Perturbative Computations in Quantum Field Theory

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Collaborative Research Center TRR 257

CPPS retreat Meinerzhagen, February 14th - 16th, 2024

Equations of motion in classical mechanics

- Example: Bead on rotating ring
- Lagrangian function

$$L = \frac{m}{2} \left[R^2 \dot{\theta}^2 + \omega^2 R^2 \sin^2 \theta \right] + m g R \cos \theta$$

• Equation of motion (Euler-Lagrange equation):

$$\frac{d}{dt}\frac{\partial L}{\partial \dot{\theta}} - \frac{\partial L}{\partial \theta} = 0$$

Yields:

$$\ddot{\theta} = \omega^2 \sin \theta \cos \theta - \frac{g}{R} \sin \theta$$



Equations of motion in field theory

- Example: QED
- Lagrangian density

$$\mathcal{L} = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu} + \bar{\psi}(i\partial \!\!\!/ - m)\psi + e \,\bar{\psi} \mathcal{A}\psi$$

- ψ : Fermion field (electron, positron)
- A^µ: Photon field
- $F^{\mu\nu} = \partial^{\mu}A^{\nu} \partial^{\nu}A^{\mu}$: Field strength tensor
- $e = \sqrt{4\pi\alpha}$: coupling. $\alpha \sim \frac{1}{137}$

• Euler-Lagrange equations:

$$\partial_{\mu} \, \frac{\partial \mathcal{L}}{\partial (\partial_{\mu} \bar{\psi})} - \frac{\partial \mathcal{L}}{\partial \bar{\psi}} = 0$$

$$\partial_{\mu} \frac{\partial \mathcal{L}}{\partial (\partial_{\mu} A_{\nu})} - \frac{\partial \mathcal{L}}{\partial A_{\nu}} = 0$$

Yields:

$$(i\partial \!\!\!/ - m)\psi = -e \, A\!\!\!/ \psi$$
$$\partial_{\mu}F^{\mu\nu} = -e \, \bar{\psi}\gamma^{\nu}\psi$$

Coupled, nonlinear equations
Cannot be solved in closed form

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Perturbative expansion

 Idea: Series expansion in small coupling

$$\alpha \sim \frac{1}{137} \qquad \alpha_s \sim 0.1 \ - \ 0.2$$

$$\mathcal{A} = \mathcal{O}(1) + \mathcal{O}(\alpha) + \mathcal{O}(\alpha^2) + \dots$$

- At each order, more Feynman diagrams with more loops
- For each loop, have one unconstrained four-momentum *ℓ*^μ, to be integrated over with measure

 $\int \frac{d^4\ell}{(2\pi)^4}$

• $a_e = (g_e - 2)/2$ of electron



 $a_e = 0.00115965218059(13)$ (experiment)

IBP reduction

- A multi-loop amplitude contains easily millions of terms (integrals)
- Not all integrals independents
- Integration-by-parts (IBP) equations

[Tkachov'81;Chetyrkin,Tkachov'81]

$$\int \frac{d^d \ell_1}{i\pi^{d/2}} \cdots \int \frac{d^d \ell_L}{i\pi^{d/2}} \frac{\partial}{\partial \ell_i^{\mu}} \frac{v_j^{\mu}}{D_1^{a_1} \cdots D_n^{a_n}} = 0$$

For the one-loop bubble

$$\frac{v=p:}{0=(a_1-a_2)F(a_1,a_2)+a_2p^2F(a_1,a_2+1)-a_1p^2F(a_1+1,a_2)}{a_2F(a_1-1,a_2+1)-a_1F(a_1+1,a_2-1)}$$



Laporta's algorithm

[Laporta'01]

- Specialize set of IBP equations to a range of integer indices
- Generate system of linear equations

0 = F(2,1) - F(1,2) $0 = (d-3)F(1,1) - p^{2}F(1,2)$ $0 = (d-4)F(1,2) - 2p^{2}F(1,3)$ $0 = 2p^{2}F(1,3) - p^{2}F(2,2) - F(1,2) - F(2,1)$ $0 = (d-5)F(2,1) - p^{2}F(2,2) - F(1,2)$ \vdots include symmetries directly discard scaleless integrals $0 = (d-5)F(2,1) - p^{2}F(2,2) - F(1,2)$

Solve system by Gaussian elimination.

Express each integral as a linear combination of independent master integrals.

IBP reduction: Laporta's algorithm

- Calcuation of higher loop-orders entails several challenges
 - number of Feynman diagrams
 - size of intermediate expressions
 - complexity of loop integrals (# loops, legs, scales)
- Many sophisticated public and private tools exist

| ۲ | AIR | [Anastasiou,Lazopoulos'04] |
|---|---------|--|
| ۰ | Reduze | [Studerus'09; Manteuffel,Studerus'12] |
| ۲ | FIRE | [Smirnov et al.'08+] |
| ۲ | LiteRed | [Lee'12+] |
| ۲ | Kira | [Maierhöfer,Usovitsch,Uwer'17; Klappert,Lange,Maierhöfer,Usovitsch'20] |
| ٩ | Crusher | [Marquard,Seidel] |

- Mostly based on Laporta's algorithm [Laporta'01]
 - Solves IBP equations for numerical values of indices with Gaussian elimination.
- Several refinements exist, e.g.
 - Parallelization
 - Methods from finite fields

[v. Manteuffel.Schabinger'14: Smirnov.Chukharev'19] [Peraro'16'19: Klappert.Klein.Lange'19'20]

Drawbacks

- Compute many more integrals than reauired
- Large storage required for results of $10^{\sim 4-6}$ integrals

- New ideas from
 - algebraic geometry

[Larsen,Zhang'14; Böhm et al.'18'19]

syzygy equations

[Kosower et al.'10'18; Schabinger et al.'11'20; Ita'15; Böhm et al.'17]

$$A \cdot B = 0$$

- Linear algebra
- New approach

[Barakat,TH et al. 2020+]

- Leave propagator powers symbolic
- Derive so-called normal-form IBP relations

Normal-form IBP equations

$$R_i = a_i D_i^- - NF_G \left(a_i D_i^- \right) \in I_{\text{IBP}}$$

• e.g. for one-loop bubble

$$NF_G(a_1D_1^-) = \frac{(d-a_1-a_2-1)(d-2a_1-2a_2)}{p^2(d-2a_1-2)}$$

NF_G
$$(a_2 D_2^-) = \frac{(d - a_1 - a_2 - 1)(d - 2a_1 - 2a_2)}{p^2(d - 2a_2 - 2)}$$

- Allow for a straightforward reduction
- All you have to store
- Easy to implement in Mathematica or FORM
- Well-suited for parallelization
- Allows for fast reduction, also of high propagator powers

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Laporta vs. Normal-form IBPs

Compare reduction with Laporta vs. Normal-form IBPs for the one-loop massless bubble

Reduction of F(3,2):



- Jobs run, depending on complexity, on Laptop, desktop, OMNI, cluster in Karlsruhe
- Need modern CAS (Mathematica, C++, Fermat, ...)
- RAM, up to $\mathcal{O}(10^{2-3} \text{GB})$
- storage space up to $\mathcal{O}(1TB)$
- Fast I/O traffic, read/write on local disk (scratch)
- Parallelizable ($\mathcal{O}(20)$ cores)
- Job runtimes of up to $\mathcal{O}(\text{weeks})$
- Training for students in CAS, programming needed
- Local resources improvable in machines with a lot of RAM, Mathematica licenses, job length
- Not so much in need of GPUs, network