

Rare $D^0 \rightarrow \pi^+ \pi^- \ell^+ \ell^-$ decays:
Contribution of S-wave dynamics in the SM
and sensitivity to New Physics

based on 2308.XXXXX

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Introduction

Flavor Physics with rare processes

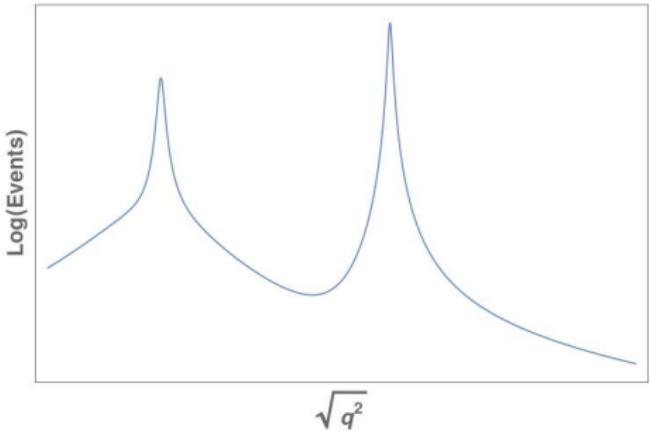
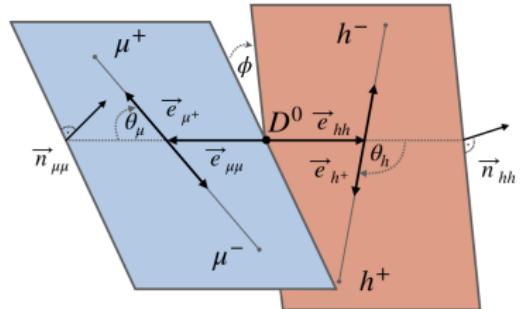
- Rare decays could provide the ground for discovering New Physics effects:
SM at loop level
- Rich history: $K \rightarrow \mu\mu$, B Physics
- Many NP tests can be constructed
- Some tests theoretically very clean (e.g. LFU tests)
- Others sensitive to hadronic effects - require careful QCD implementation
(e.g. in P'_5)
- Rare decays with charm provide complementary tests from the up-type sector

See e.g. Burdman et al., Fajfer et al., Bharucha et al.,...

Rare processes with charm

New LHCb results on $D^0 \rightarrow \pi^+ \pi^- \mu^+ \mu^-$ and $D^0 \rightarrow K^+ K^- \mu^+ \mu^-$ [PRL 128(22):221801]

- 5D phase space: q^2 (leptons), p^2 (pions), $\theta_h, \theta_\ell, \phi$
⇒ 5 1D plots of differential decay rates (LHCb data for q^2, p^2)



- Angular observables: different integrations
 $\sum_i c_i \int_{\theta_{\ell,i}}^{\theta_{\ell,i+1}} d\cos\theta_\ell \sum_j c'_j \int_{\phi_i}^{\phi_{i+1}} d\phi d^5\Gamma$ (e.g. A_{FB} is a null test!)
- CP-asymmetric ones: further tests (not the focus of the talk)

Rare processes with charm

New LHCb results on $D^0 \rightarrow \pi^+ \pi^- \mu^+ \mu^-$ and $D^0 \rightarrow K^+ K^- \mu^+ \mu^-$ [PRL 128(22):221801]

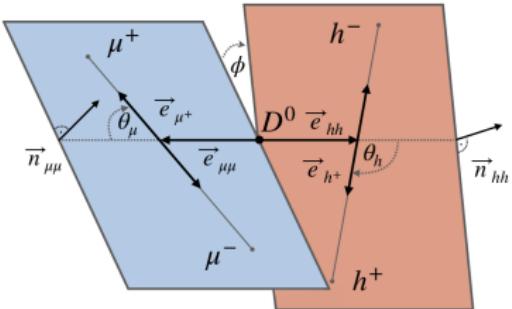
- 5D phase space: q^2 (leptons), p^2 (pions), $\theta_h, \theta_\ell, \phi$
 \Rightarrow 5 1D plots of differential decay rates (LHCb data for q^2, p^2)

Table 2: Angular observables ($\langle S_i \rangle$) for (top) $D^0 \rightarrow \pi^+ \pi^- \mu^+ \mu^-$ and (bottom) $D^0 \rightarrow K^+ K^- \mu^+ \mu^-$ decays in the dimuon-mass regions. The first uncertainty is statistical and the second systematic.

$m(\mu^+ \mu^-)$ [MeV/ c^2]	$\langle S_0 \rangle$ [%]	$\langle S_1 \rangle$ [%]	$\langle S_2 \rangle$ [%]	$\langle S_3 \rangle$ [%]	$\langle S_4 \rangle$ [%]	$\langle S_5 \rangle$ [%]	$\langle S_6 \rangle$ [%]	$\langle S_7 \rangle$ [%]	$\langle S_8 \rangle$ [%]
< 525	5 ± 14 ± 4	-6 ± 16 ± 2	21 ± 16 ± 2	-20 ± 14 ± 1	-14 ± 14 ± 1	8 ± 14 ± 1	16 ± 17 ± 1	26 ± 16 ± 2	-
525–565	-	-	-	-	-	-	-	-	-
565–600	-2.1 ± 4.1 ± 1.1	-0.1 ± 1.8 ± 1.5	3.7 ± 0.9 ± 1.3	-0.1 ± 4.1 ± 0.6	2.5 ± 1.0 ± 0.6	0.8 ± 1.1 ± 1.0	12.9 ± 0.6 ± 1.0	-0.1 ± 4.9 ± 0.9	-
600–650	-10.7 ± 5.8 ± 1.1	7.7 ± 4.1 ± 1.0	-4.7 ± 5.8 ± 1.1	4.7 ± 5.8 ± 0.7	9.6 ± 5.8 ± 1.0	1.4 ± 6.0 ± 0.7	-1.7 ± 6.0 ± 0.8	-	-
650–1020	-2.0 ± 3.7 ± 1.6	-37.0 ± 4.3 ± 1.1	-2.0 ± 3.7 ± 1.5	2.0 ± 3.7 ± 0.8	6.5 ± 3.7 ± 1.4	-3.6 ± 3.7 ± 1.1	2.6 ± 4.3 ± 0.9	10.6 ± 4.3 ± 1.0	-
1020–1100	1.7 ± 3.4 ± 1.5	-35.3 ± 4.0 ± 1.7	-8.3 ± 4.0 ± 2.5	-6.9 ± 3.4 ± 1.2	1.1 ± 3.4 ± 0.8	2.7 ± 3.4 ± 2.0	0.7 ± 4.1 ± 0.9	7.8 ± 4.0 ± 1.7	-
> 1100	-	-	-	-	-	-	-	-	-
Full range	-3.4 ± 2.1 ± 1.0	-10.4 ± 2.5 ± 0.9	-4.6 ± 2.5 ± 1.6	-2.9 ± 2.1 ± 0.6	3.7 ± 2.1 ± 0.5	-0.6 ± 2.1 ± 0.9	3.8 ± 2.5 ± 0.5	5.1 ± 2.5 ± 0.5	-

Table 3: Angular observables ($\langle A_i \rangle$) for (top) $D^0 \rightarrow \pi^+ \pi^- \mu^+ \mu^-$ and (bottom) $D^0 \rightarrow K^+ K^- \mu^+ \mu^-$ decays in the dimuon-mass regions. The first uncertainty is statistical and the second systematic.

$m(\mu^+ \mu^-)$ [MeV/ c^2]	$\langle A_0 \rangle$ [%]	$\langle A_1 \rangle$ [%]	$\langle A_2 \rangle$ [%]	$\langle A_3 \rangle$ [%]	$\langle A_4 \rangle$ [%]	$\langle A_5 \rangle$ [%]	$\langle A_6 \rangle$ [%]	$\langle A_7 \rangle$ [%]	$\langle A_8 \rangle$ [%]	$\langle A_9 \rangle$ [%]
< 525	-10 ± 14 ± 2	2 ± 16 ± 1	-7 ± 16 ± 2	16 ± 14 ± 1	0 ± 14 ± 1	-10 ± 14 ± 2	3 ± 17 ± 2	-25 ± 16 ± 2	-	-
525–565	-	-	-	-	-	-	-	-	-	-
565–780	-1.1 ± 4.1 ± 1.9	3.7 ± 4.8 ± 0.7	0.6 ± 4.9 ± 0.7	-0.0 ± 4.1 ± 1.1	-4.8 ± 4.1 ± 1.0	-3.5 ± 4.1 ± 1.0	-1.8 ± 4.9 ± 1.2	1.6 ± 4.9 ± 1.1	-	-
780–850	-7.7 ± 5.8 ± 0.6	3.9 ± 6.0 ± 0.8	1.2 ± 6.0 ± 0.7	-3.3 ± 5.8 ± 1.0	0.4 ± 5.8 ± 0.6	-2.6 ± 5.8 ± 0.6	-5.1 ± 6.9 ± 1.5	-2.9 ± 6.8 ± 1.0	-	-
850–1020	2.3 ± 3.7 ± 0.7	-2.6 ± 4.3 ± 2.1	7.6 ± 4.3 ± 0.9	-3.6 ± 3.7 ± 1.2	4.5 ± 3.7 ± 1.1	3.5 ± 3.7 ± 0.9	2.7 ± 4.3 ± 1.3	1.4 ± 4.3 ± 1.2	-	-
1020–1100	-4.8 ± 3.4 ± 0.9	-2.6 ± 4.0 ± 1.2	-2.4 ± 4.0 ± 1.0	-2.3 ± 3.4 ± 1.2	3.2 ± 3.4 ± 1.1	-1.3 ± 3.4 ± 0.8	5.1 ± 4.3 ± -1.3	-5.9 ± 4.0 ± 1.8	-	-
> 1100	-	-	-	-	-	-	-	-	-	-
Full range	-2.6 ± 2.1 ± 0.7	0.9 ± 2.5 ± 0.5	1.3 ± 2.5 ± 0.5	-1.7 ± 2.1 ± 0.9	-0.5 ± 2.1 ± 0.9	-1.5 ± 2.1 ± 0.6	0.3 ± 2.5 ± 1.0	-2.1 ± 2.5 ± 0.9	-	-

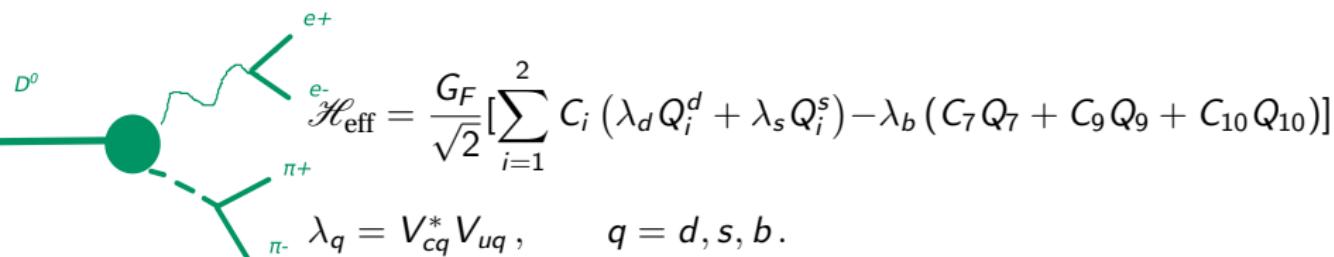


- Angular observables: different integrations

$$\sum_i c_i \int_{\theta_{\ell,i}}^{\theta_{\ell,i+1}} d \cos \theta_\ell \sum_j c'_j \int_{\phi_i}^{\phi_{i+1}} d\phi d^5 \Gamma \quad (\text{e.g. } A_{FB} \text{ is a null test!})$$

- CP-asymmetric ones: further tests (not the focus of the talk)

What is interesting about rare charm decays?



$$Q_1^d = (\bar{d}c)_V - A(\bar{u}d)_V - A,$$

$$Q_2^d = (\bar{u}c)_V - A(\bar{d}d)_V - A,$$

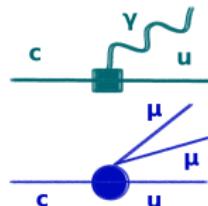
$$Q_1^s = (\bar{s}c)_V - A(\bar{u}s)_V - A,$$

$$Q_2^s = (\bar{u}c)_V - A(\bar{s}s)_V - A,$$

$$Q_7 = -\frac{g_{em}}{8\pi^2} m_c \bar{u} \sigma^{\mu\nu} (1 + \gamma_5) F^{\mu\nu} c,$$

$$Q_9 = \frac{\alpha_{em}}{2\pi} (\bar{u} \gamma_\mu (1 - \gamma_5) c) (\bar{\ell} \gamma^\mu \ell),$$

$$Q_{10} = \frac{\alpha_{em}}{2\pi} (\bar{u} \gamma_\mu (1 - \gamma_5) c) (\bar{\ell} \gamma^\mu \gamma_5 \ell),$$



- Strong GIM suppression $\Rightarrow C_{9,10,7}$ very small! ($\neq B$ Physics: $|C_9|, |C_{10}| \approx 4.2$)

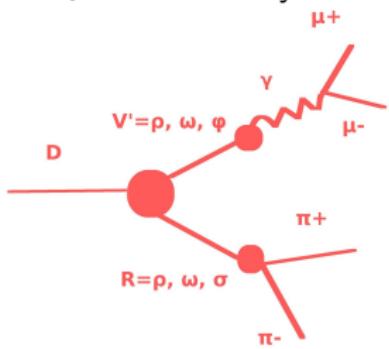
- Result:

- (+) Some observables vanish in the SM ("null tests") [de Boer, Hiller Phys.Rev.D 98 (2018) 3]
- (-) SM contributions are only long-distance/non-perturbative QCD \Rightarrow difficult territory to navigate!
- Typical picture: interplay between SM and potential NP contributions
- To probe NP Wilson coefficients of e.g. SMEFT, we need good control of SM contributions

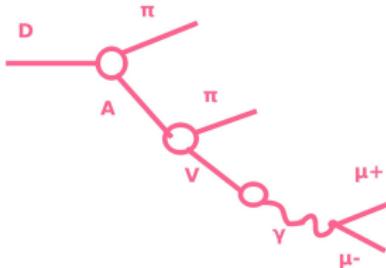
Framework

Rare D decays: model

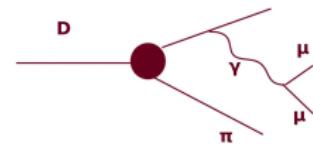
Quasi-two-body



Cascade

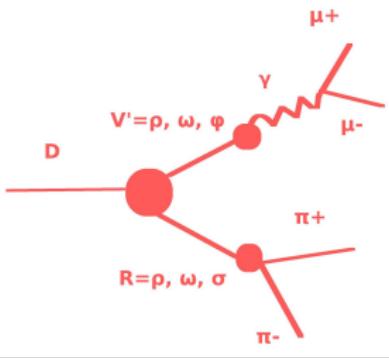


Bremmstrahlung



Rare D decays: model

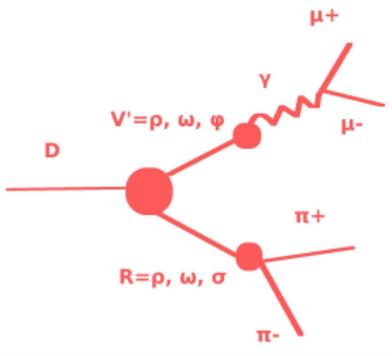
Quasi-two-body



$$\langle \pi^+ \pi^- \ell^+ \ell^- | S | D^0 \rangle = \langle \pi^+ \pi^- \ell^+ \ell^- | \int d^4x d^4w d^4y d^4z \sum_{\mathcal{R}, \mathcal{V}} T \{ \mathcal{H}_{em}^{\text{lept}}(z) \mathcal{H}_{V\gamma}(y) \mathcal{H}_{R\pi\pi}(w) \mathcal{H}_{eff}(x) \} | D^0 \rangle$$

Rare D decays: model

Quasi-two-body



$$\langle \pi^+ \pi^- \ell^+ \ell^- | S | D^0 \rangle = \langle \pi^+ \pi^- \ell^+ \ell^- | \int d^4x d^4w d^4y d^4z \sum_{\mathcal{R}, \mathcal{V}} T \{ \mathcal{H}_{em}^{\text{lept}}(z) \mathcal{H}_{V\gamma}(y) \mathcal{H}_{R\pi\pi}(w) \mathcal{H}_{eff}(x) \} | D^0 \rangle$$

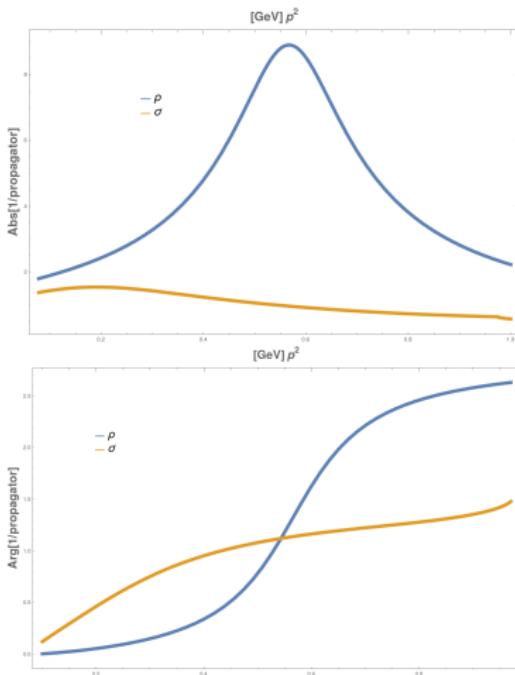
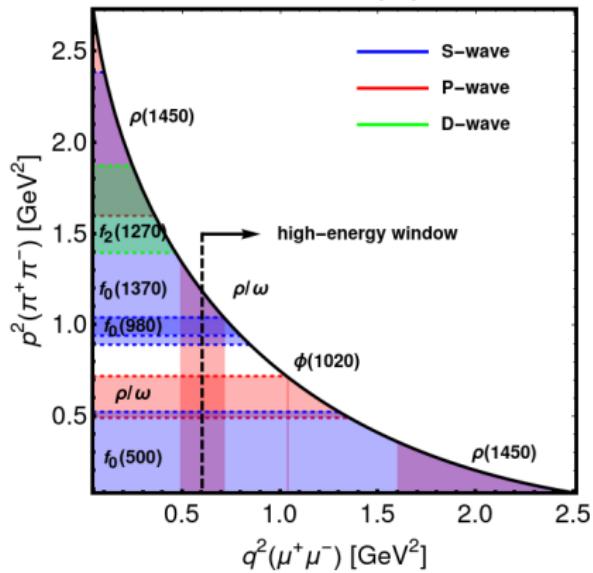
New: we include the scalar resonance $f_0(500) = \sigma$ in $D^0 \rightarrow \sigma \rightarrow \pi^+ \pi^-$

(observed in $D \rightarrow \pi\pi e\nu_e$ [BESIII], evidence in $D \rightarrow 4\pi$ [d'Argent+, CLEO data])

Instead of Breit-Wigner, Bugg's parameterization [Bugg, Phys.Lett.B 572 (2003) 1-7]: includes rescattering effects

Resonances

$D^0 \rightarrow \pi^+ \pi^- \mu^+ \mu^-$



Factorisation-inspired

Starting point: Cappiello, Cata, d'Ambrosio, JHEP 04 (2013) 135

$$Q_2^d = (\bar{u}c)_{V-A}(\bar{d}d)_{V-A}$$

$$\rho, \omega, \sigma : \{u, d\}, \quad \phi : s \quad \text{Zweig rule}$$

$$Q_2^s = (\bar{u}c)_{V-A}(\bar{s}s)_{V-A}$$

Tree -type 1 ('W')

$$D \rightarrow R \rightarrow \pi\pi$$

$$0 \rightarrow V \rightarrow \ell\ell$$

Tree -type 2 ('J')

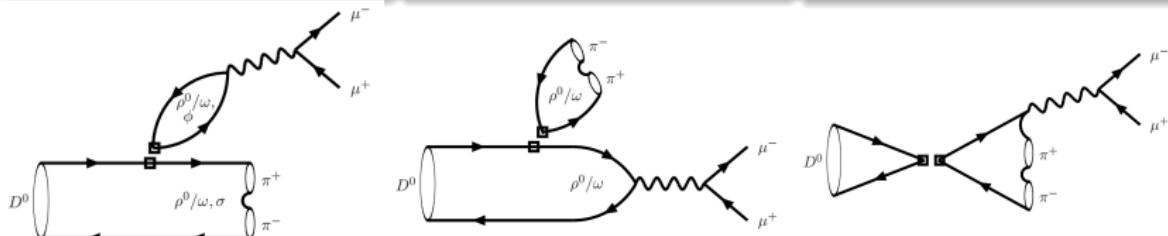
$$D \rightarrow V \rightarrow \ell\ell$$

$$0 \rightarrow R \rightarrow \pi\pi$$

Annihilation

$$D \rightarrow 0$$

$$0 \rightarrow R(\rightarrow \pi\pi)V(\rightarrow \ell\ell)$$



- J-type important: same CKM powers as W-type, unflavored final mesons (\neq B-Physics)
- Correct factorization with the introduction of **strong phases** and 'fudge factors'
- Annihilation: only $\rho\rho, \sigma\rho - \sim m_d$, partially reabsorbed in free parameters

Angular observables /

Defined e.g. as

$$I_2 = \int_{-\pi}^{\pi} d\phi \left[\int_{-1}^{-0.5} d \cos \theta_\ell + \int_{0.5}^1 d \cos \theta_\ell - \int_{-0.5}^{0.5} d \cos \theta_\ell \right] \frac{d^5 \Gamma}{dp^2 dq^2 d\Omega}$$

results in

$$I_i = f(\underbrace{C_9^{\text{eff}:P}(q^2), C_9^{\text{eff}:S}(q^2)}_{\text{SM, P- and S- wave}}, \underbrace{C_9^{NP}, C_{10}^{NP}, C_9'^{NP}, C_{10}'^{NP}}_{\text{NP only}}) \times (\text{Long-distance part})$$

Null tests

$$I_5, I_6 \text{ (Forward-backward asymmetry)}, I_7 \stackrel{SM}{=} 0$$

Further integration over θ_h

$$\langle I_i \rangle_- \equiv \left[\int_0^{+1} d \cos \theta_{P_1} - \int_{-1}^0 d \cos \theta_{P_1} \right] I_i, \quad \langle I_i \rangle_+ \equiv \int_{-1}^{+1} d \cos \theta_{P_1} I_i$$

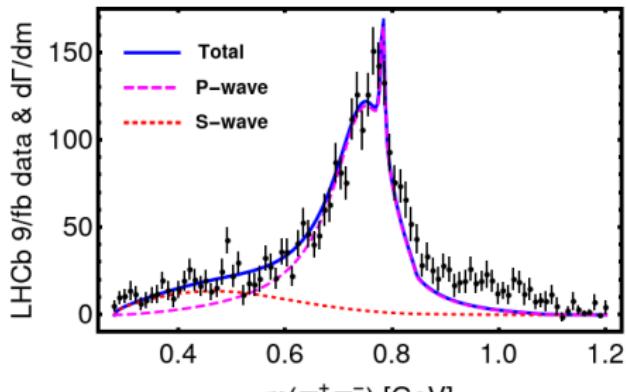
results in observables that either

- depend on the P-wave only: $\langle I_3 \rangle_+, \langle I_6 \rangle_+, \langle I_9 \rangle_+, \langle I_4 \rangle_-, \langle I_5 \rangle_-, \langle I_7 \rangle_-, \langle I_8 \rangle_-$
- depend on the P-wave and the S-wave but not their interference: $\langle I_1 \rangle_+, \langle I_2 \rangle_+$
- depend on the interference of S and P wave (rest of the observables)

CP-averaged ones $\langle S_i \rangle$; no CP violation $\rightarrow \langle S_i \rangle = \langle I_i \rangle_{\pm} (q^2\text{-binned})$

Results (preliminary)

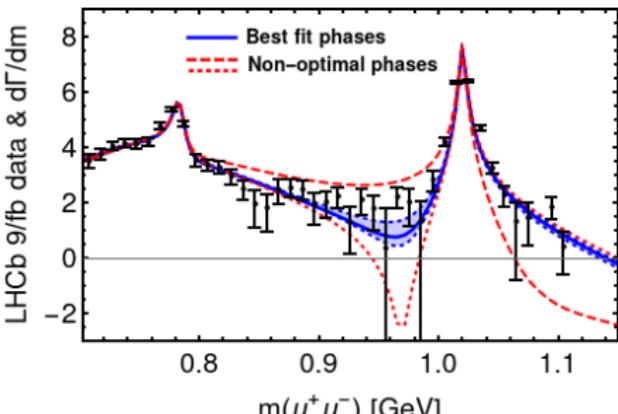
Differential mass distributions: p^2 , q^2



- Double peak from ρ/ω
- We fit the relative phase of $\rho - \omega$: dynamics specific to this decay (not a semileptonic one)

Inclusion of σ is paramount!

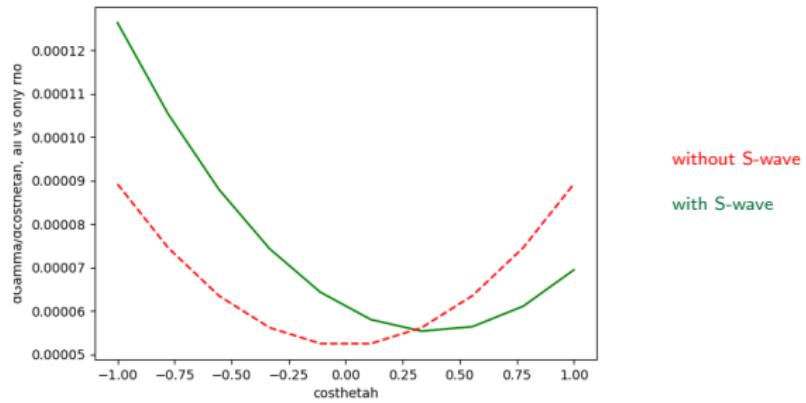
$$\chi^2_{min; w/o \sigma} - \chi^2_{min} = (8.6)^2$$



- Peaks: ρ/ω and ϕ
- We fit for the phase difference of ρ/ω and ϕ ; significantly affects high q^2
- σ contributes to the diff. decay rate by 10 – 35% depending on q^2 (not visible)

- q^2 -binned observables I_2 , I_3 reproduced well ; S-wave contributes more than 30% to I_2
- I_4 not reproducible with current framework

Tests for S-wave: $d\Gamma/d \cos \theta_h$, angular observables



(not yet experimentally provided)

$$\frac{d\Gamma}{d \cos \theta_h} \propto |F_S|^2 + |F_{0,P}|^2 \cos^2 \theta_h + (|F_{||}|^2 + |F_{\perp}|^2) \sin^2 \theta_h \\ + \text{Re}\{F_S F_{0,P}^*\} \cos \theta_h$$

Integration by p^2 regions would further highlight the presence of σ / help determine strong dynamics

Observables $\langle I_2 \rangle_-$, $\langle I_4 \rangle_+$, $\langle I_8 \rangle_+$ calculated to be at a measurable level

Angular observables: NP

Null tests

I_5, I_6 (Forward-backward asymmetry) , $I_7 \stackrel{SM}{=} 0$

Examine the case of $C_{10} \neq 0$, $C_9, C'_9, C'_{10} = 0$

- NP observables sensitive to P-wave only: $\langle I_5 \rangle_-, \langle I_6 \rangle_+, \langle I_7 \rangle_-$
 - ① $\langle I_5 \rangle_-, \langle I_6 \rangle_+$ up to a few % for current upper bounds on C_{10}
 - ② $\langle I_7 \rangle_- \approx 0$ very suppressed with the current long-distance parameterization (same for $\langle I_8 \rangle_-, \langle I_9 \rangle_+$)
- NP observables with S-wave:
 - ① $\langle I_5 \rangle_+$ (S-P interf.) $\approx \mathcal{O}(\langle I_5 \rangle_-)$ (P-wave)
 - ② $\langle I_7 \rangle_+$ (S-P interf.) at the same level as $\langle I_5 \rangle_+, >> \langle I_7 \rangle_-$

Conclusions

Summary

- Rare D decays are dominated by long-distance QCD dynamics; focus on $D^0 \rightarrow \pi^+ \pi^- \mu^+ \mu^-$
- Factorisation-inspired + phases + room for magnitude adjustments
- New: inclusion of S-wave dynamics through the $\sigma = f_0(500)$ resonance, encoding rescattering in the relevant energy range

Results:

- Predictions for $d\Gamma/dp^2, dq^2$ (dihadron/dilepton mass) significantly improved; good agreement with data for most observables; determination of relative strong phases
- S-wave presence very prominent in decay rate over pion angle; accessible observables can probe it further
- S-wave inclusion gives access to complementary NP tests
- Some angular observables not correctly reproduced, regardless of S-wave presence; refinement (e.g. cascade decays?) needed before drawing conclusions regarding NP

Summary

- Rare D decays are dominated by long-distance QCD dynamics; focus on $D^0 \rightarrow \pi^+ \pi^- \mu^+ \mu^-$
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Thank you!

BACKUP

Angular observables - definition

$$I_2 = \int_{-\pi}^{\pi} d\phi \left[\int_{-1}^{-0.5} d \cos \theta_\mu + \int_{0.5}^1 d \cos \theta_\mu - \int_{-0.5}^{0.5} d \cos \theta_\mu \right] \frac{d^5 \Gamma}{dq^2 dp^2 d\vec{\Omega}},$$

$$I_3 = \frac{3\pi}{8} \left[\int_{-\pi}^{-\frac{3\pi}{4}} d\phi + \int_{-\frac{\pi}{4}}^{\frac{\pi}{4}} d\phi + \int_{\frac{3\pi}{4}}^{\pi} d\phi - \int_{-\frac{3\pi}{4}}^{-\frac{\pi}{4}} d\phi - \int_{\frac{\pi}{4}}^{\frac{3\pi}{4}} d\phi \right] \int_{-1}^1 d \cos \theta_\mu \frac{d^5 \Gamma}{dq^2 dp^2 d\vec{\Omega}},$$

$$I_4 = \frac{3\pi}{8} \left[\int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} d\phi - \int_{-\pi}^{-\frac{\pi}{2}} d\phi - \int_{\frac{\pi}{2}}^{\pi} d\phi \right] \left[\int_0^1 d \cos \theta_\mu - \int_{-1}^0 d \cos \theta_\mu \right] \frac{d^5 \Gamma}{dq^2 dp^2 d\vec{\Omega}},$$

$$I_5 = \left[\int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} d\phi - \int_{-\pi}^{-\frac{\pi}{2}} d\phi - \int_{\frac{\pi}{2}}^{\pi} d\phi \right] \int_{-1}^1 d \cos \theta_\mu \frac{d^5 \Gamma}{dq^2 dp^2 d\vec{\Omega}},$$

$$I_6 = \int_{-\pi}^{\pi} d\phi \left[\int_0^1 d \cos \theta_\mu - \int_{-1}^0 d \cos \theta_\mu \right] \frac{d^5 \Gamma}{dq^2 dp^2 d\vec{\Omega}},$$

$$I_7 = \left[\int_0^{\pi} d\phi - \int_{-\pi}^0 d\phi \right] \int_{-1}^1 d \cos \theta_\mu \frac{d^5 \Gamma}{dq^2 dp^2 d\vec{\Omega}},$$

$$I_8 = \frac{3\pi}{8} \left[\int_0^{\pi} d\phi - \int_{-\pi}^0 d\phi \right] \left[\int_0^1 d \cos \theta_\mu - \int_{-1}^0 d \cos \theta_\mu \right] \frac{d^5 \Gamma}{dq^2 dp^2 d\vec{\Omega}},$$

$$I_9 = \frac{3\pi}{8} \left[\int_{-\frac{\pi}{2}}^{-\frac{\pi}{2}} d\phi + \int_0^{\frac{\pi}{2}} d\phi - \int_{-\frac{\pi}{2}}^0 d\phi - \int_{\frac{\pi}{2}}^{\pi} d\phi \right] \int_{-1}^1 d \cos \theta_\mu \frac{d^5 \Gamma}{dq^2 dp^2 d\vec{\Omega}}.$$

Numerical results - comparison to experiment

$\sqrt{q^2}$ region	$\Gamma[10^{-5}]$	$\frac{\Gamma_\sigma}{\Gamma}(\%)$	$(\int dp^2 I_2)_{+, \sigma} / (\int dp^2 I_2)_{+, \text{total}}(\%)$
$r^{(\text{low})}$	0.08	17	26
$r^{(\eta)}$	0.03	15	28
$r^{(\rho:\text{inf})}$	1.09	32	61
$r^{(\rho:\text{sup})}$	0.59	37	74
$r^{(\phi:\text{inf})}$	1.00	27	79
$r^{(\phi:\text{sup})}$	0.85	22	75
$r^{(\text{high})}$	0.01	13	62

$\sqrt{q^2}$ region	$\langle S_2 \rangle(\%)$	$\langle S_3 \rangle(\%)$	$\langle S_4 \rangle(\%)$	$\langle S_2 \rangle [\%]$	$\langle S_3 \rangle [\%]$	$\langle S_4 \rangle [\%]$
				$\langle S_2 \rangle [\%]$	$\langle S_3 \rangle [\%]$	$\langle S_4 \rangle [\%]$
$r^{(\text{low})}$	-15	-2	15	$5 \pm 14 \pm 4$	$-6 \pm 16 \pm 2$	$21 \pm 16 \pm 2$
$r^{(\eta)}$	-17	-4	23	—	—	—
$r^{(\rho:\text{inf})}$	-17	-6	21	$-2.4 \pm 4.1 \pm 1.1$	$-9.1 \pm 4.8 \pm 1.5$	$3.7 \pm 4.9 \pm 1.3$
$r^{(\rho:\text{sup})}$	-17	-7	20	$-10.7 \pm 5.8 \pm 1.1$	$7.7 \pm 6.9 \pm 1.0$	$-4.7 \pm 6.9 \pm 1.5$
$r^{(\phi:\text{inf})}$	-12	-12	25	$-2.0 \pm 3.7 \pm 1.6$	$-17.4 \pm 4.3 \pm 1.5$	$-9.9 \pm 4.3 \pm 3.5$
$r^{(\phi:\text{sup})}$	-10	-13	27	$1.7 \pm 3.4 \pm 1.5$	$-15.3 \pm 4.0 \pm 1.7$	$-18.3 \pm 4.0 \pm 2.5$
$r^{(\text{high})}$	-11	-23	44	—	—	—

New Physics observables

$\sqrt{q^2}$ region	Γ (SM+NP) (10^{-5})	$\langle I_5 \rangle_{(+,S)}(\%)$	$\langle I_5 \rangle_{(-,P)}(\%)$	$\langle I_6 \rangle_{(+,P)}(\%)$	$\langle I_7 \rangle_{(+,S)}(\%)$
0.212-0.525	$0.08 + 0.09 \tilde{C}_{10} ^2$	2.1	-2.4	-1.8	-2.1
0.525-0.565	$0.03 + 0.02 \tilde{C}_{10} ^2$	2.2	-2.6	-2.6	-2.0
0.565-0.78	$1.09 + 0.11 \tilde{C}_{10} ^2$	0.7	0.3	0.4	-0.9
0.78- 0.95	$0.59 + 0.10 \tilde{C}_{10} ^2$	-0.2	2.0	2.7	-0
0.95-1.02	$1.00 + 0.04 \tilde{C}_{10} ^2$	0.3	-0.3	-0.5	0.4
1.02-1.1	$0.85 + 0.03 \tilde{C}_{10} ^2$	0.02	0.5	0.9	-0.1
1.1-1.59	$0.01 + 0.03 \tilde{C}_{10} ^2$	-1.8	4.6	8.3	-0.2

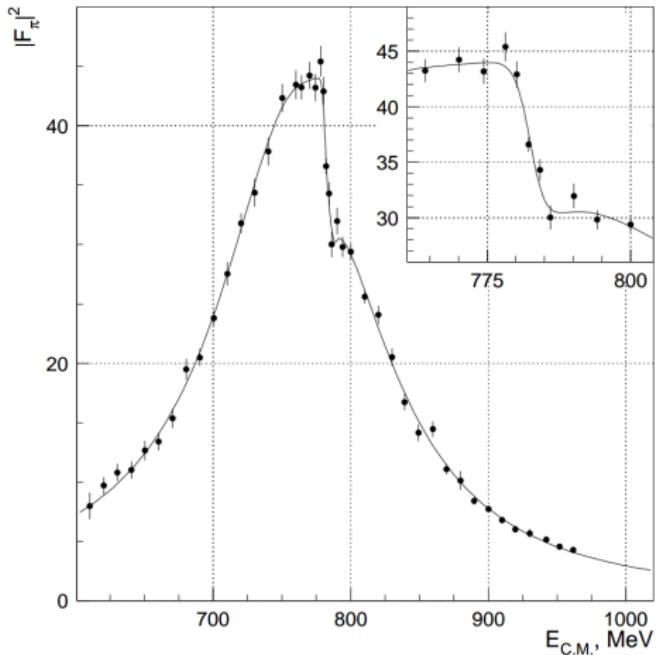
Angular observables given at $\tilde{C}_{10} = 0.43$

SM S-wave observables

For a certain choice of relative phases:

$\sqrt{q^2}$ region (GeV)	$\langle I_2 \rangle_{(+,S)}(\%)$	$\langle I_4 \rangle_{(+,S)}(\%)$	$\langle I_8 \rangle_{(+,S)}(\%)$
0.212-0.525	10	-3	2
0.525-0.565	11	-4	3
0.565-0.78	-1	2	4
0.78- 0.95	-10	6	2
0.95-1.02	14	-9	1
1.02-1.1	-5	3	4
1.1-1.59	-20	15	2

CMD-2



Compare to CMD-2: