: SM predictions.

Grey areas: SM predictions

Because of tensions between the measurements, the errors on  $\Gamma_s$  and  $\Delta\Gamma_s$  have been scaled (the ellipses representing the results of each experiment are shown before scaling, while the combined ellipses include the scale factors).

#### High precision test of SM in $B_s^0 \to J/\psi \phi$ : experiments w.r.t SM

#### CP violation phase $\phi_s$ potentially sensitive to New Physics

- SM prediction  $\phi_s = -0.0367^{+0.00072}_{-0.00082}$  rad M. Bona et al. (UTfit), Rend. Lincei Sci. Fis. Nat. 34, 27 (2023) is cca 17 more precise than World combined 2024  $\phi_s = -0.060 \pm 0.014$  rad. HFLAV 2024  $\checkmark$
- There is a room for New physics. An answer is on experimental side: Run3 LHC and HL-LHC

#### Case of $\Delta\Gamma_s$ and $\Gamma_s$

- SM prediction ΔΓ<sub>s</sub>= + 0.091±0.015ps<sup>-1</sup> Albrecht, Bernlochner, Lenz, Rusov 2023 
   slightly higher than World Combined value.
- A potential New Physics enhancement of  $\phi_s$  would also decrease  $\Delta\Gamma_s$ ,

- Lenz1 et al. 🗹
- "Precise experimental knowledge on B-lifetimes and their Ratios, will provide bounds on New Physics LenzOct2021
- Currently tensions at level  $2\sigma$  in  $\Delta\Gamma_s$  between ATLAS CMS and  $5\sigma$  in  $\Gamma_s$  between ATLAS LHCb. More data to be used to improve lifetime measurement precision and control of systematic effects.
- All experiments still to add Run3 and HL-LHC in preparations.

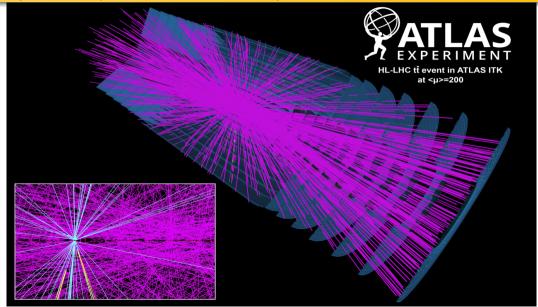
# ATLAS B-decays and lifetimes at HL-LHC

HL-LHC



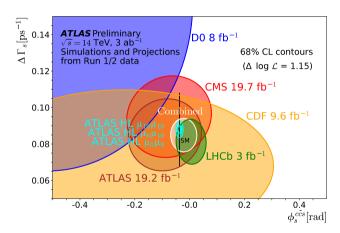
- Increase  $> 10 \text{ x} \int \text{Ldt of LHC} \rightarrow 3000\text{-}4000 \text{ fb}^{-1}$
- Peak luminosity 5 7.5 x 10<sup>34</sup> cm<sup>-2</sup> s<sup>-1</sup>
- Average amount of pp interactions 140-200 per BX with a time space 25 ns
- These conditions require Detector Upgrades.

#### High Luminosity-LHC - ATLAS track density in Inner detector



### ATLAS HL-LHC prospects $B_s^0 o J/\psi \phi$

- ATL-PHYS-PUB-2018-041
- Inner Detector upgrade: proper decay time resolution improved by 21% w.r.t. Run 2
- Three trigger scenarios for muon momenta thresholds
- $\phi_s$  precision improves (9 20) times w.r.t.Run1, or (4 9) times w.r.t. current result combining Run1 and Run2 99.7 fb<sup>-1</sup>



 $B^0_{(s)} o \mu^+\mu^-$  HL-LHC Prospects, ATL-PHYS-PUB-2025-016

#### 3 trigger scenarios for thresholds $p_T(\mu_1)$ , $p_T(\mu_2)$

Quantity	$\mid \mathcal{B}(B_s^0 \to \mu^+ \mu^-) [10^{-9}] \mid \mathcal{B}(B^0 \to \mu^+ \mu^-) [10^{-1}]$				$[10^{-10}]$	$ au_{\mu\mu}  ext{ [ps]}$			
SM value	3.66 ± 0.14 [1]			1.03 ± 0.05 [1]			1.624 ± 0.009 [7]		
ATLAS 2015–2016 measurements	3.2	1+0.96+0.49 1-0.91-0.30	[2]	< 4.3 at 95% CL. [2]			$0.99^{+0.42}_{-0.07} \pm 0.17$ [8]		
Projected uncertainty	Stat.	Syst.	Total	Stat.	Syst.	Total	Stat.	Syst.	Total
$\mu10\mu10$	0.17	0.36	0.40	0.36	0.32	0.48	+0.09 -0.06	0.06	+0.11 -0.08
μ10μ6	0.11	0.32	0.34	0.23	0.26	0.35	+0.06 -0.03	0.05	+0.08 -0.06
μ6μ6	0.09	0.32	0.33	0.19	0.26	0.32	+0.05 -0.03	0.05	+0.07 -0.05

