New observables in non-leptonic B decays

Gilberto Tetlalmatzi-Xolocotzi

Based on: *JHEP 06 (2023) 108 (2301.10542)*

2404.01186

In collaboration with: A. Biswas, S. Descotes-Genon, Q. Matias

CPPS, Theoretische Physik 1,
Universität Siegen

Université Paris-Saclay, CNRS/IN2P3,
IJCLab,

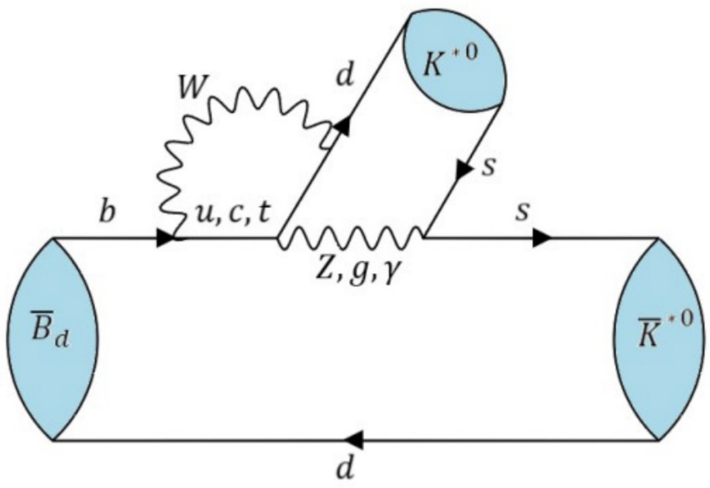
Center for Particle Physics Siegen

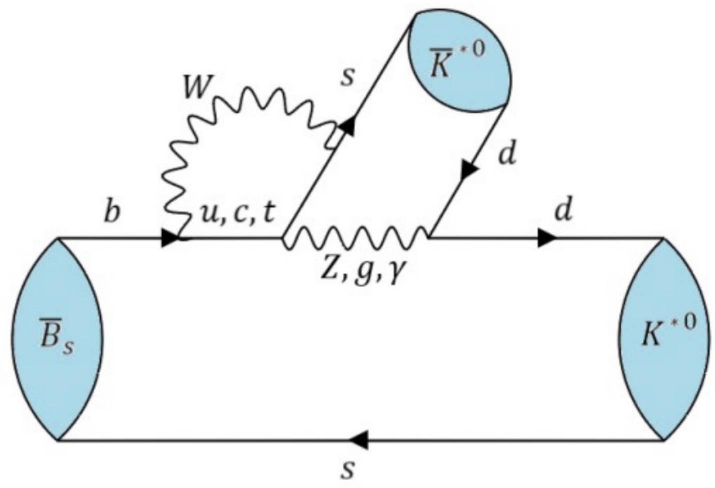
TRR 257 - Particle Physics Phenomenology after the Higgs Discovery

Universität Siegen

IJCLab Irène Joliot-Curie Laboratoire de Physique des 2 Infinis

Relevant decay processes





Penguin induced processes: b ► d and b ► s

$$\bar{A}_f = \lambda_u^{(q)} T_q + \lambda_c^{(q)} P_q = \lambda_u^{(q)} \Delta_q - \lambda_t^{(q)} P_q \qquad q = d\,,\, s$$

$$\lambda_u^{(q)} + \lambda_c^{(q)} + \lambda_t^{(q)} = 0$$

Free from infrared divergences $\quad \Delta_q = T_q - P_q$

Amplitude structure

The amplitudes are calculated using QCDF

Beneke et. al. 0308039 [hep-ph]

$$T(\bar{B}_d \to \bar{K}^{\,0} K^{\,0}) = A_{\bar{K}\,K}\,[\alpha_4^u - \frac{1}{2}\alpha_{4,EW}^u + \beta_3^u + \beta_4^u - \frac{1}{2}\beta_{3,EW}^u - \frac{1}{2}\beta_{4,EW}^u]$$
$$+A_{K\,\bar{K}}\,[\beta_4^u - \frac{1}{2}\beta_{4,EW}^u]\,,$$
$$P(\bar{B}_d \to \bar{K}^{\,0} K^{\,0}) = A_{\bar{K}\,K}\,[\alpha_4^c - \frac{1}{2}\alpha_{4,EW}^c + \beta_3^c + \beta_4^c - \frac{1}{2}\beta_{3,EW}^c - \frac{1}{2}\beta_{4,EW}^c]$$
$$+A_{K\,\bar{K}}\,[\beta_4^c - \frac{1}{2}\beta_{4,EW}^c]\,,$$

The infrared divergences in *T* and *P* have the same structure

$$\alpha_4^p(M_1 M_2) = a_4^p(M_1 M_2) + r_\chi^{M_2}\, a_6^p(M_1 M_2)$$

Amplitude structure

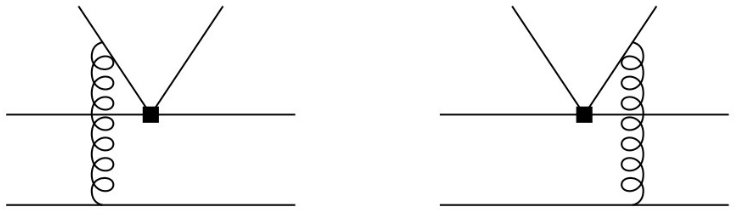
$$a_i^p(M_1 M_2) = \left(C_i + \frac{C_{i\pm1}}{N_c}\right) N_i(M_2) + \frac{C_{i\pm1}}{N_c}\,\frac{C_F\alpha_s}{4\pi}\left[V_i(M_2) + \frac{4\pi^2}{N_c} H_i(M_1 M_2)\right] + P_i^p(M_2)\,,$$

Vertex corrections

Amplitude structure

$$a_i^p(M_1M_2) = \left(C_i + \frac{C_{i\pm1}}{N_c}\right) N_i(M_2) + \frac{C_{i\pm1}}{N_c}\frac{C_F\alpha_s}{4\pi}\left[V_i(M_2) + \frac{4\pi^2}{N_c} H_i(M_1M_2)\right] + P_i^p(M_2)\,,$$

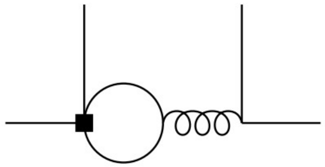
Spectator scattering

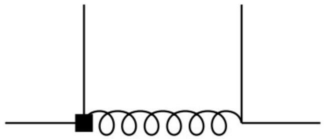


Amplitude structure

$$a_i^p(M_1M_2) = \left(C_i + \frac{C_{i\pm1}}{N_c}\right)N_i(M_2) + \frac{C_{i\pm1}}{N_c}\frac{C_F\alpha_s}{4\pi}\left[V_i(M_2) + \frac{4\pi^2}{N_c}H_i(M_1M_2)\right] + P_i^p(M_2),$$

Penguin contributions

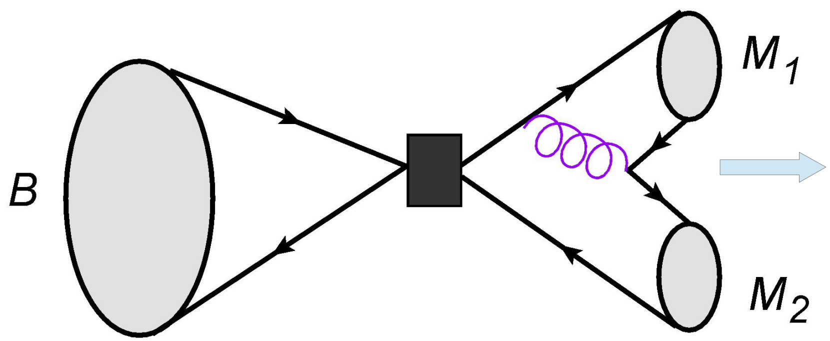




Annihilation contributions

Weak annihilation topolgies are affected by LCDA end point singularities

$$\int_0^1 \frac{dy}{y} \to X_A$$



β, b

Educated Ansatz

$$X_A = (1 + \rho_A e^{i\varphi_A}) \ln \frac{m_B}{\Lambda_h}$$

$$0 \le \rho_A \le 1, \quad 0 \le \varphi_A \le 2\pi$$

Data can be used to get bounds for these contributuions

GTX, T. Huber 2111.06418 [hep-ph]

For decays of *B* mesons into pseudoscalar-pseudoscalar final states

$$|\beta_i| \le 0.40 \quad \text{for} \quad |\alpha_1| \approx 1$$

SU(3) flavour exact

Another possibility is to construct observables with low sensitivity to these contributions

Optimized Observables

$$L_{K\bar{K}} = \rho(m_{K^0}, m_{K^0}) \frac{\mathcal{B}(\bar{B}_s \to K^0 \bar{K}^0)}{\mathcal{B}(\bar{B}_d \to K^0 \bar{K}^0)} = \frac{|A^s|^2 + |\bar{A}^s|^2}{|A^d|^2 + |\bar{A}^d|^2}$$

U-spin related channels

Optimized Observables

$$L_{K\bar{K}} = \rho(m_{K^0}, m_{K^0}) \frac{\mathcal{B}(\bar{B}_s \to K^0 \bar{K}^0)}{\mathcal{B}(\bar{B}_d \to K^0 \bar{K}^0)} = \frac{|A^s|^2 + |\bar{A}^s|^2}{|A^d|^2 + |\bar{A}^d|^2}$$

↑

Phase space factor

Optimized Observables

$$L_{K\bar{K}} = \rho(m_{K^0}, m_{K^0}) \frac{\mathcal{B}(\bar{B}_s \to K^0\bar{K}^0)}{\mathcal{B}(\bar{B}_d \to K^0\bar{K}^0)} = \frac{|A^s|^2 + |\bar{A}^s|^2}{|A^d|^2 + |\bar{A}^d|^2}$$

$$L_{K^*\bar{K}^*} = \rho(m_{K^{*0}}, m_{K^{*0}}) \frac{\mathcal{B}(\bar{B}_s \to K^{*0}\bar{K}^{*0})}{\mathcal{B}(\bar{B}_d \to K^{*0}\bar{K}^{*0})} \frac{f_L^{B_s}}{f_L^{B_d}} = \frac{|A_0^s|^2 + |\bar{A}_0^s|^2}{|A_0^d|^2 + |\bar{A}_0^d|^2}$$

S. Descotes, J. Matias, et al 2011.07867 [hep-ph]

Only the longitudinal component is factorizable (vector-vector final states)

$$L_{K^*\bar{K}^*} = \kappa \left|\frac{P_s}{P_d}\right|^2 \left[\frac{1 + |\alpha^s|^2 \left|\frac{\Delta_s}{P_s}\right|^2 + 2\mathrm{Re}\left(\frac{\Delta_s}{P_s}\right)\mathrm{Re}(\alpha^s)}{1 + |\alpha^d|^2 \left|\frac{\Delta_d}{P_d}\right|^2 + 2\mathrm{Re}\left(\frac{\Delta_d}{P_d}\right)\mathrm{Re}(\alpha^d)}\right]$$

$$\begin{aligned}
\alpha^d &= (-0.0136^{+0.0095}_{-0.0096}) + i(0.4181^{+0.0085}_{-0.0064}),\\
\alpha^s &= (0.00863^{+0.00040}_{-0.00036}) + i(-0.01829^{+0.00037}_{-0.00042}),\\
\kappa &= \left|\frac{\lambda_c^s + \lambda_u^s}{\lambda_c^d + \lambda_u^d}\right|^2 = 22.92^{+0.52}_{-0.30}.
\end{aligned}$$

SU(3) $\left|\frac{P_s}{P_d}\right| = 1 \pm 0.3$

Naive Factorization $\left|\frac{P_s}{P_d}\right| = 0.91^{+0.20}_{-0.17}$

QCD Factorization $\left|\frac{P_s}{P_d}\right| = 0.92^{+0.20}_{-0.18}$

Optimized Observables

Under the assumption of *U-Spin symmetry* in the sources of infrared divergences we expect that they fluctuate in the same direction in numerator and denominator

Theory prediction

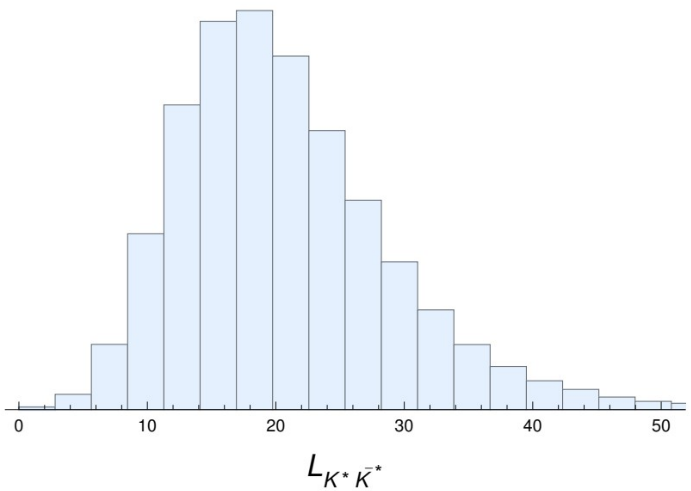
$$\text{naive } SU(3) : L_{K^*\bar{K}^*} = 23^{+16}_{-12} \qquad 1.9\sigma\,,$$

$$\text{fact } SU(3) : L_{K^*\bar{K}^*} = 19.2^{+9.3}_{-6.5} \qquad 3.0\sigma\,,$$

$$\text{QCD fact} : L^{\rm SM}_{K^*\bar{K}^*} = 19.53^{+9.14}_{-6.64} \qquad 2.6\sigma\,,$$

S. Descotes, J. Matias, et al
2011.07867 [hep-ph]

A. Biswas, S. Descotes, J. Matias, GTX
2301.10542 [hep-ph]



Optimized Observables

Experimental measurement

$$\frac{\mathcal{B}_{B_d\to K^{*0}\bar{K}^{*0}}}{\mathcal{B}_{B_s\to K^{*0}\bar{K}^{*0}}} = 0.0758 \pm 0.0057(\text{stat}) \pm 0.0025(\text{syst}) \pm 0.0016 \left(\frac{f_s}{f_d}\right)$$

LHCb [1905.06662, 0708.2248]

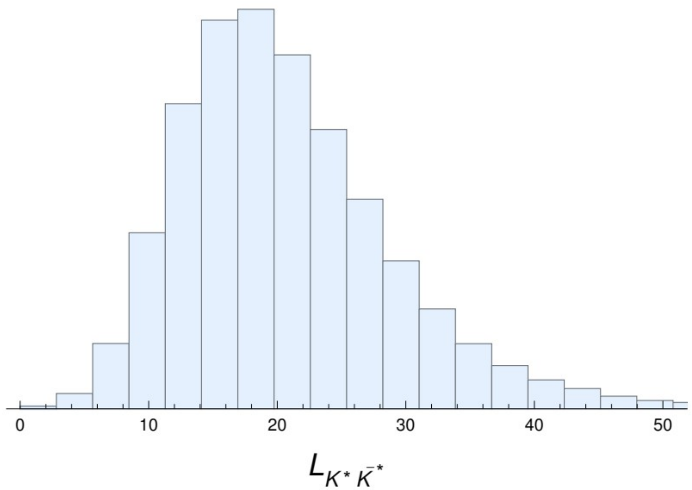
Measured longitudinal polarisation fractions	
$f_L(\bar{B}_d \to K^{0*}\bar{K}^{0*})$	$f_L(\bar{B}_s \to K^{0*}\bar{K}^{0*})$
0.73 ± 0.05	0.240 ± 0.040

LHCb [1905.06662], BABAR [0708.2248]
S. Descotes, J. Matias, et al 2011.07867 [hep-ph]

$$L^{\text{exp}}_{K^*\bar{K}^*} = 4.43 \pm 0.92$$

$2.6\ \sigma$ tension

LHCb [1503.05362]



Optimized Observables

Experimental measurement

$$\frac{\mathcal{B}_{B_d\to K^{*0}\bar{K}^{*0}}}{\mathcal{B}_{B_s\to K^{*0}\bar{K}^{*0}}} = 0.0758 \pm 0.0057(\text{stat}) \pm 0.0025(\text{syst}) \pm 0.0016 \left(\frac{f_s}{f_d}\right)$$

LHCb [1905.06662, 0708.2248]

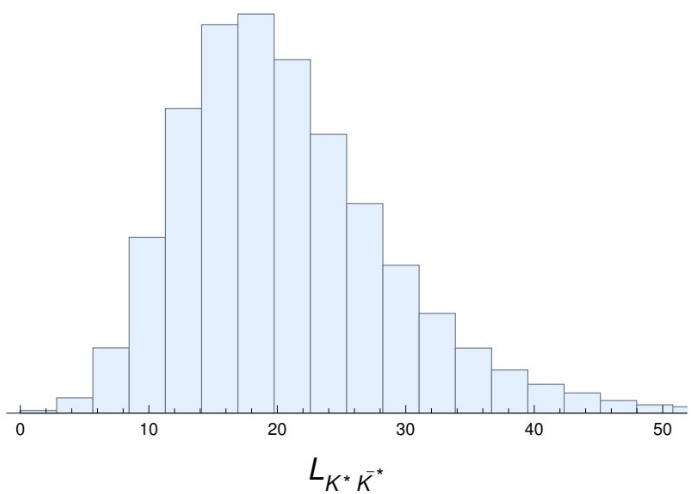
Measured longitudinal polarisation fractions	
$f_L(\bar{B}_d \to K^{0*}\bar{K}^{0*})$	$f_L(\bar{B}_s \to K^{0*}\bar{K}^{0*})$
0.73 ± 0.05	0.240 ± 0.040

LHCb [1905.06662], BABAR [0708.2248]
S. Descotes, J. Matias, et al 2011.07867 [hep-ph]

$$L^{\text{exp}}_{K^*\bar{K}^*} = 4.43 \pm 0.92$$

$2.6\ \sigma$ tension

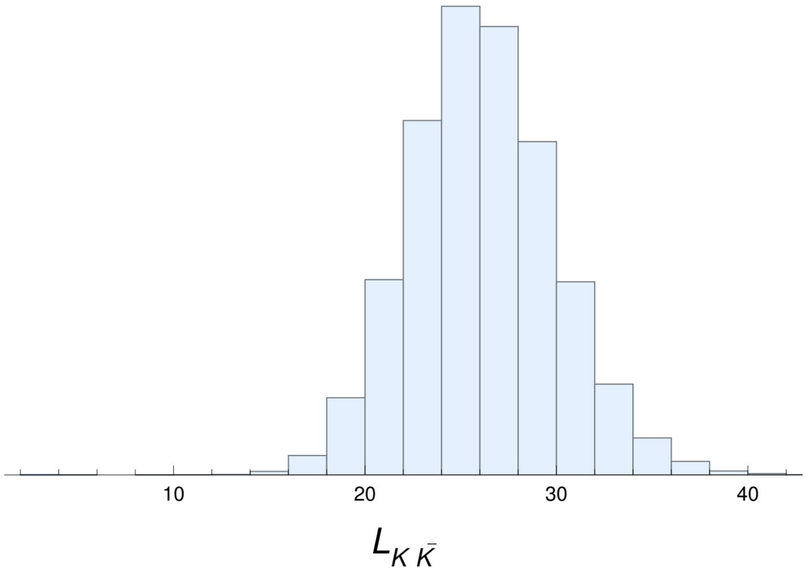
LHCb [1503.05362]



Optimized Observables

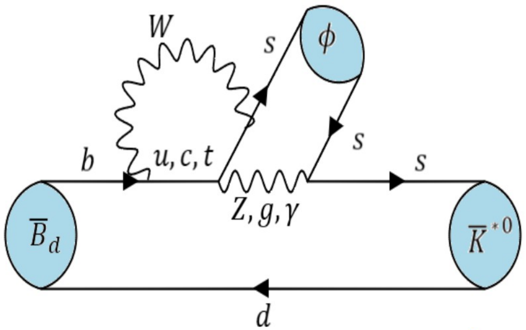
$$L_{K\bar{K}} = \rho(m_{K^0}, m_{K^0}) \frac{\mathcal{B}(\bar{B}_s \to K^0 \bar{K}^0)}{\mathcal{B}(\bar{B}_d \to K^0 \bar{K}^0)} = \frac{|A^s|^2 + |\bar{A}^s|^2}{|A^d|^2 + |\bar{A}^d|^2}$$

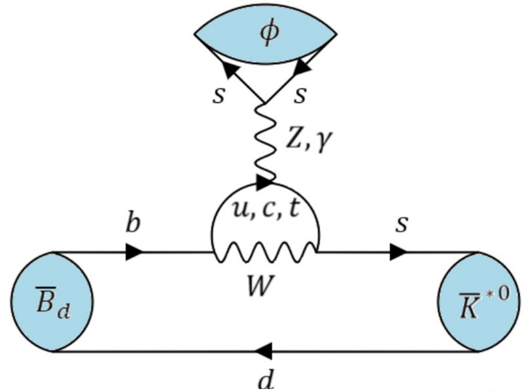
$L_{K\bar{K}}^{\mathrm{SM}} = 26.00^{+3.88}_{-3.59}$ $\quad$ $L_{K\bar{K}}^{\mathrm{exp}} = 14.58 \pm 3.37$ $\quad$ $2.4\ \sigma$ $\quad$ tension



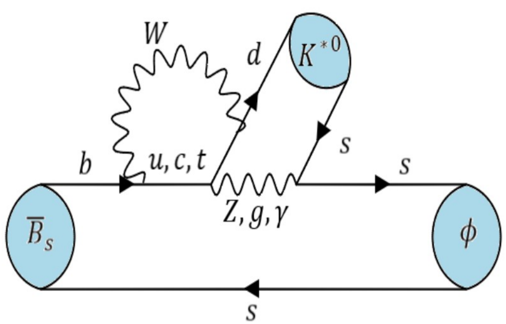
Optimized Observables

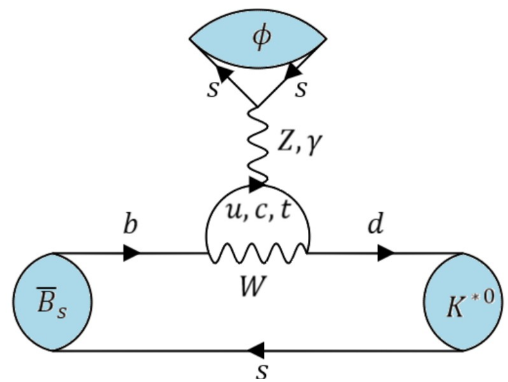
$$B_{d(s)} \to K^{(*)0}\phi$$ processes





$$A(\bar{B}_d \to \bar{K}^{*0}\phi) = \sum_{p=u,c} \lambda_p^{(s)} \left(a_4^p + a_3 + a_5 - \frac{1}{2}(a_7^p + a_9^p + a_{10}^p) \right) A_{K^*\phi}$$
$$+ \left(\lambda_u^{(s)} + \lambda_c^{(s)} \right) \left(b_3 - \frac{1}{3} b_3^{\mathrm{EW}} \right) B_{K^*\phi}$$

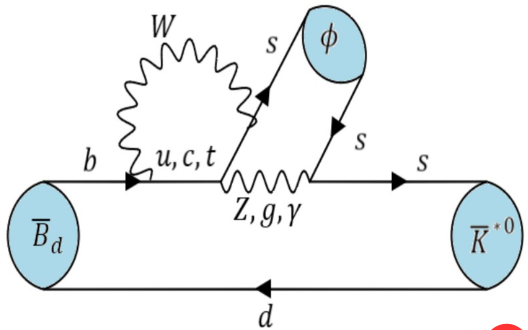


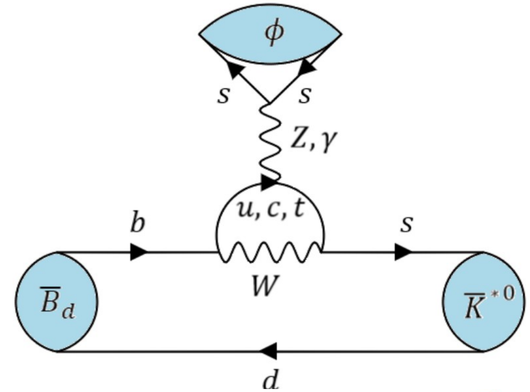


$$A(\bar{B}_s \to K^{*0}\phi) = \sum_{p=u,c} \lambda_p^{(d)} \left[\left(a_4^p - \frac{1}{2} a_{10}^p \right) A_{\phi K^*} + \left(a_3 + a_5 - \frac{1}{2}(a_7^p + a_9^p) \right) A_{K^*\phi} \right]$$
$$+ \left(\lambda_u^{(d)} + \lambda_c^{(d)} \right) \left(b_3 - \frac{1}{3} b_3^{\mathrm{EW}} \right) B_{K^*\phi},$$

Optimized Observables

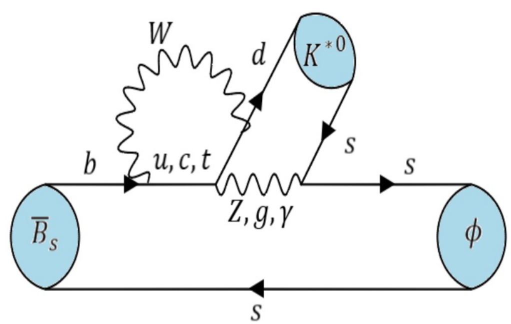
$B_{d(s)} \to K^{(*)0}\phi$ processes

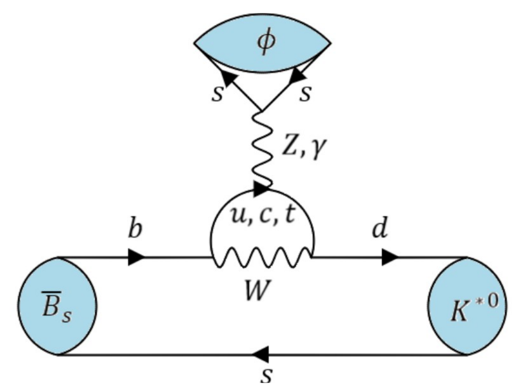




$$A(\bar{B}_d \to \bar{K}^{*0}\phi) = \sum_{p=u,c} \lambda_p^{(s)} \left(a_4^p + a_3 + a_5 - \frac{1}{2}(a_7^p + a_9^p + a_{10}^p) \right) A_{K^*\phi}$$
$$+ \left(\lambda_u^{(s)} + \lambda_c^{(s)}\right)\left(b_3 - \frac{1}{3} b_3^{\rm EW}\right) B_{K^*\phi}$$

Penguin dominance
$p=c$





$$A(\bar{B}_s \to K^{*0}\phi) = \sum_{p=u,c} \lambda_p^{(d)} \left[\left(a_4^p - \frac{1}{2} a_{10}^p \right) A_{\phi K^*} + \left(a_3 + a_5 - \frac{1}{2}(a_7^p + a_9^p) \right) A_{K^*\phi} \right]$$
$$+ \left(\lambda_u^{(d)} + \lambda_c^{(d)}\right)\left(b_3 - \frac{1}{3} b_3^{\rm EW}\right) B_{K^*\phi},$$

Optimized Observables

$$L_{K^*\phi} = \rho(m_{K^{*0}}, m_\phi)\frac{\mathcal{B}(\bar{B}_d \to \bar{K}^{*0}\phi)}{\mathcal{B}(\bar{B}_s \to K^{*0}\phi)}\frac{f_L^{B_d}}{f_L^{B_s}}$$

Measured longitudinal polarisation fractions	
$f_L(\bar{B}_d \to \bar{K}^{0*}\phi)$	$f_L(\bar{B}_s \to K^{0*}\phi)$
0.497 ± 0.017	0.51 ± 0.17

BABAR 0808.3586, BELLE 1308.1830, LHCb 1403.2888, LHCb 1306.2239

$$L_{K^*\phi}^{\rm th} = 22.04^{+7.06}_{-4.88} \qquad L_{K^*\phi}^{\rm exp} = 8.80^{+6.07}_{-2.97} \qquad 1.48\ \sigma$$

A. Biswas, S. Descotes, J. Matias, GTX
2404.01186 [hep-ph]

Extra Observables ($B \longrightarrow VP$ and $B \longrightarrow PV$)

$$\mathcal{B}(\bar{B}_d \to \bar{K}^0 \phi)^{\rm th} = (4.28^{+2.71}_{-1.50}) \times 10^{-6}$$

$$\mathcal{B}(\bar{B}_d \to \bar{K}^0 \phi)^{\rm exp} = (7.3 \pm 0.7) \times 10^{-6}$$

BABAR 1201.5897, BELLE 0307014

$$\mathcal{B}(B_s \to K^{*0} \bar{K}^0)^{\rm exp} + \mathcal{B}(B_s \to \bar{K}^{*0} K^0)^{\rm exp} = (1.98 \pm 0.28 \pm 0.50) \times 10^{-5}$$

$$\mathcal{B}(\bar{B}_s \to K^{*0} \bar{K}^0 + c.c.)^{\rm th} = (8.35^{+5.02}_{-2.51}) \times 10^{-6}$$

A. Biswas ,S. Descotes, J. Matias and GTX
2404.01186 [hep-ph]

Complementary observables are required experimentally to define the corresponding optimized ratios

Phenomenological analysis in the second part of the talk.....