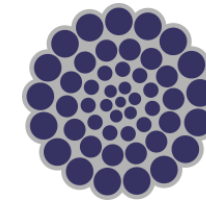


Non-stabilizerness (Magic) in the Asymmetric Quantum Rabi Model

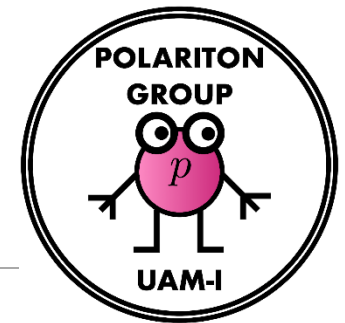


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CIENCIAS Y TECNOLOGÍAS

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WORKSHOP: GEOMETRY OF QUANTUM DYNAMICS

AUGUST 30, 2024

Contents

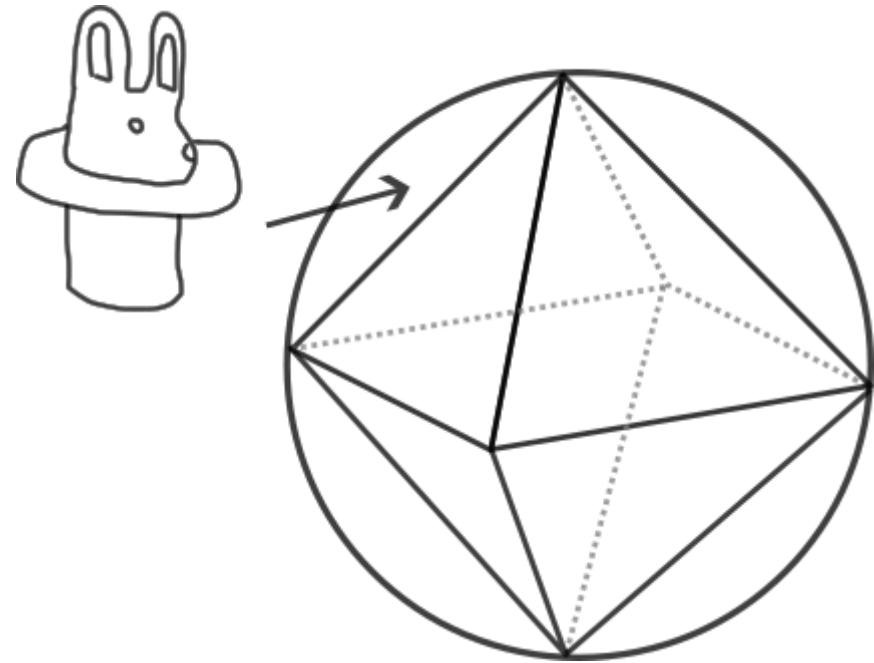
1. Non-Stabilizerness or Magic
2. The Asymmetric Quantum Rabi Model
3. Magic + AQRM = Magical AQRM
4. Results



Non-Stabilizerness or Magic

What is magic?

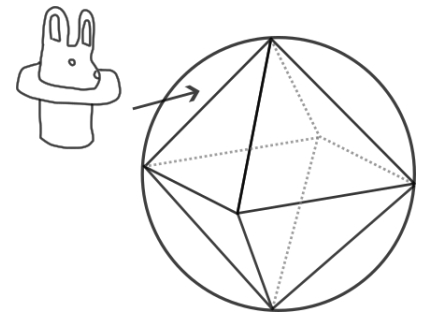
- ❑ Magic or Non-Stabilizerness is a resource for quantum computation, which differentiates quantum from classical computation.
- ❑ A resource theory for non-stabilizerness is necessary in order to work within a theoretical frame.
- ❑ Resource theories of magic in discrete and continuous regimes have been developed.



New J. Phys. 16, 013009 (2014), Phys. Rev. A 98, 052350 (2018)

Quantifying magic

- ❑ Quantifying magic has been an issue to tackle since its proposal in the fault tolerant quantum computation context.
- ❑ Distillation of magic states was the first process that used the resource theory of non-stabilizerness.
- ❑ Efforts to understand the meaning of magic in other branches of Physics, like many body Physics, are currently under deep scrutiny.



Appl. Phys. B 86 367 (2007), PRL 118 090501 (2017), PRL 128 050402 (2022)

Quantifying magic

□ Relative entropy of magic

$$r_M(\rho) = \min [S(\rho || \sigma)]_{\sigma \in STAB(H_d)}$$

ρ - state of the system

$STAB(H_d)$ - set of stabilizer states of the space

d - dimension of the Hilbert space

New J. Phys. 16, 013009 (2014)

Quantifying magic

□ Sum negativity

$$sn(\rho) = \sum_{W_\rho(u) < 0} \frac{1}{2} |W_\rho(u)| = \frac{1}{2} \left[\sum_u |W_\rho(u)| - 1 \right],$$

□ Mana

$$mana(\rho) = \mathcal{M}(\rho) = \log [2sn(\rho) + 1]$$

$W_\rho(u)$ - element of the discrete Wigner function of ρ

New J. Phys. 16, 013009 (2014)

Quantifying magic

There is not an ultimate measure of magic!!

□ Dai's M Witness

$$M(\rho) = \sum_{\mu, \nu} \sqrt{\text{tr}(\rho D_{\mu, \nu}) \text{tr}(\rho D_{\mu, \nu}^\dagger)}$$

$D_{\mu, \nu}$ - displacement operator of Heisenberg-Weyl algebra

Monotones and witnesses

not increasing under free operations

New J. Phys. 16, 013009 (2014), Int. J. Theor. Phys. 61 35 (2022)

Quantifying magic

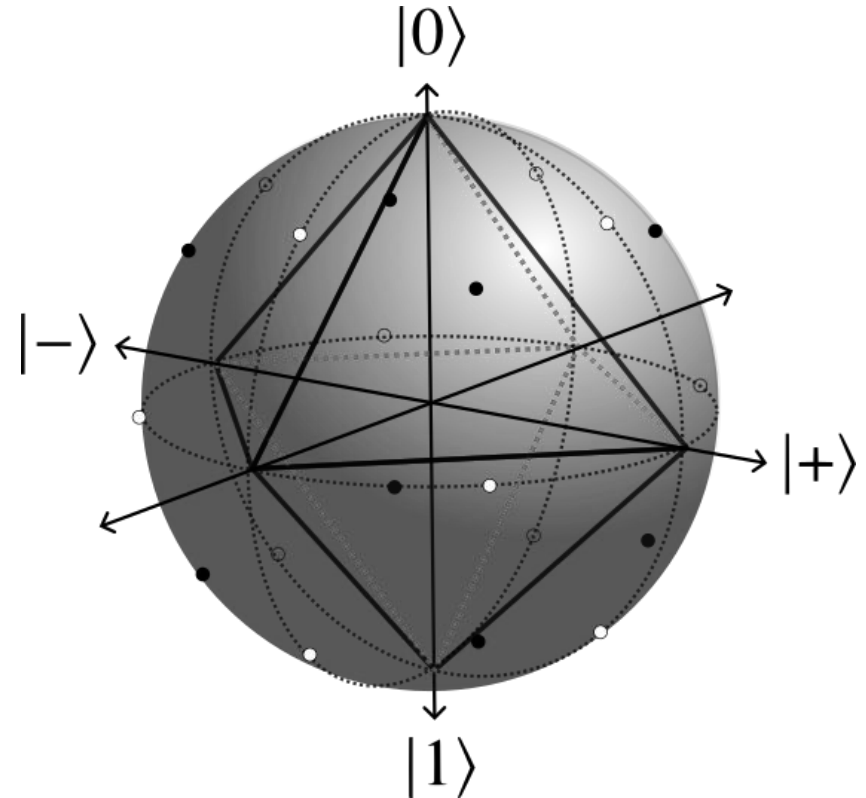
□ Bloch sphere and octahedron

The white dots are the H -type states

$$|H\rangle\langle H| = \frac{1}{2} \left[I + \frac{1}{\sqrt{2}} (\sigma_x + \sigma_z) \right]$$

The black dots are the T -type states

$$|T\rangle\langle T| = \frac{1}{2} \left[I + \frac{1}{\sqrt{3}} (\sigma_x + \sigma_y + \sigma_z) \right]$$



Quantifying magic

- What about the continuous regime?
- As an extrapolation from the discrete theory, the negativity of the continuous variable Wigner function can be considered as a resource.
- Wigner Logarithmic Negativity (WLN) can be written as

$$WLN(\rho) = \log(\mathcal{N}(\rho) + 1), \quad \mathcal{N}(\rho) = \int d\mathbf{r} |W_\rho(\mathbf{r})| - 1$$

Phys. Rev. A 98, 052350 (2018)

Quantifying magic

Quantifying discrete magic

- The discrete Wigner function is not well defined in all dimensions, only in prime odd dimensions.
- One of the ways to surpass this issue is to take additional constrictions to the systems.
- We take the discrete Wigner function for SU(2) proposed in M. A. Marchioli and D. Galleti, J. Phys. A: Math. and Theo. 52 405305 (2019):

$$W(\mu, \nu) = \frac{1}{2} \text{tr} \left\{ \rho \left[\mathbf{1} + (-1)^\nu \sigma_x + (-1)^{\mu+\nu+1} \sigma_y + (-1)^\mu \sigma_z \right] \right\}$$
$$0 \leq \mu, \nu \leq 1$$

Quantifying magic

- ❑ Quantifying continuous magic
 - ❑ To propose measures of continuous magic besides Wigner logarithmic negativity.
 - ❑ In our case we want to quantify magic in Fock basis for bosons.
 - ❑ To extrapolate measures from the discrete regime to the continuous (as first approach we take just a methodology of discretization, not a theory).
 - ❑ To tackle the hybrid scenario to see the behavior of magic in subspaces of a bipartite system.

The Asymmetric Quantum Rabi Model

The Quantum Rabi Model

- ❑ The model describes the coherent interaction between one two-level atom and a quantized mode of light.
- ❑ It is one of the more simplest hybrid systems.
- ❑ Searching for hidden symmetries.
- ❑ It has promising applications, for instance, atom chip connected to a waveguide and superconducting qubits.

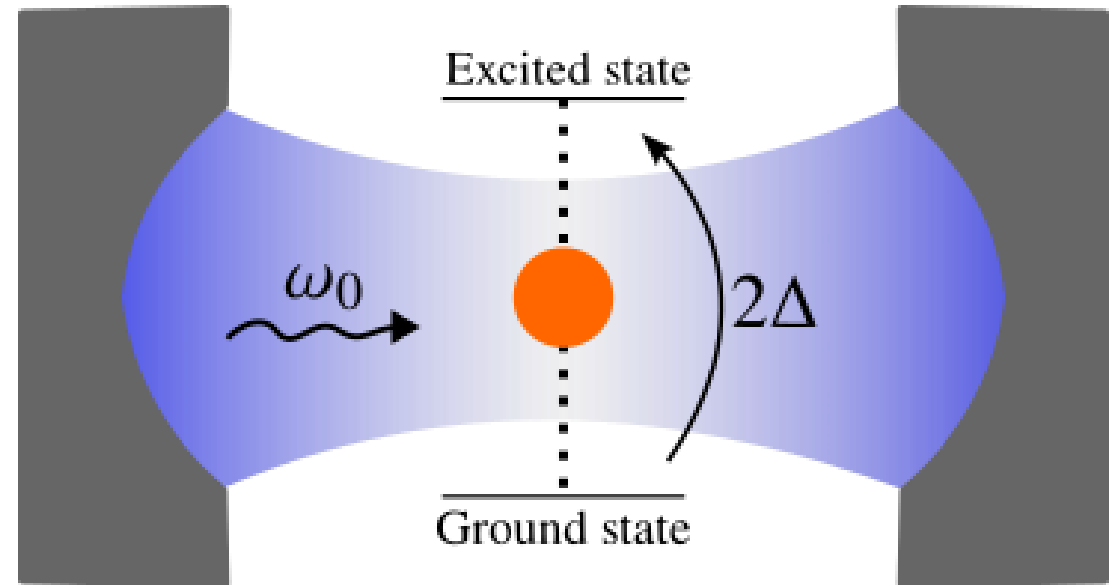


Figure: Scheme of the Quantum Rabi model

PRL 107 100401 (2011), Symmetry 11, 1259 (2019)

The Quantum Rabi Model

□ The **Quantum Rabi Model** describes a two-level system coupled to a quantum harmonic oscillator (mode of the electromagnetic field),

$$H_R = \Delta \sigma_z + \omega_0 a^\dagger a + g \sigma_x (a^\dagger + a)$$

2Δ - the level splitting between the two levels

ω_0 - the frequency of the quantized light field

g - the coupling strength between the light field and the two-level system ($\hbar = 1$).

PRL 107 100401 (2011), Symmetry 11, 1259 (2019)

The Asymmetric Quantum Rabi Model

□ The **Asymmetric Quantum Rabi Model** Hamiltonian,

$$H_A = \Delta\sigma_z + \epsilon\sigma_x + \omega_0 a^\dagger a + g\sigma_x (a^\dagger + a)$$

which has a bias extra term, proportional to ϵ , that is related with spontaneous tunnelling between the two levels.

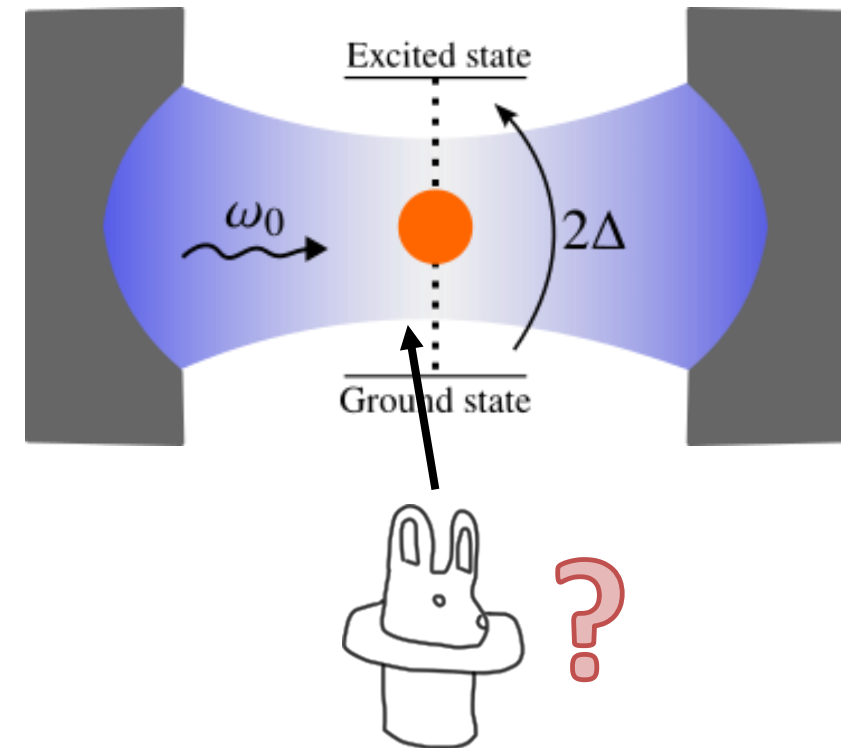
- $\epsilon\sigma_x$ - Breaks symmetry of the QRM — Restauration of parity symmetry when ϵ/ω_0 is half integer.
- To study the hidden symmetries of the AQRM could give us information about light matter interaction and its possible applications, for instance, superconducting qubits.

Phys. Rev. A 103 023719 (2021)

Magic + AQRM =
Magical AQRM

Quantifying Magic in the AQRM

1. We numerically obtain the energy spectrum of the complete system for various values of the constants of the model, Δ , ϵ , ω_0 and g .
2. As a first approach we partially trace the systems.
 - ☐ Tracing over harmonic oscillator.
 - ☐ Tracing over two-level system.
 - ☐ We quantify discrete magic for both subsystems.



Results

Considerations

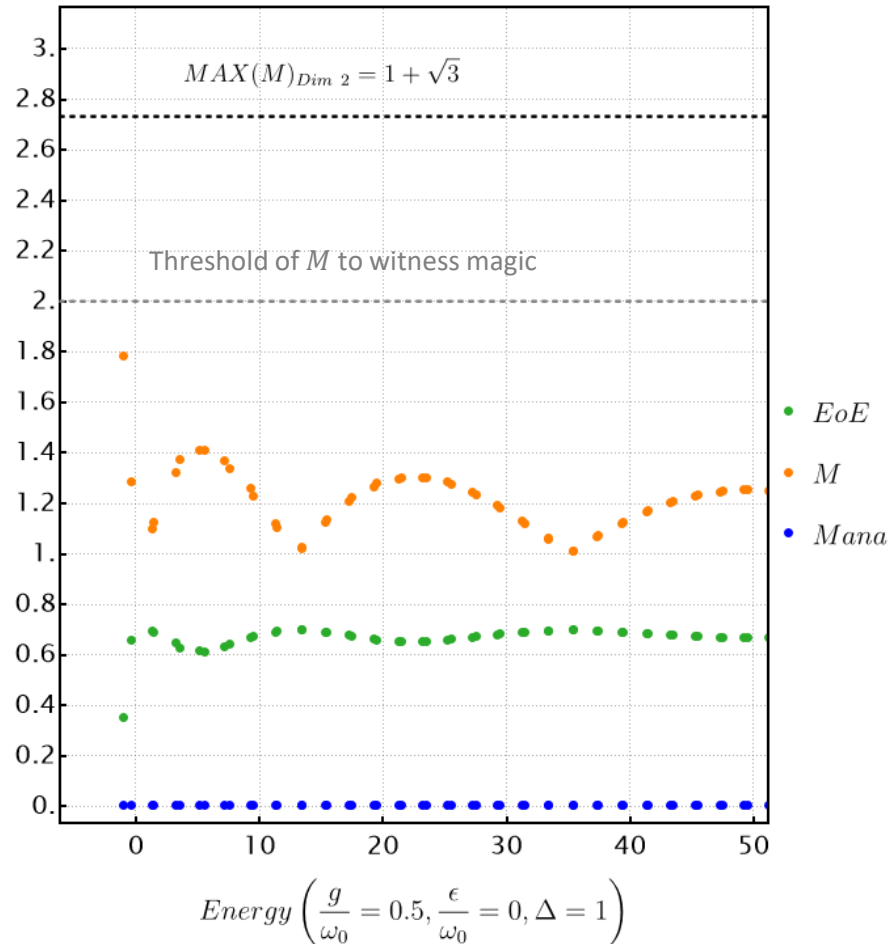
- To make our numerical computations we normalize the Hamiltonian with respect to ω_0 .

$$\frac{H_A}{\omega_0} = a^\dagger a + \frac{g}{\omega_0} (a^\dagger + a) \sigma_x + \frac{\Delta \sigma_z + \epsilon \sigma_x}{\omega_0}$$

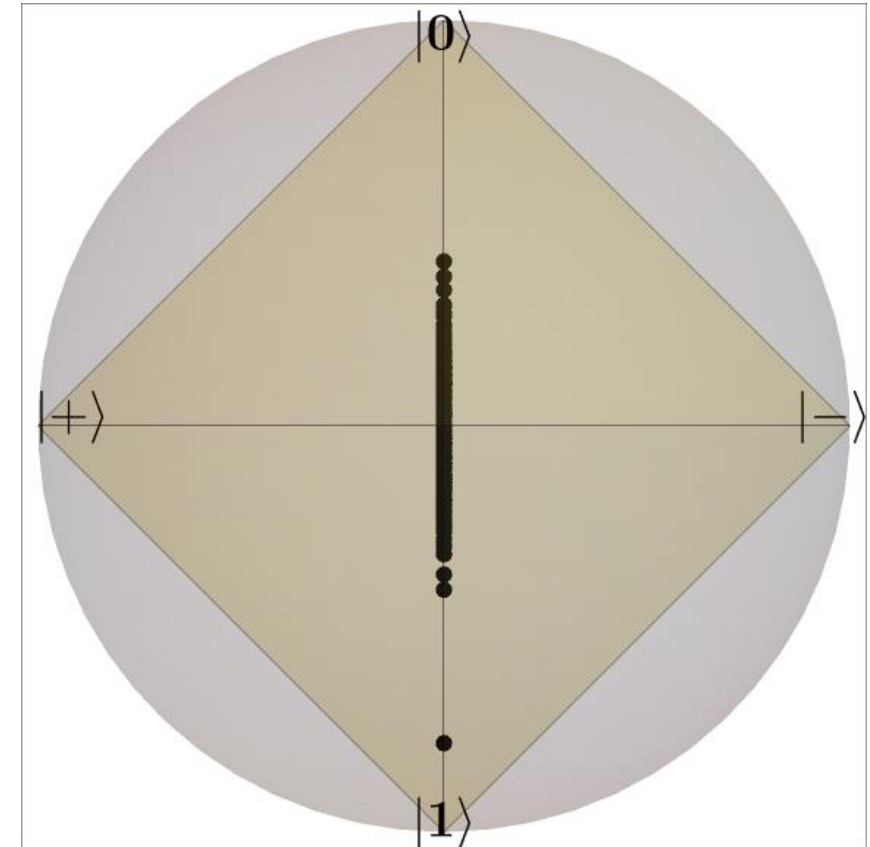
Then we diagonalize the Hamiltonian to obtain the eigen-energies and eigen vectors to construct the state of the system, $\rho_A(E_k)$. We trace such system to obtain the partially traced subsystems.

- $\rho_{light} = Tr_{atom} (\rho_A(E_k))$
- $\rho_{atom} = Tr_{light} (\rho_A(E_k))$

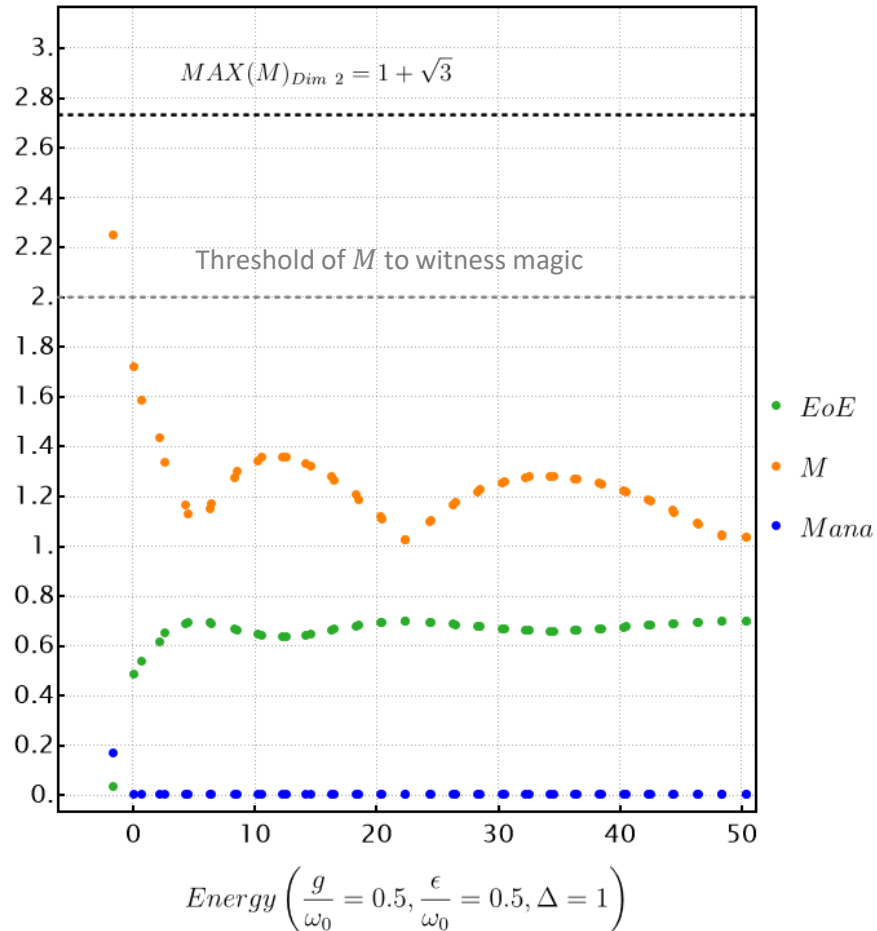
Magic in the atomic subspace



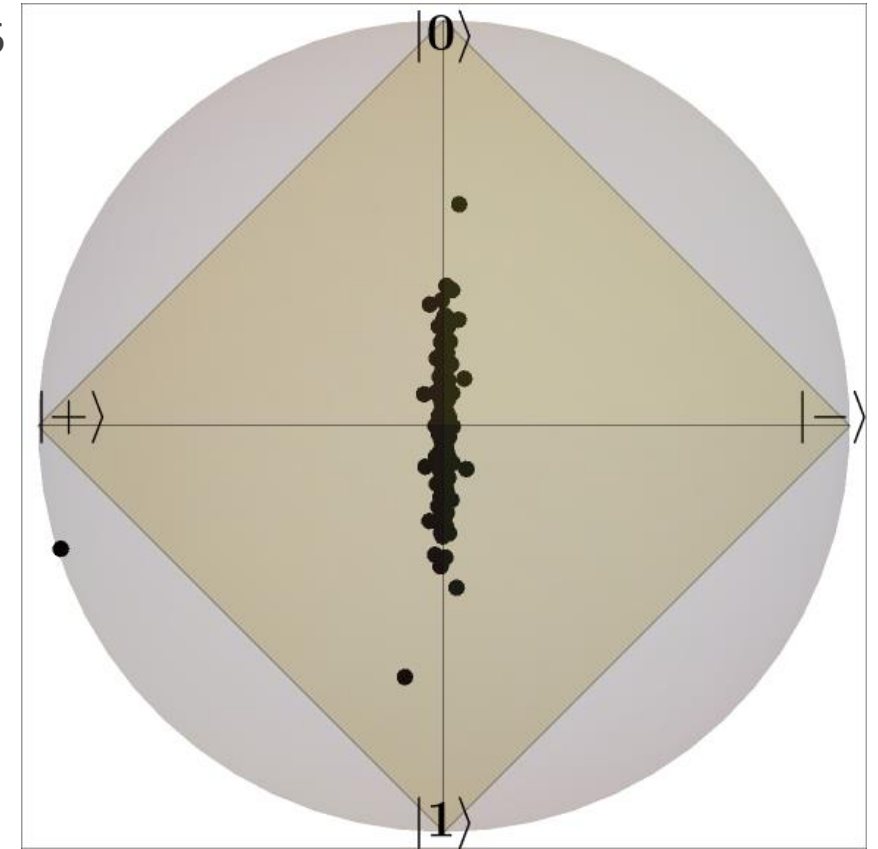
Fixed g/ω_0 , $\epsilon/\omega_0 = 0$



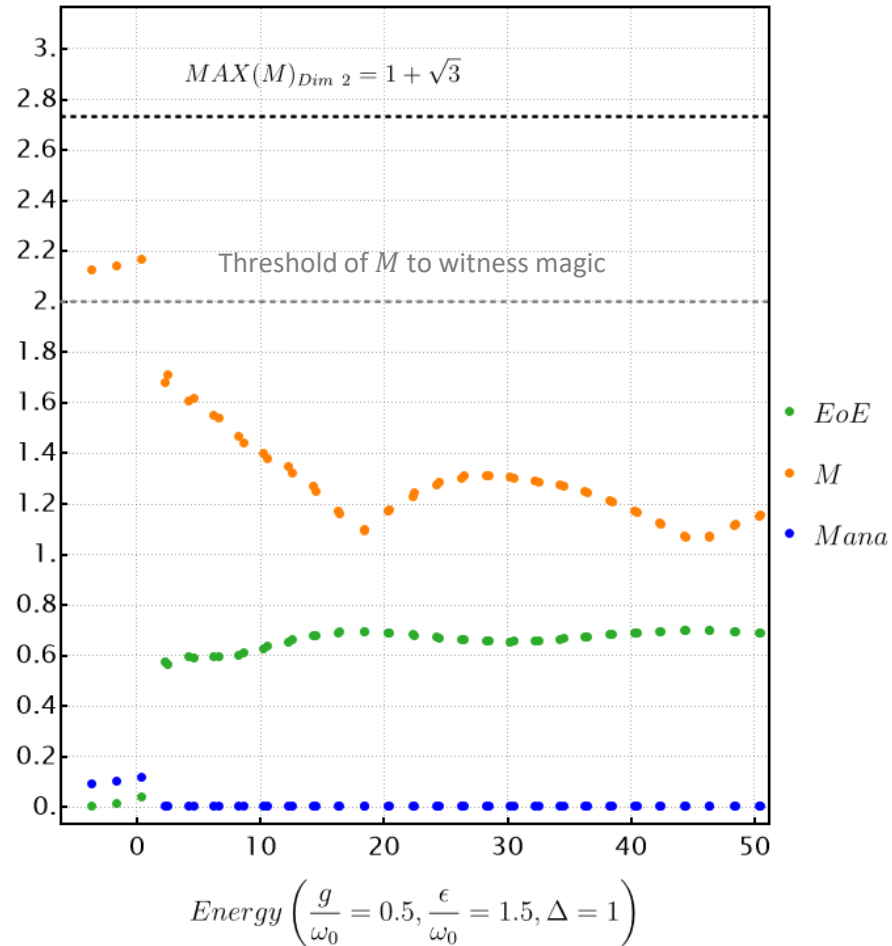
Magic in the atomic subspace



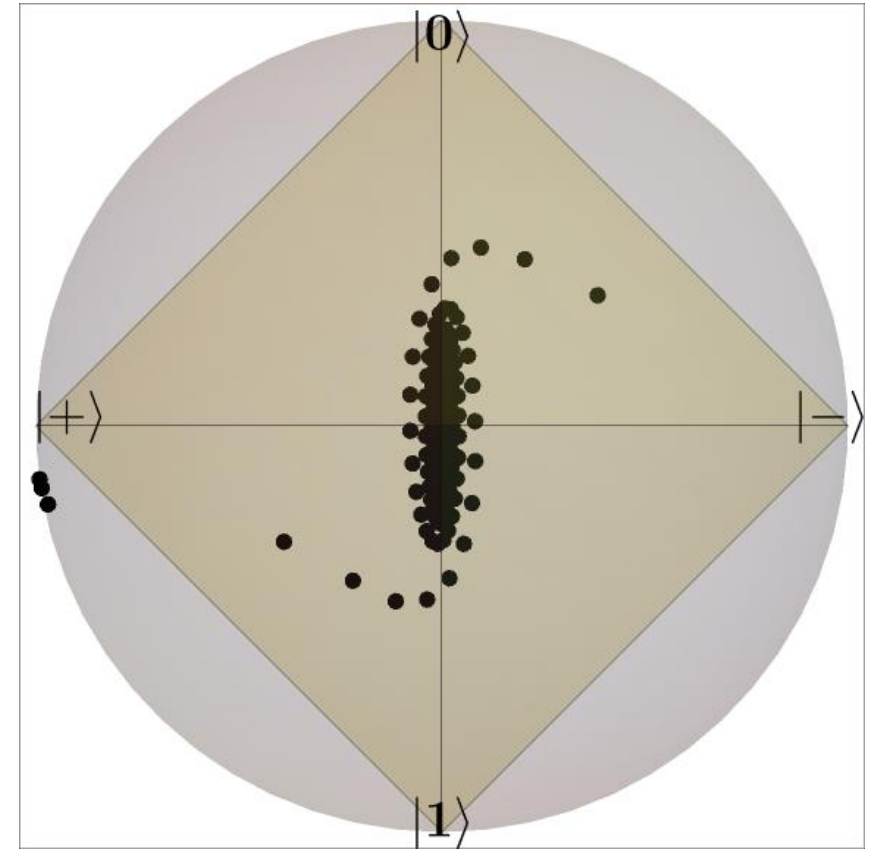
Fixed $g/\omega_0, \epsilon/\omega_0 = 0.5$



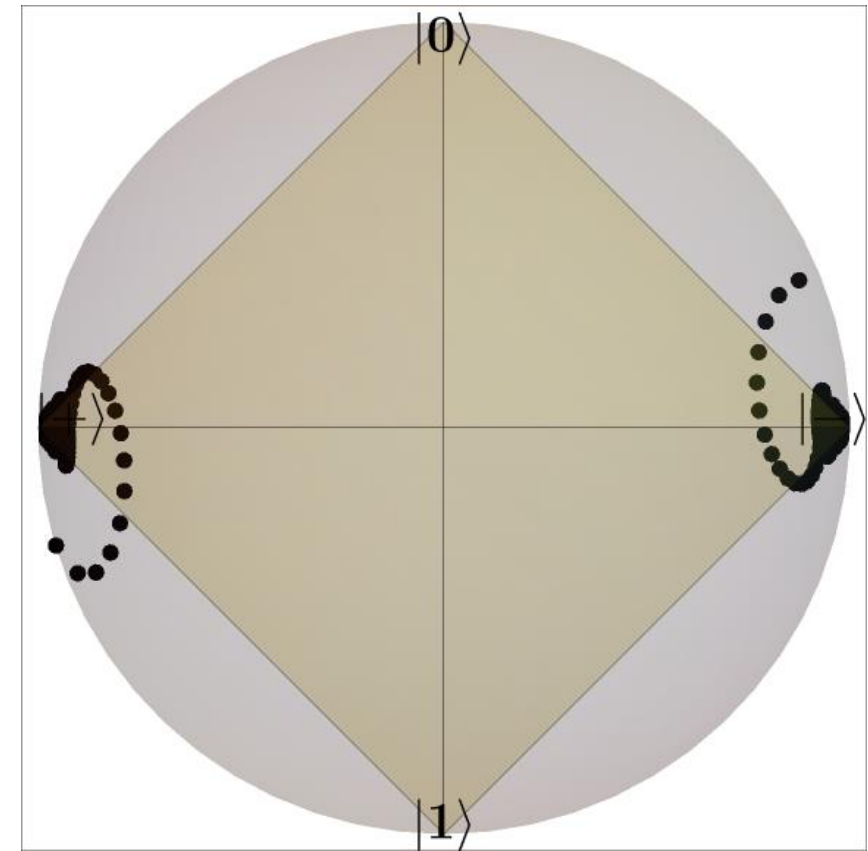
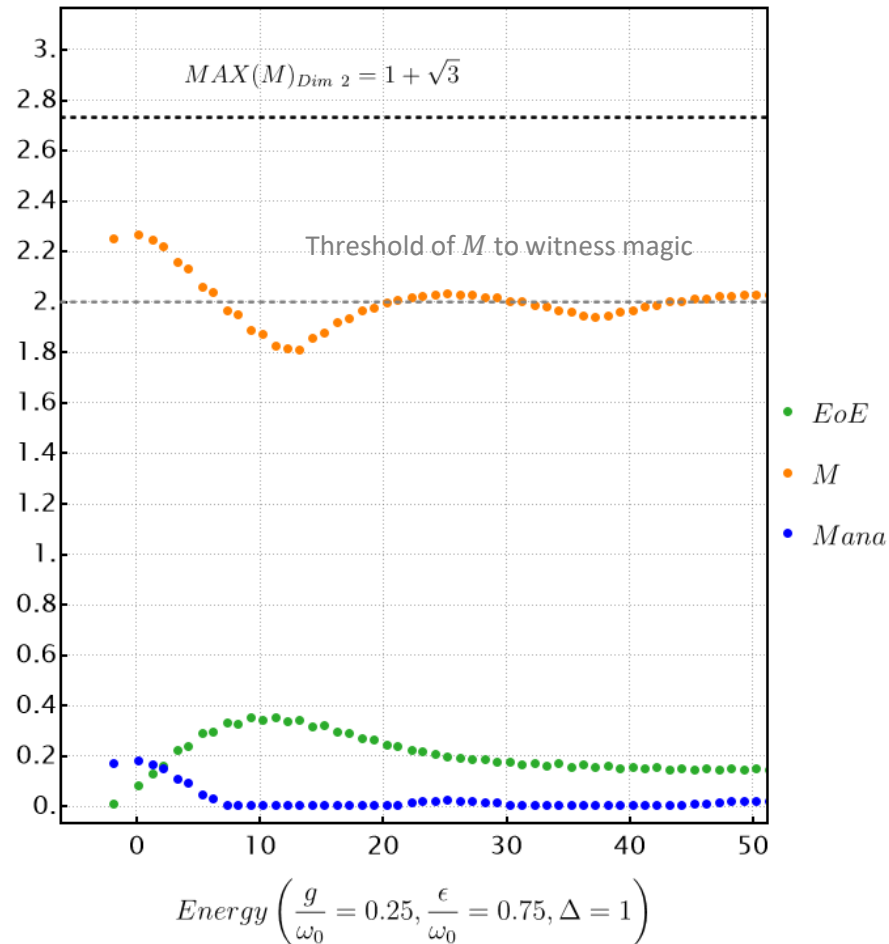
Magic in the atomic subspace



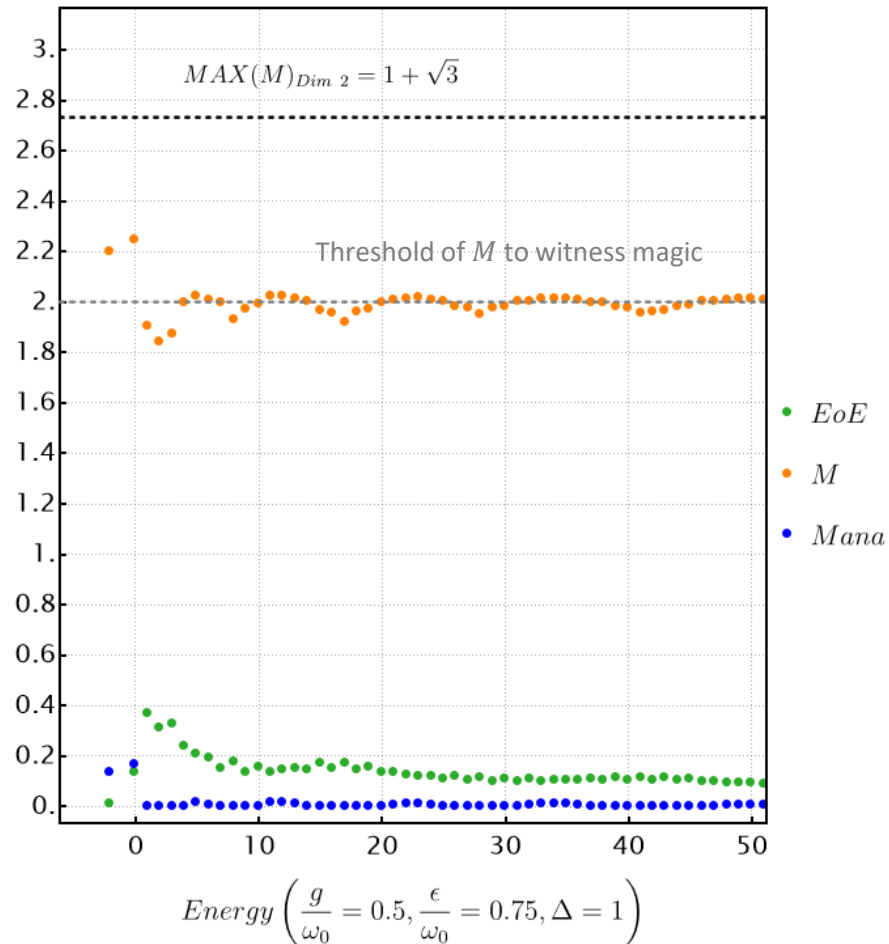
Fixed $g/\omega_0, \epsilon/\omega_0 = 1.5$



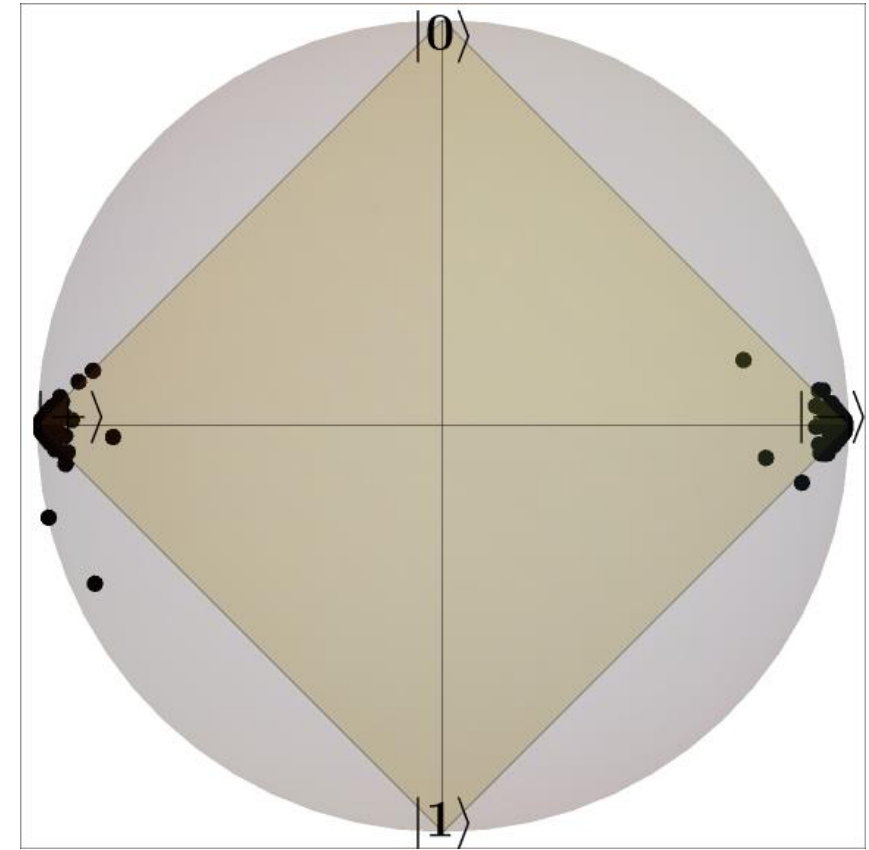
Magic in the atomic subspace



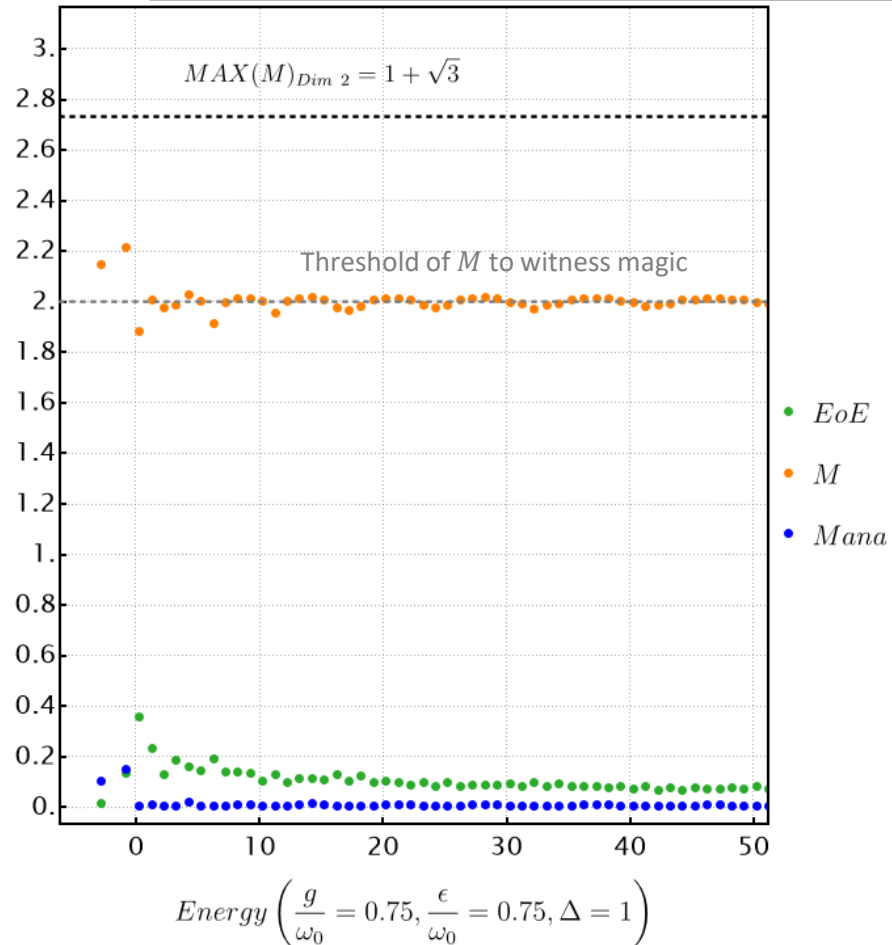
Magic in the atomic subspace



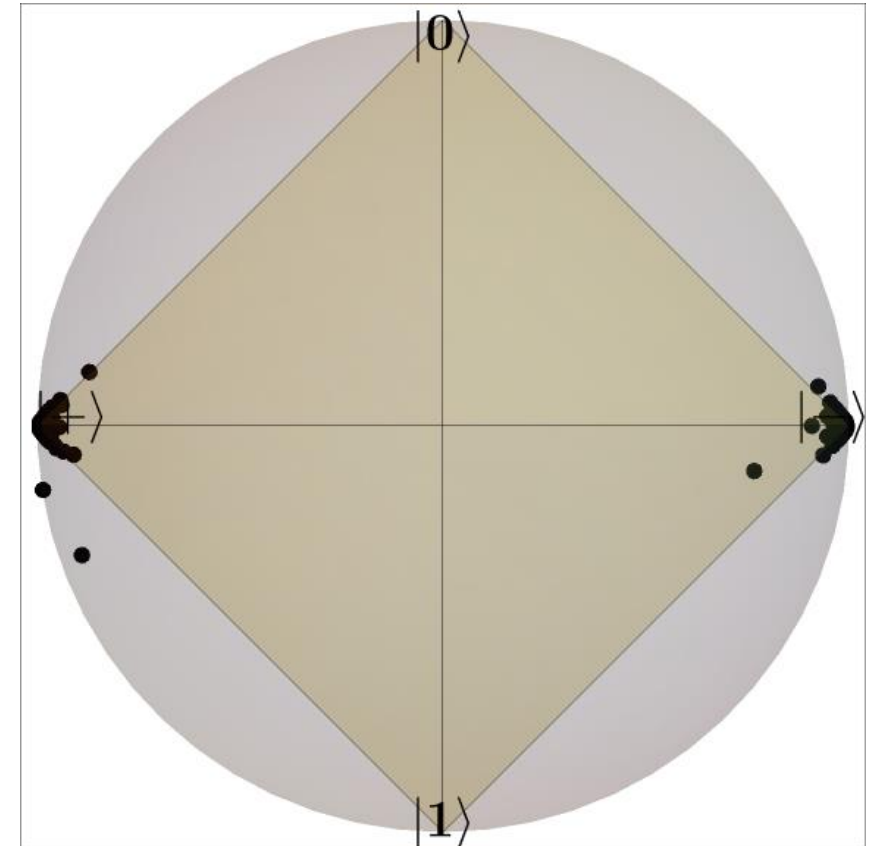
Fixed $\epsilon/\omega_0, g/\omega_0 = 0.5$



Magic in the atomic subspace



Fixed $\epsilon/\omega_0, g/\omega_0 = 0.75$



Magic in the bosonic subspace

□ Remember the Fock Wigner function,

$$W(\hat{\rho}, \alpha, \bar{\alpha}) = \text{tr} \left(\hat{\rho} \hat{W}(\alpha, \bar{\alpha}) \right), \quad \hat{W}(\alpha, \bar{\alpha}) = \hat{D}(\alpha) \hat{\Pi} \hat{D}^\dagger(\alpha)$$

where $\hat{\Pi} = \sum_{n=0}^{\infty} (-1)^n |n\rangle\langle n|$ is the parity operator and

$\hat{D}(\alpha) = \exp(\alpha \hat{a}^\dagger - \bar{\alpha} \hat{a})$ the Fock displacement operator. And

$$\hat{a} = \sqrt{\frac{m\omega}{2\hbar}} \left(\hat{q} + \frac{i\hat{p}}{m\omega} \right), \quad \alpha(q, p) = \sqrt{\frac{m\omega}{2\hbar}} \left(q + \frac{ip}{m\omega} \right)$$

Fortschr. Phys. 65 1600092 (2017)

Magic in the bosonic subspace

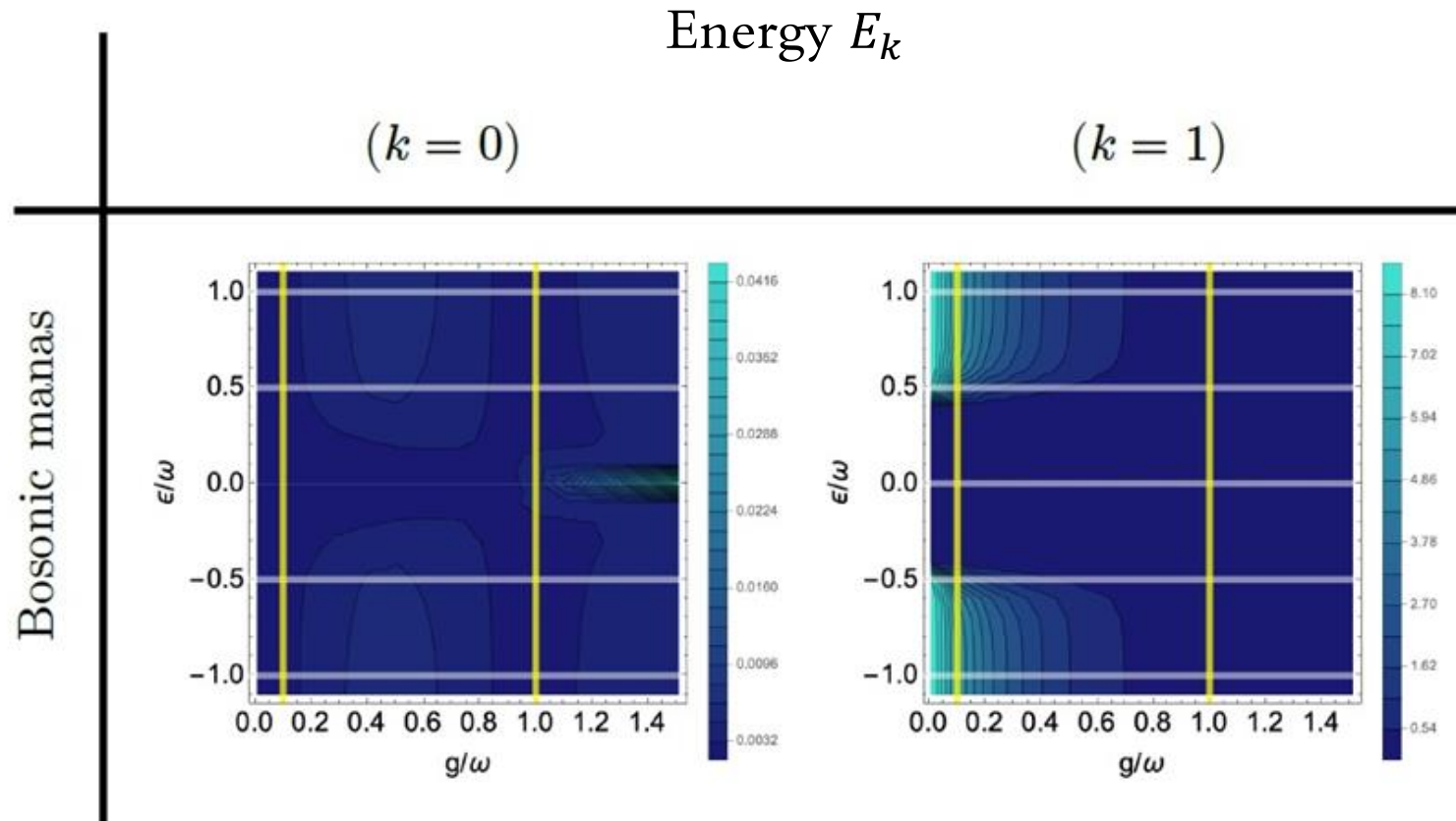
To numerically quantify magic in bosonic space we made a grid of the phase space of the traced subsystem,

$$\text{Grid} = \left\{ (\alpha_R, \alpha_I) : \alpha_R = [-\ell, \ell], \alpha_I = [-\ell, \ell] \right\}$$

And, we obtain plots for a given energy E_k ,

$$\left\{ (\alpha_R, \alpha_I, W[E_k, \alpha, \bar{\alpha}]) : \alpha \in \text{Grid} \right\}$$

Magic in the bosonic subspace

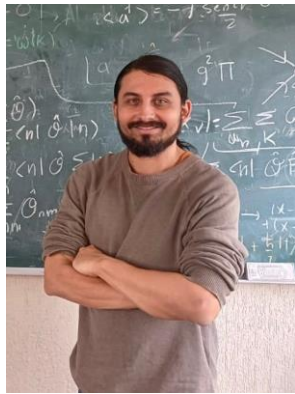


Conclusions

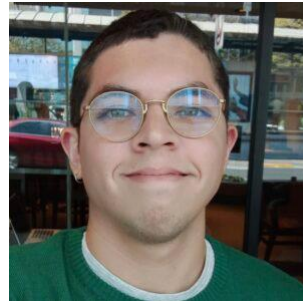
- We characterize magic in the AQRM subsystems, namely, atomic and bosonic.
- We compare the behaviour of a variety of measures of magic in each subsystem.
- We also check the behaviour of magic, comparing mana for fixed parameters in both subsystems.
- Our endeavours pave the way towards a hybrid measure of magic in AQRM.

Acknowledgments

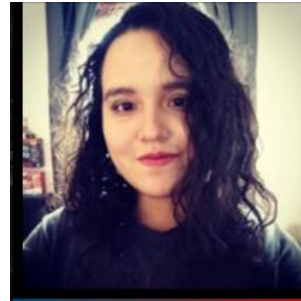
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Thank you!

