Probing the \mathcal{CP} violation of the top-Yukawa coupling at future colliders

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To explain baryogenesis, CP-violating effects beyond those predicted by the Standard Model are essential. With the discovery of the Higgs boson, attention has turned to exploring the CP structure within the Higgs sector. While CP-violating Higgs couplings to vector bosons are loop-suppressed in most BSM models, such models often predict the largest amount of CP violation in Higgs-fermion couplings. The top-Yukawa coupling is particularly important, as it is the largest and most accessible one.

In this study, we use the Higgs Characterization Model and parameterize the Higgs-top coupling via

$$\mathcal{L}_{\text{top-Yuk}} = -\frac{y_t^{SM}}{\sqrt{2}} \bar{t} \left(c_t + i\gamma_5 \tilde{c}_t \right) tH, \qquad (1)$$

where y_t^{SM} is the SM top-Yukawa interaction and $c_t = 1$ as well as $\tilde{c}_t = 0$ for the SM. Events are generated using MadGraph5_aMC@NLO and are passed to Pythia8 and Delphes3 for parton showering, hadronisation, and detector simulation.

The goal of this study is to evaluate the constraints that future colliders could impose on the parameters c_t and \tilde{c}_t , and therefore on their potential sensitivity to deviations from the SM in the Higgs-top coupling. While an e^+e^- collider like the FCC-ee benefits from a very clean background, which allows precise measurements of final states with many quarks, a proton-proton collider like the FCC-hh offers the advantage of a massive amount of events and data.

In order to examine the CP structure of the top-Yukawa coupling, we generate data for the scenarios $c_t = 1$, $\tilde{c}_t = 0$ and $c_t = 0$, $\tilde{c}_t = 1$. This process results in separate data sets corresponding to purely CP-even and purely CP-odd contributions, respectively. The distribution of events across various kinematic variables is analyzed to construct likelihood functions. For example, Figure 1 shows the kinematic distribution for E_H , the energy of the produced Higgs boson, which is notably sensitive to the CP structure. These likelihoods are combined to explore the parameter space of c_t and \tilde{c}_t , enabling us to set limits on the couplings. By combining different kinematic distributions in a machine learning algorithm, we can enhance the distinction between CP-even and CPodd scenarios, which allows us to set tighter constraints on c_t and \tilde{c}_t and increase the sensitivity to CP-violating effects. An example of the resulting constraints on c_t and \tilde{c}_t



Figure 1: Distributions of the kinematic variable E_H for the two scenarios $c_t = 1$, $\tilde{c}_t = 0$ and $c_t = 0$, $\tilde{c}_t = 1$. The distributions illustrate the differences between the $C\mathcal{P}$ -even and $C\mathcal{P}$ -odd cases.



Figure 2: Limits derived from the classifier trained to distinguish between the scenarios $c_t = 1$, $\tilde{c}_t = 0$ and $c_t = 0$, $\tilde{c}_t = 1$. The 1-, 2-, and 3- σ confidence level contours are marked by white, grey, and black dashed lines, respectively.